

Free particle: ($V=0$)

$$-\frac{\hbar^2}{2m} \frac{d^2 u}{dr^2} + \frac{\hbar^2 \ell(\ell+1)}{2mr^2} u = Eu, \quad \boxed{R(r) = u(r)/r}$$

$$\frac{d^2 u}{dr^2} - \frac{\ell(\ell+1)}{r^2} u + \frac{2mE}{\hbar^2} u = 0$$

$$\text{Let } k = \sqrt{\frac{2mE}{\hbar^2}} \quad (E > 0)$$

$$\text{Then } \frac{d^2 u}{dr^2} - \frac{\ell(\ell+1)}{r^2} u + k^2 u = 0.$$

$$\text{Let } \rho = kr$$

$$\rightarrow \frac{d^2 u}{d\rho^2} - \frac{\ell(\ell+1)}{\rho^2} u + u = 0$$

$$\text{and } \frac{1}{\rho} \frac{d^2(\rho R)}{d\rho^2} - \frac{\ell(\ell+1)}{\rho^2} R + R = 0$$

For $l=0$ we know the solution:

$$\frac{d^2 u}{d\rho^2} + u = 0 \rightarrow u(\rho) = \sin(\rho), \cos(\rho)$$

$$R(\rho) = \frac{\sin(\rho)}{\rho}, \frac{\cos(\rho)}{\rho}$$

For $l \neq 0$ the solutions ^{for R} are spherical Bessel functions: (Table 4.4)

$$l=0$$

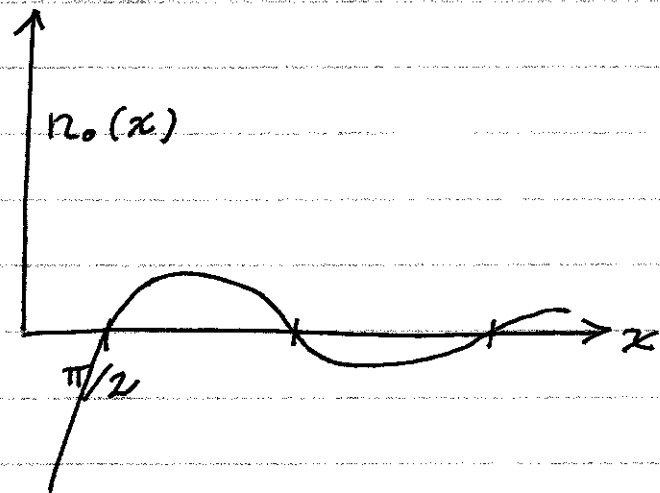
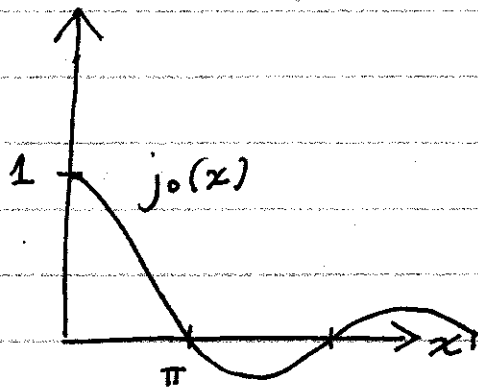
$$j_0(x) = \frac{\sin x}{x}$$

$$n_0(x) = -\frac{\cos(x)}{x}$$

$$l=1$$

$$j_1(x) = \frac{\sin x}{x^2} - \frac{\cos x}{x}$$

$$n_1(x) = -\frac{\cos(x)}{x^2} - \frac{\sin x}{x}$$



The $j_l(x)$ will be non-singular as $x \rightarrow 0$:

$$j_l(x) \sim x^l \quad \text{as } x \rightarrow 0$$

$$n_l(x) \sim \frac{1}{x^{l+1}} \quad \text{as } x \rightarrow 0$$

Infinite spherical well:

