

- **Exam 3: Friday December 3rd, 8:20pm to 10:20pm**
- **You must go to the following locations based on the 1st letter of your last name:**

A to F	WEIL270 (Weil Hall)
G to M	WM 100 (Williamson Hall)
N to Z	FAB103/105 (Fine Arts B)

- **Review sessions: Tuesday Nov. 30 (Hill) and Thurs. Dec. 2 (Woodard), 6:15 to 8:10pm in NPB 1001 (HERE!)**
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- **Final Exam (cumulative): Tuesday December 14th, 12:30pm to 2:30pm.**
 - **Room assignments: A to K in NPB1001 (in here); L to Z in Norman Hall 137.**
 - **Two more review sessions: Dec. 7 and Dec. 9, 6:15 to 8:10pm in NPB1001 (HERE!)**

Class 38 - Waves I

Chapter 16 - Monday November 29th

- QUICK review of traveling waves
- Wave speed and the wave equation
- HiTT problem(s)
- Energy in traveling waves
- Introduction to wave interference

Reading: pages 413 to 437 (chapter 16) in HRW

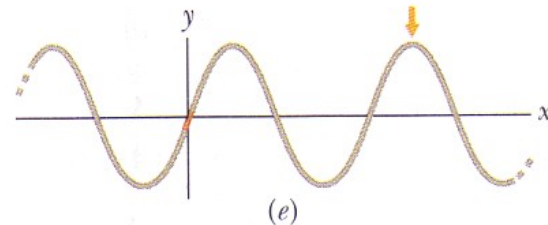
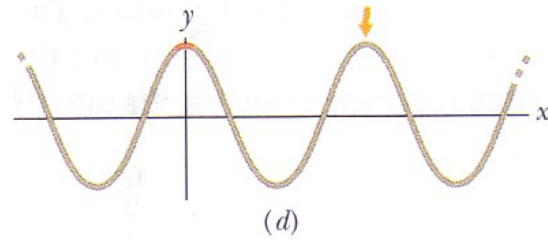
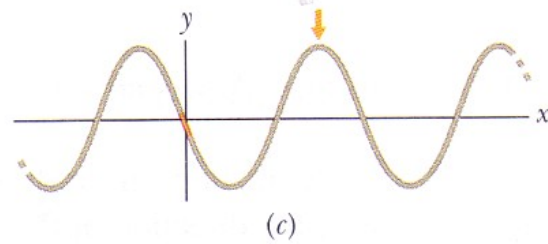
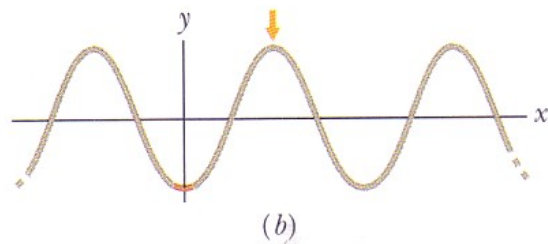
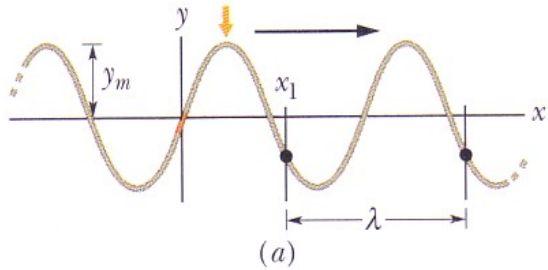
Read and understand the sample problems

Assigned problems from chapter 16 (due Dec. 2nd!):

