## Proof of the Uncertainty Principle

## Introduction

This is a simplified proof of the uncertainty principle. We will do a more general proof later, but I think it is useful to do a proof of a special case now if the proof is transparent. At the end of this document I show how this special case can be generalized to include all wave functions.

## Special Case

The special case we consider is that the expectation value of the position and momentum are zero. In this case the uncertainty principle reduces to

$$
\begin{aligned}
& (\Delta x)^{2}=\left\langle x^{2}\right\rangle-\langle x\rangle^{2}=\left\langle x^{2}\right\rangle=\int d x \psi^{*} x^{2} \psi=\int d x(x \psi)(x \psi) \\
& (\Delta p)^{2}=\left\langle p^{2}\right\rangle-\langle p\rangle^{2}=\left\langle p^{2}\right\rangle=\int d x \psi^{*}\left(-\hbar^{2} \frac{d^{2}}{d x^{2}}\right) \psi=\hbar^{2} \int d x\left(\frac{d \psi}{d x}\right)\left(\frac{d \psi}{d x}\right) \\
& \Delta x \Delta p \geq \frac{\hbar}{2}
\end{aligned}
$$

## Cauchy-Schwarz Inequality

You are familiar with this inequality from three dimensional vectors. If $\vec{A}$ and $\vec{B}$ are two three dimensional vectors, then $\vec{A} \cdot \vec{B}=|\vec{A}||\vec{B}| \cos (\theta)$ and in particular

$$
\begin{equation*}
|\vec{A} \cdot \vec{B}| \leq|\vec{A}||\vec{B}| \tag{1}
\end{equation*}
$$

This is the Cauchy-Schwarz inequality. Let us prove it in a way that generalizes to the case at hand.

The basic idea is that there is a part of $\vec{A}$ which is in the direction of $\vec{B}$ and a part of $\vec{A}$ which is perpendicular to $\vec{B}$. Thus, we can break $\vec{A}$ up into two vectors:

$$
\begin{align*}
\vec{A} & =\vec{A}_{\|}+\vec{A}_{\perp}  \tag{2}\\
\vec{B} \cdot \vec{A}_{\perp} & =0 . \tag{3}
\end{align*}
$$

The parallel and perpendicular vectors are given by

$$
\begin{align*}
\vec{A}_{\|} & =\frac{\vec{B}}{|\vec{B}|}\left(\frac{\vec{A} \cdot \vec{B}}{|\vec{B}|}\right)  \tag{4}\\
\vec{A}_{\perp} & =\vec{A}-\vec{A}_{\|} . \tag{5}
\end{align*}
$$

Now because $\vec{A}_{\|}$and $\vec{A}_{\perp}$ are orthogonal

$$
\begin{equation*}
|\vec{A}|^{2}=\vec{A}_{\|} \cdot \vec{A}_{\|}+\vec{A}_{\perp} \cdot \vec{A}_{\perp} \tag{6}
\end{equation*}
$$

This is basically the Pythagorean theorem. The length of vector $\vec{A}$ is greater than or equal to the length of vector $\vec{A}_{\|}$.

$$
\begin{equation*}
|\vec{A}|^{2} \geq \vec{A}_{\|} \cdot \vec{A}_{\|}=\frac{(\vec{A} \cdot \vec{B})^{2}}{|\vec{B}|^{2}} \tag{7}
\end{equation*}
$$

This is the Cauchy-Schwarz inequality:

$$
\begin{equation*}
|\vec{A}|^{2}|\vec{B}|^{2} \geq(\vec{A} \cdot \vec{B})^{2} \tag{8}
\end{equation*}
$$

## Cauchy-Schwarz inequality for functions

We will cover the results of this section rigorously in approximately a month. Thus, if this does not live up to your level of rigor, just wait until then. Consider two functions: $f(x)$ and $g(x)$. We can define a kind of dot product for these functions as follows

$$
\begin{align*}
" f \cdot f " & =\int d x f^{*}(x) f(x)  \tag{9}\\
" g \cdot g " & =\int d x g^{*}(x) g(x)  \tag{10}\\
" f \cdot g " & =\int d x f^{*}(x) g(x) . \tag{11}
\end{align*}
$$

The complex conjugate is important here so that " $f \cdot f$ " is greater than or equal to zero. Following a line of reasoning exactly the same as in the previous section we can prove the Cauchy-Schwarz identity in this case

$$
\begin{equation*}
\left|\int d x f^{*}(x) g(x)\right|^{2} \leq\left(\int d x f^{*}(x) f(x)\right)\left(\int d x g^{*}(x) g(x)\right) \tag{12}
\end{equation*}
$$

## Uncertainty Principle

Let $f(x)=x \psi(x)$ and $g(x)=d \psi / d x$. Then

$$
\begin{align*}
\int d x f^{*}(x) f(x) & =\int d x|x \psi(x)|^{2}=(\Delta x)^{2}  \tag{13}\\
\int d x g^{*}(x) g(x) & =\int d x\left|\frac{d \psi}{d x}\right|^{2}=\frac{(\Delta p)^{2}}{\hbar^{2}}  \tag{14}\\
\int d x f^{*}(x) g(x) & =\int d x\left(x \psi^{*}(x)\right)\left(\frac{d \psi}{d x}\right)  \tag{15}\\
\int d x g^{*}(x) f(x) & =\int d x(x \psi(x))\left(\frac{d \psi^{*}}{d x}\right)  \tag{16}\\
\operatorname{Re}\left\{\int d x f^{*}(x) g(x)\right\} & =\int d x \frac{1}{2}\left(\psi^{*} \frac{d \psi}{d x}+\frac{d \psi}{d x} \psi^{*}\right)=\int d x x \frac{1}{2} \frac{d|\psi|^{2}}{d x}=-\frac{1}{2} . \tag{17}
\end{align*}
$$

Because

$$
\begin{equation*}
\left|\operatorname{Re}\left\{\int d x f^{*}(x) g(x)\right\}\right| \leq\left|\int d x f^{*}(x) g(x)\right| \tag{18}
\end{equation*}
$$

applying the Cauchy-Schwarz identity yields

$$
\begin{equation*}
\left(-\frac{1}{2}\right)^{2} \leq(\Delta x)^{2} \frac{(\Delta p)^{2}}{\hbar^{2}} \tag{19}
\end{equation*}
$$

which is the uncertainty relation.

## General Case

At the beginning of this document assumed that the expectation values of position and momentum were zero. Suppose that the expectation value of the position is ${\underset{\sim}{x}}_{o}$ and the expectation value of the momentum is $p_{o}$. We can define a new wave function, $\tilde{\psi}$,

$$
\begin{equation*}
\tilde{\psi}(x)=e^{-i p_{o} x / \hbar} \psi\left(x+x_{o}\right) \tag{20}
\end{equation*}
$$

which has zero expectation values for both position and momentum, and which also has the same $\Delta x$ and $\Delta p$ as $\psi$. Thus, we can without loss of generality assume that the expectation values of position and momentum are zero.