$$V(r) = 0 \quad r < a$$

$$V(r) = \infty \quad r > a$$

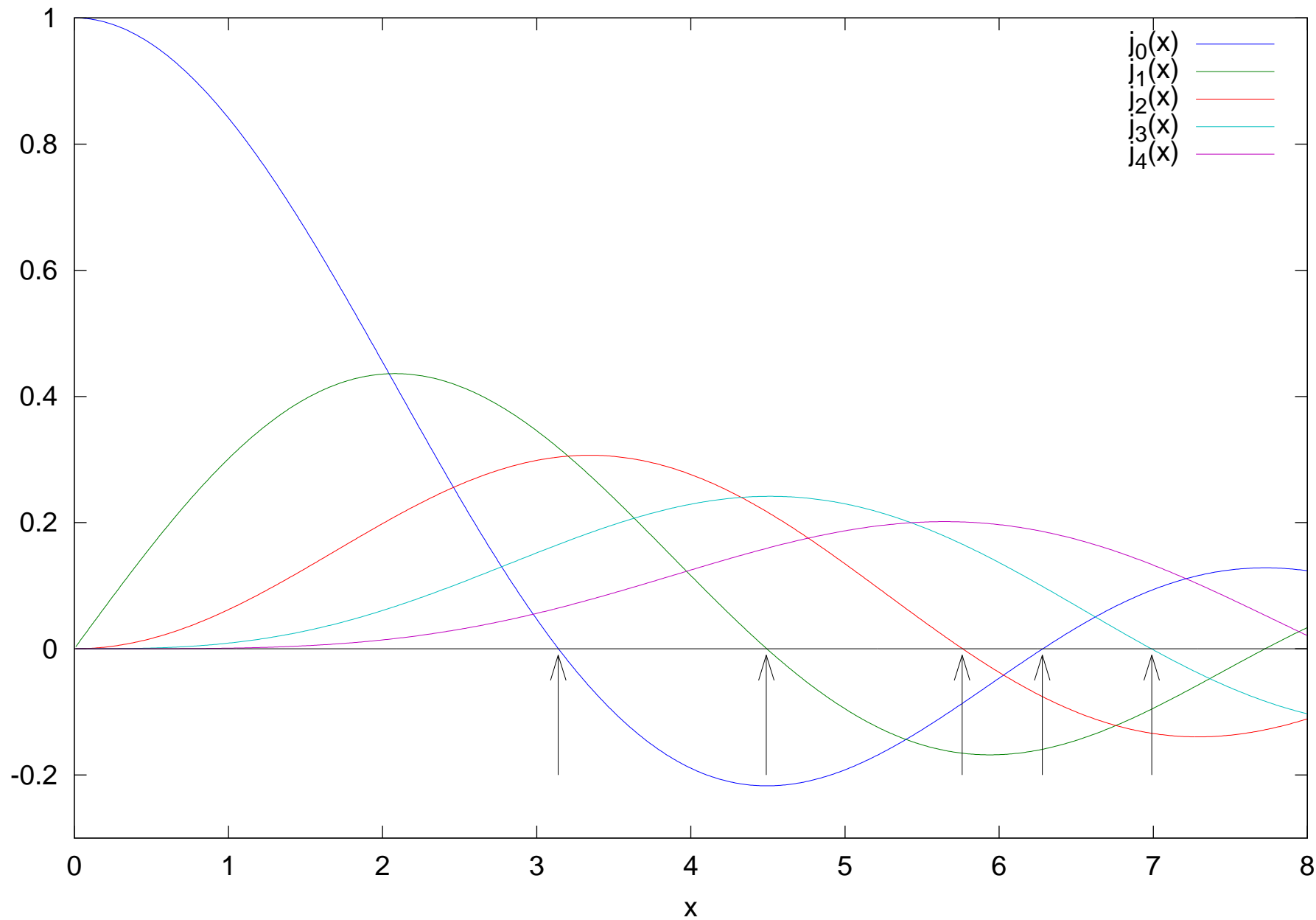
$$R(a) = 0, \quad R(r < a) \propto j_l(ka) \rightarrow j_l(ka) = 0$$

$$\text{with } E = \frac{\hbar^2 k^2}{2m}$$

The zeros of the first few j_l are plotted on the following page:

①	first zero: $j_0(ka) = 0$, $ka \approx 3.14$	$\swarrow \pi$	degeneracy $2 \times 1 = 2$
②	2nd zero: $j_1(ka) = 0$, $ka \approx 4.49$		$2 \times 3 = 6$
③	3rd zero: $j_2(ka) = 0$, $ka \approx 5.76$		$2 \times 5 = 10$
④	4th zero: $j_0(ka) = 0$, $ka \approx 6.28$	$\swarrow 2\pi$	$2 \times 1 = 2$
⑤	5th zero: $j_3(ka) = 0$, $ka \approx 6.99$		$2 \times 7 = 14$
			\uparrow spin

First 5 Spherical Bessel Functions as Defined in Book



In the nuclear shell model one imagines filling a well like this one with nucleons (protons & neutrons).

When a shell is full, the nucleus is particularly stable because it requires extra energy to add another nucleon. This is analogous to why the noble gases (He, Ne, Ar, ...) are not chemically reactive.

The number of nucleons when a filled shell occurs are the so-called magic numbers.

<u>shells filled</u>	<u>Infinite 3D well</u>	<u>Observed (by stability,)</u> <i>more tightly bound</i>
1	$2 = 2$	2
1, 2	$2 + 6 = 8$	8
1, 2, 3	$2 + 6 + 10 = 18$	20
1, 2, 3, 4	$2 + 6 + 10 + 2 = 20$	28

With refinements to the potential & including some interactions one can obtain the observed magic numbers.