6, 20, 22, 24, 30, 34, 42, 44, 66, 70, 78, 82

Review - wavelength and frequency

Transverse wave



Displacement $y(x, t) = y_m \sin(kx \mp \omega t + \phi)$

Amplitude y_m

angular wavenumber k

angular frequency ω

Phase ϕ

Phase shift

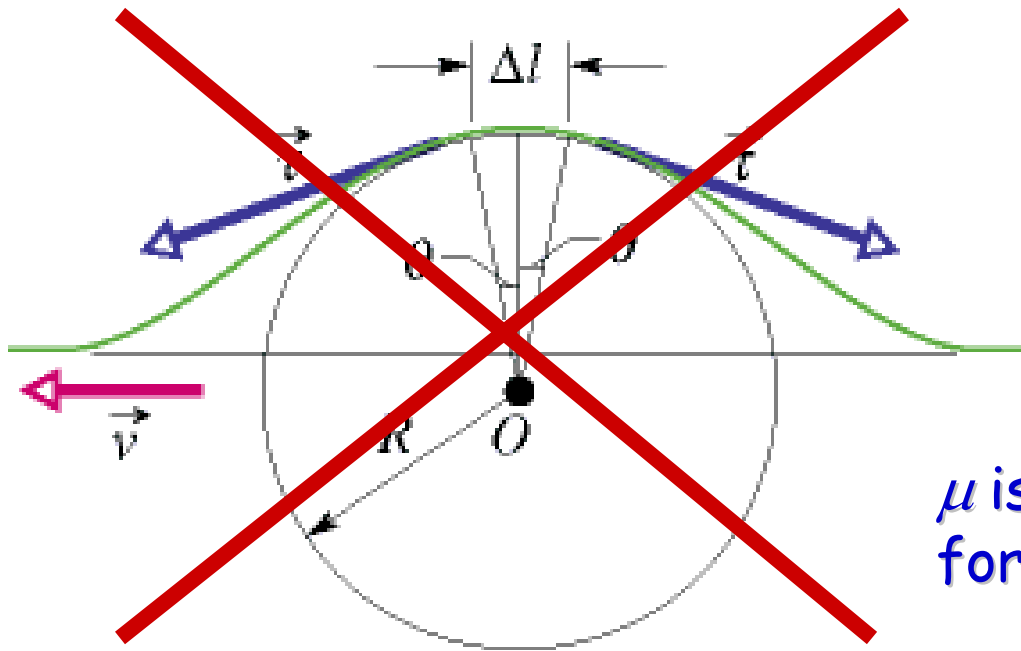
$k = \frac{2\pi}{\lambda}$ k is the angular wavenumber.

$\omega = \frac{2\pi}{T}$ ω is the angular frequency.

frequency $f = \frac{1}{T} = \frac{\omega}{2\pi}$

velocity $v = \pm \frac{\omega}{k} = \pm \frac{\lambda}{T} = \pm f \lambda$

Traveling waves on a stretched string



Dimensional analysis

$$\Delta m = \mu \Delta l$$

μ is the string's linear density, or force per unit length.

- Tension τ provides the restoring force ($\text{kg}\cdot\text{m}\cdot\text{s}^{-2}$) in the string. Without tension, the wave could not propagate.
- The mass per unit length μ ($\text{kg}\cdot\text{m}^{-1}$) determines the response of the string to the restoring force (tension), through Newton's 2nd law.
- Look for combinations of τ and μ that give dimensions of speed ($\text{m}\cdot\text{s}^{-1}$).

$$v = C \sqrt{\frac{\tau}{\mu}}$$

The wave equation

$$\frac{\tau}{\mu} \frac{\partial^2 y}{\partial x^2} = \frac{\partial^2 y}{\partial t^2}$$

Derivation on page 425 of HRW
Or in class notes from Nov. 24th

• General solution:

$$y(x,t) = y_m \sin(kx \pm \omega t) \quad \text{or} \quad y(x,t) = y_m f(kx \pm \omega t)$$

