

EXCESS BROADENING
DUE TO DIFFUSION IN
THE ABSORBER

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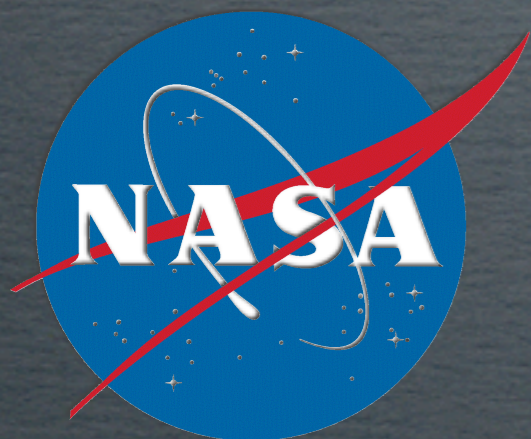
TAREK SAAB
DAVID ELAM

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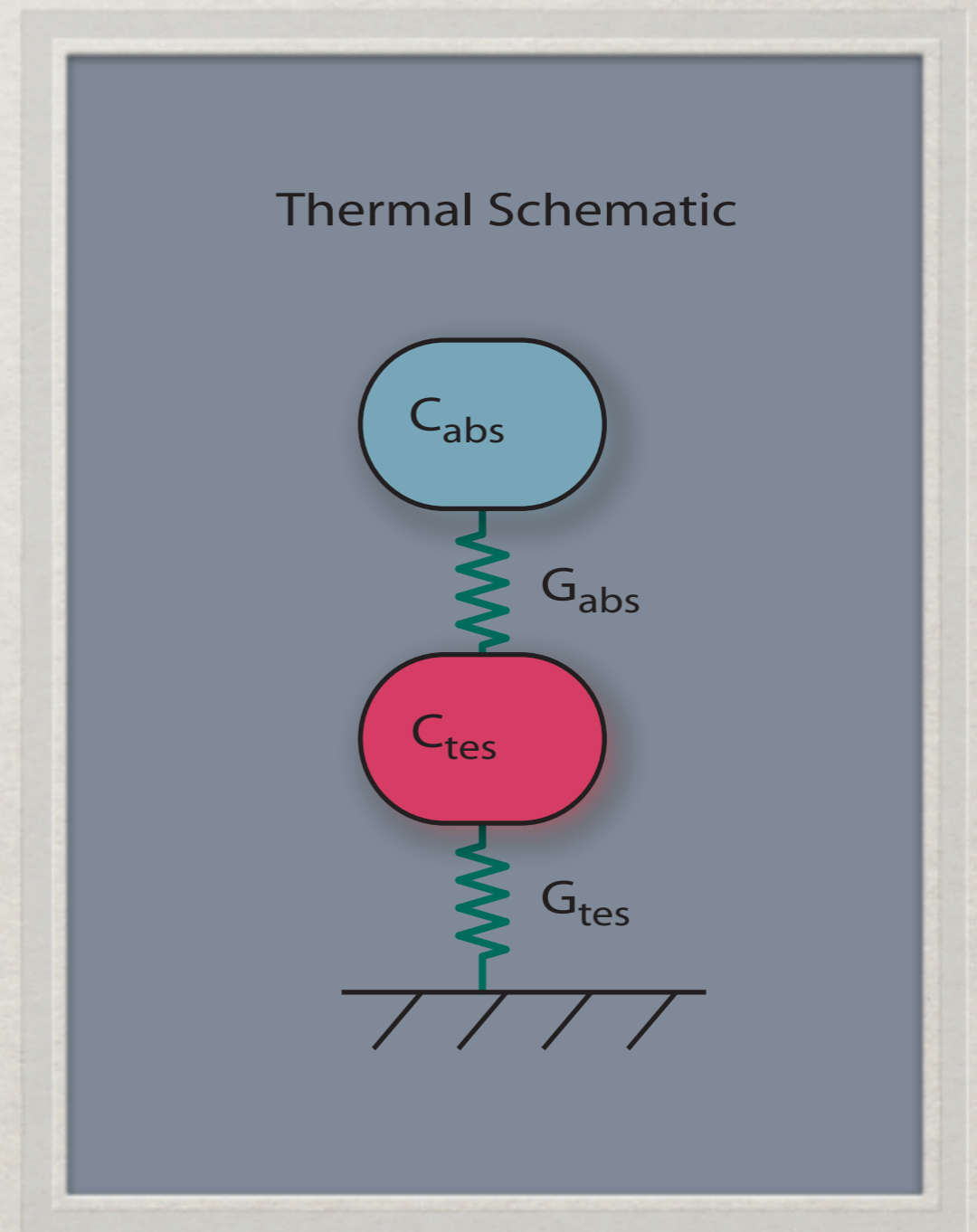
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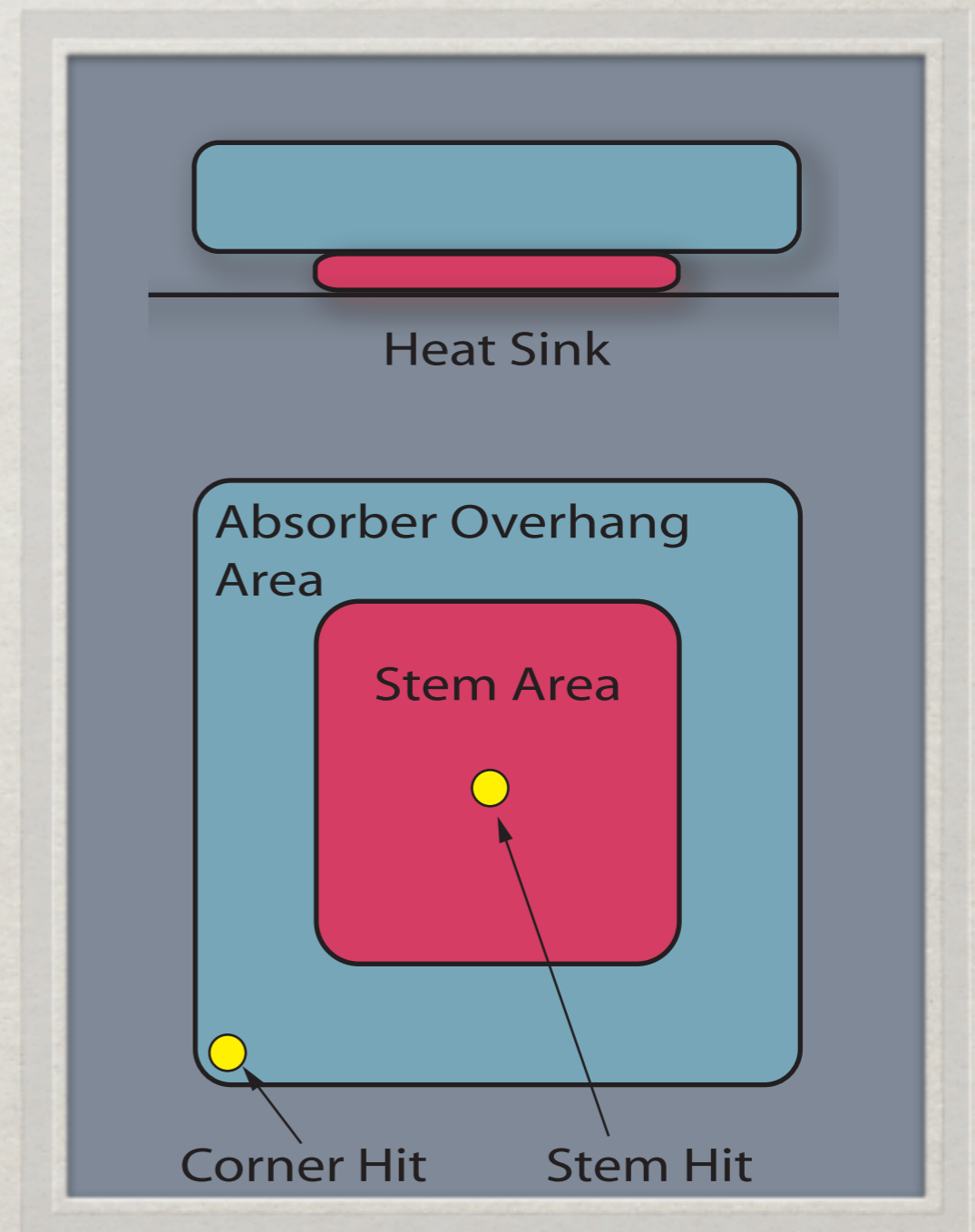
THE IDEAL MICROCALORIMETER

