ADMX Collaboration Meeting May 4-6, 2022

Low Frequency Project

Pierre Sikivie

- Motivation
- Reentrant Cavity simulations
- Prototype reentrant cavity

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caustic ring



Inner Galactic rotation curve



from Massachusetts-Stony Brook North Galactic Pane CO Survey (Clemens, 1985)

GAIA sky map



L. Duffy and PS, 2008

S. Chakrabarty, Y. Han, A. Gonzalez & PS, 2007.10509





z

ρ

S. Chakrabarty, Y. Han, A. Gonzalez & PS, 2007.10509

with respect to a non-rotating Galactic reference frame

Flow	$d \ [10^{-24} \ { m g/cm^3}]$	$v_ ho^{ m G}~[m km/s]$	$v_{\phi}^{ m G}~[m km/s]$	$v_z^{ m G}~[{ m km/s}]$
Big	20.0	-104.4	509.4	6.1
Little	2.0	-0.2	520.0	4.5
Up	9.6	-115.3	505.1	44.8
Down	8.4	-116.4	505.4	-38.1

 \hat{z} is toward the North Galactic Pole

 $\hat{\phi}$ is in the direction of Galactic rotation

 $\hat{\rho}~$ is in the direction away from the Galactic Center

Uncertainties on the densities are large

but

$$d_{\rm Big} > 10.0 \times 10^{-24} \ \frac{{\rm gr}}{{\rm cc}} = 6 \ \frac{{\rm GeV}}{{\rm cc}}$$

• The evidence for caustic rings requires that the dark matter has vorticity

• Ordinary cold dark matter, such as WIMPs and sterile neutrinos, cannot have vorticity. Their velocity field is curl free.

 Axion dark matter acquires vorticity because it forms a rethermalizing Bose-Einstein condensate. Indeed the lowest energy state for given angular momentum is a state of rotation. The axion dark matter fluid first thermalizes (by gravitational self-interactions) and forms a BEC when the universe is approximately 1 year old (Q. Yang and PS, 2009; O. Erken et al., 2012)

• The horizon at age of one year expands to the size of the largest disk galaxies today.

• Assuming the above is not an accident, we have a new handle on the axion mass.

 $m_a \sim \mu \mathrm{eV} \left(\frac{0.27}{\Omega_a}\right)^{\frac{3}{2}}$

242 MHz

(2π) 120 MHz $\lesssim m_a \lesssim (2\pi)$ 3 GHz



Rybka - ICHEP July 2018

A reentrant cavity search for axion dark matter

ADMX Letter of Intent August 20, 2020

Abstract:

We propose to search for axion dark matter in the frequency range 100 - 625 MHz ($0.4 < m_a < 2.6 \mu eV$) using a reentrant cavity haloscope. First a small prototype will be built and tested in data-taking mode inside the magnet of the 1980's UF Axion Project. Based on this experience, a larger reentrant cavity haloscope will be built for the ADMX site at UW and operated there.

Higher Order Reentrant Post Modes in Cylindrical Cavities

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(Dated: 4 December 2018)

Reentrant cavities are microwave resonant devices employed in a number of different areas of physics. They are appealing due to their simple frequency tuning mechanism, which offers large tuning ranges. Reentrant cavities are, in essence, 3D lumped LC circuits consisting of a conducting central post embedded in a resonant cavity. The lowest order reentrant mode (which transforms from the TM_{010} mode) has been extensively studied in past publications. In this work we show the existence of higher order reentrant post modes (which transform from the TM_{01n} mode family). We characterize these new modes in terms of their frequency tuning, filling factors and quality factors, as well as discuss some possible applications of these modes in fundamental physics tests. The appendix contains a comment on a paper related to this work.