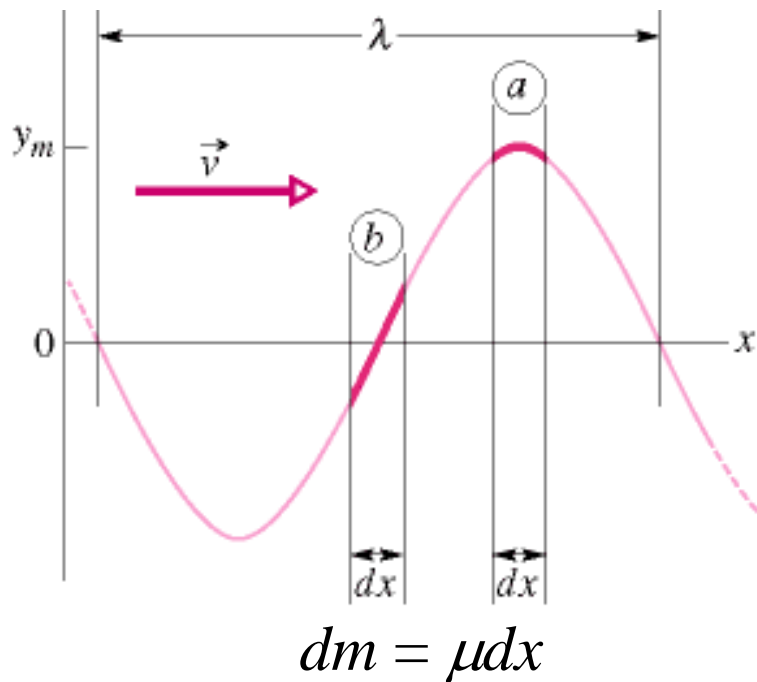
$$\frac{\partial^2 y}{\partial x^2} = -k^2 y(x,t) \quad \frac{\partial^2 y}{\partial t^2} = -\omega^2 y(x,t)$$

$$\Rightarrow -\frac{\tau}{\mu} k^2 = -\omega^2 \quad \text{or} \quad \frac{\omega^2}{k^2} = v^2 = \frac{\tau}{\mu}$$

$$\text{i.e.} \quad v = \sqrt{\frac{\tau}{\mu}}$$

Energy in traveling waves

$$y(x, t) = y_m \sin(kx - \omega t)$$



Similar expression for elastic potential energy

$$\text{Kinetic energy: } dK = \frac{1}{2} dm v_y^2$$

$$v_y = \frac{\partial y}{\partial t} = -\omega y_m \cos(kx - \omega t)$$

$$dK = \frac{1}{2} (\mu dx) (-\omega y_m)^2 \cos^2(kx - \omega t)$$

Divide both sides by dt , where $dx/dt = v_x$

$$\frac{dK}{dt} = \frac{1}{2} \mu v_x \omega^2 y_m^2 \cos^2(kx - \omega t)$$

$$\frac{dU}{dt} = \frac{1}{2} \mu v_x \omega^2 y_m^2 \cos^2(kx - \omega t)$$

$$P_{avg} = 2 \times \frac{1}{2} \mu v \omega^2 y_m^2 \langle \cos^2(kx - \omega t) \rangle = 2 \times \frac{1}{2} \mu v \omega^2 y_m^2 \times \frac{1}{2} = \frac{1}{2} \mu v \omega^2 y_m^2$$

Energy is pumped in an oscillatory fashion down the string

Note: I dropped the subscript on v since it represents the wave speed

The principle of superposition for waves

- It often happens that waves travel simultaneously through the same region, *e.g.*
 - Radio waves from many broadcasters
 - Sound waves from many musical instruments
 - Different colored light from many locations from your TV
- Nature is such that all of these waves can exist without altering each others' motion
- Their effects simply add
- This is a result of the **principle of superposition**, which applies to all **harmonic waves**, *i.e.* waves that obey the linear wave equation

$$v^2 \frac{\partial^2 y}{\partial x^2} = \frac{\partial^2 y}{\partial t^2}$$

- And have solutions: $y(x,t) = y_m f(kx \pm \omega t)$ or $y_m \sin(kx \pm \omega t)$

The principle of superposition for waves

•If two waves travel simultaneously along the same stretched string, the resultant displacement y' of the string is simply given by the summation

$$y'(x,t) = y_1(x,t) + y_2(x,t)$$

where y_1 and y_2 would have been the displacements had the waves traveled alone.