I. INTRODUCTION

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3 Dec 201

[physics.ins-det]

arXiv:1611.08939v4

The cylindrical reentrant cavity is a device that can provide high-Q microwave modes with large tuning ranges. It consists of a metal cavity with a conducting post or ring located centrally within the cavity. The gap between the top of this post or ring and the top of the cavity adjusts the mode frequency and at certain gap spacings traps the electric field within the gap. Such cavity designs have been extensively studied and allow for a standard reentrant mode tuning range on the order of GHz without the need for physically large cavities.¹⁻⁶

Since this structure was first investigated in connection with the development of klystrons, it has been widely used in the construction of microwave oscillators and particle accelerators^{7,8}, and is often chosen as a structure in the study of metamaterials^{9,10}. Some recent work has discussed potential applications in telecommunications¹¹, and detection of gravitational waves^{12,13} and dark matter¹⁴. It is an interesting perspective to view the standard reentrant mode as a perturbed TM_{010} mode. The standard reentrant mode transforms into the empty cavity TM_{010} mode as the central post is removed from the cavity¹⁵. In a similar way, there exist higher order reentrant modes, which can be viewed as perturbed TM_{01n} modes.

In this work we unequivocally demonstrate the existence of these higher order reentrant modes and characterize them in terms of their tuning range, and quality factor. Section II presents a theoretical study of these modes based on finite element analysis, and experimental data follows in Section III.

II. FINITE ELEMENT ANALYSIS OF THE REENTRANT CAVITY



FIG. 1: Cylindrical cross-section of the cavity from the COMSOL model in the r-z plane (encompassing the origin) in cylindrical coordinates. Grey represents the empty cavity space, while the metallic cavity region is represented by grid-lined white area. The height of cavity is ~0.4 m and the diameter of the inner wall is 0.1337 m. The diameter of the central post is 0.024 m.

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Reentrant Cavity





Possible Reentrant Cavity Design



Dielectrically loaded cavity

simulations by Mohamed

see Ben's talk

sapphire

copper

Axially symmetric reentrant cavities simulations

relax2a2.f, relax2a6.f, relax2a7.f, relax2a8.f, relax2a9.f



$$a(\rho, z) = \rho B(\rho, z)$$

 $\vec{B} = B(\rho, z)\hat{\phi}$ $\vec{E} = \frac{i}{\omega} \left[\left(\frac{\partial B}{\partial \rho} + \frac{B}{\rho} \right) \hat{z} - \frac{\partial B}{\partial z} \hat{\rho} \right]$

Reentrant Cavity Form factor

$$C(f) = C_0 \left(\frac{f}{f_0}\right)^2 \alpha(\frac{f}{f_0})$$



f = 230 MHz $f_0 = 570 \text{ MHz}$ $\alpha = 0.4$ C = 0.045

 $f/f_0 = 0.406$, alpha = 0.265





TM0 mode



For posts both of length 25 cm:



For posts of lengths 25 cm on the left, 24 cm on the right:



 $f/f_0 = 0.405$, alpha = 0.304

Sensitivity of an ADMX reentrant cavity search

 $P = g^2 \rho \ B^2 V \ C \ \frac{1}{m} \ Q$ $P = P_0 Q$ $\frac{d \log(f)}{dt} = \frac{1}{f} \frac{1}{(s/n)^2} \left(\frac{P_0}{T_n}\right)^2 \frac{4}{9} Q_a Q_a$

 $\delta v < 70 \text{ m/s} \simeq 2 \times 10^{-7} c$

 $\frac{\delta f}{f} \sim \frac{v \, \delta v}{c^2} < 2 \times 10^{-10}$

For example

$$f = 300 \text{ MHz}$$

 $\delta f < 60 \text{ mHz}$

No mode crossings.

With the ADMX magnet at CENPA, how low can we go?

$$f_0 = 574 \text{ MHz}$$
$$C(f) = 0.4 \left(\frac{f}{f_0}\right)^2$$

$$Q(f) \sim \frac{R}{\delta} \propto f^{\frac{1}{3}}$$

 $T_n \sim f^0$ (constant)

sensitivity
$$\propto \frac{g^2}{m_a} \rho \frac{\sqrt{Q_a}}{s/n} C$$

$$373 \ \left(\frac{f}{f_0}\right)^{2+1+\frac{1}{6}} = 1$$

for Big Flow density of 6 GeV/cc

$$f = 88.5 \text{ MHz}$$

Medres sensitivity for our standard 0.45 GeV/cc density



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Prototype reentrant cavity design by Joe Gleason





















 $f/f_0 = 0.324$, alpha = 0.348













Conclusion:

a low frequency search is "low lying fruit"

Thank you