•This is the **principle of superposition**.

Overlapping waves algebraically add to produce a **resultant wave** (or **net wave**).

Overlapping waves do not in any way alter the travel of each other

Link 3

Interference of waves

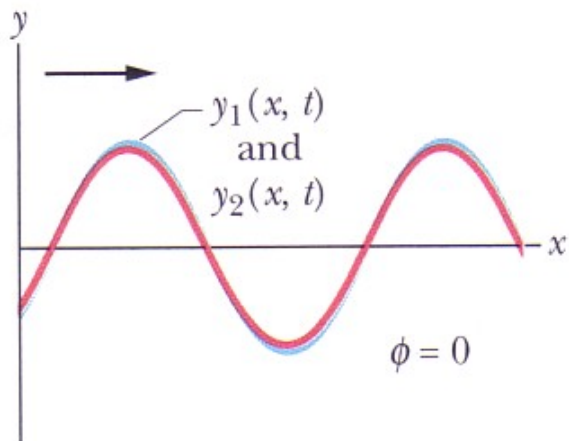
- Suppose two sinusoidal waves with the same frequency and amplitude travel in the same direction along a string, such that

$$y_1 = y_m \sin(kx - \omega t)$$

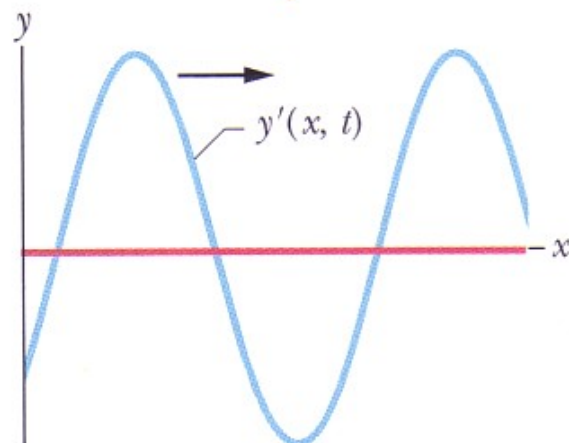
$$y_2 = y_m \sin(kx - \omega t + \phi)$$

- The waves will add.
- If they are in phase (*i.e.* $\phi = 0$), they combine to double the displacement of either wave acting alone.
- If they are out of phase (*i.e.* $\phi = \pi$), they combine to cancel everywhere, since $\sin(\alpha) = -\sin(\alpha + \pi)$.
- This phenomenon is called **interference**.

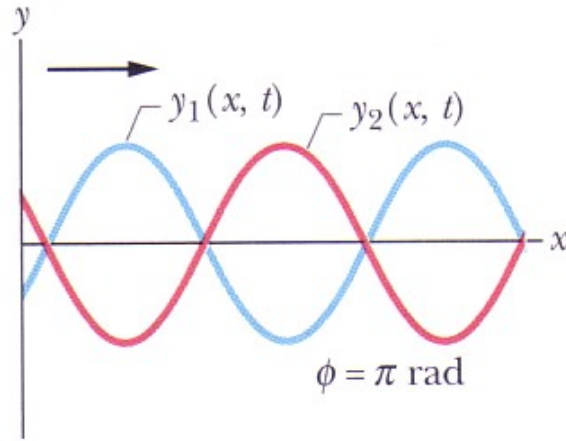
Interference of waves



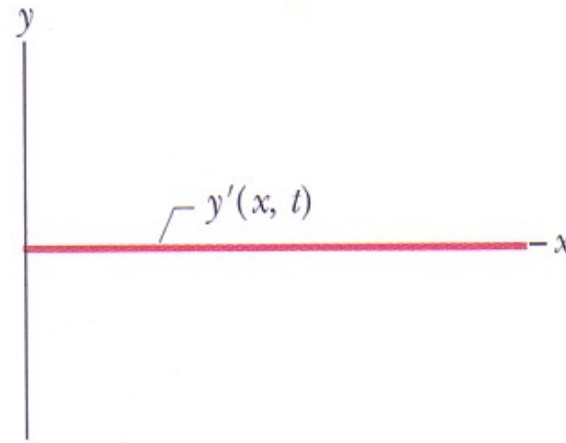
(a)



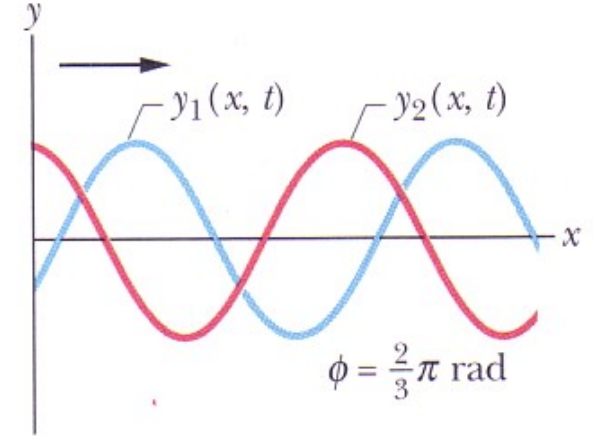
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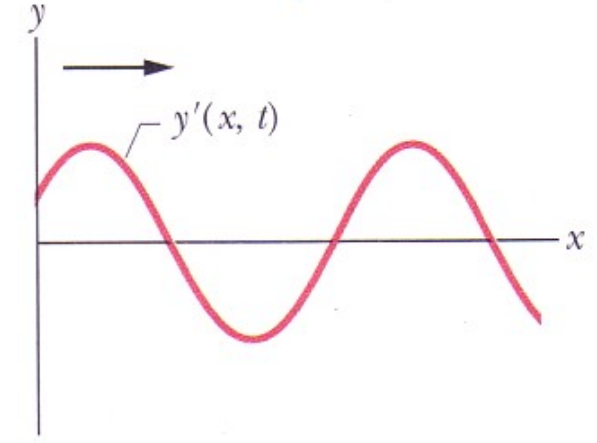
(b)



(e)



(c)



(f)

Interference of waves

•Mathematical proof:

$$y_1 = y_m \sin(kx - \omega t)$$

$$y_2 = y_m \sin(kx - \omega t + \phi)$$

Then:

$$y'(x, t) = y_1(x, t) + y_2(x, t)$$

$$= y_m \sin(kx - \omega t) + y_m \sin(kx - \omega t + \phi)$$

But:

$$\sin \alpha + \sin \beta = 2 \sin \frac{1}{2}(\alpha + \beta) \cos \frac{1}{2}(\alpha - \beta)$$

So:

$$y'(x, t) = \underbrace{[2y_m \cos \frac{1}{2}\phi]}_{\text{Amplitude}} \underbrace{\sin(kx - \omega t + \frac{1}{2}\phi)}_{\text{Wave part}}$$

Phase
shift



Amplitude

Wave part

Interference of waves

$$y'(x, t) = \left[2y_m \cos \frac{1}{2}\phi \right] \sin \left(kx - \omega t + \frac{1}{2}\phi \right)$$

If two sinusoidal waves of the same amplitude and frequency travel in the same direction along a stretched string, they interfere to produce a resultant sinusoidal wave traveling in the same direction.

- If $\phi = 0$, the waves interfere **constructively**, $\cos \frac{1}{2}\phi = 1$ and the wave amplitude is $2y_m$.
- If $\phi = \pi$, the waves interfere **destructively**, $\cos(\pi/2) = 0$ and the wave amplitude is 0, *i.e.* no wave at all.
- All other cases are intermediate between an amplitude of 0 and $2y_m$.
- Note that the phase of the resultant wave also depends on the phase difference.