- ✱ Absorber connected to TES connected to heat sink
- ✱ Assume instant thermalization in Absorber / TES
- ✱ Response / Resolution is determined by $C_{\text{abs}} / G_{\text{abs}} / C_{\text{tes}} / G_{\text{tes}}$ and the $R(T,I)$



PRACTICAL μ CAL : THE MUSHROOM DESIGN

- ✱ Cantilevered design driven by practical considerations :
- ✱ Maximizing focal plan area
- ✱ Preventing wiring and substrate hits



PRACTICAL μ CAL : PHYSICAL CONSIDERATIONS

- ✻ The ideal microcalorimeter model no longer applies
- ✻ Diffusion time scales in the absorber, comparable to the thermal time scales, can affect detector response
- ✻ Different “path lengths” from interaction point to heat sink result in position dependent pulse shape
- ✻ This leads to a degradation in resolution that increases linearly with energy

CHARACTERIZING THE PROBLEM

- ✱ We did a numerical simulation of diffusion in a mushroom absorber

$$\frac{\partial E_{\text{abs}}(x, y, t)}{\partial t} = D \nabla^2 E_{\text{abs}}(x, y, t) - \frac{G_{\text{abs}}}{C_{\text{abs}}} E_{\text{abs}}(x, y, t) + \frac{G_{\text{abs}}}{C_{\text{tes}}} E_{\text{tes}}(x, y, t)$$

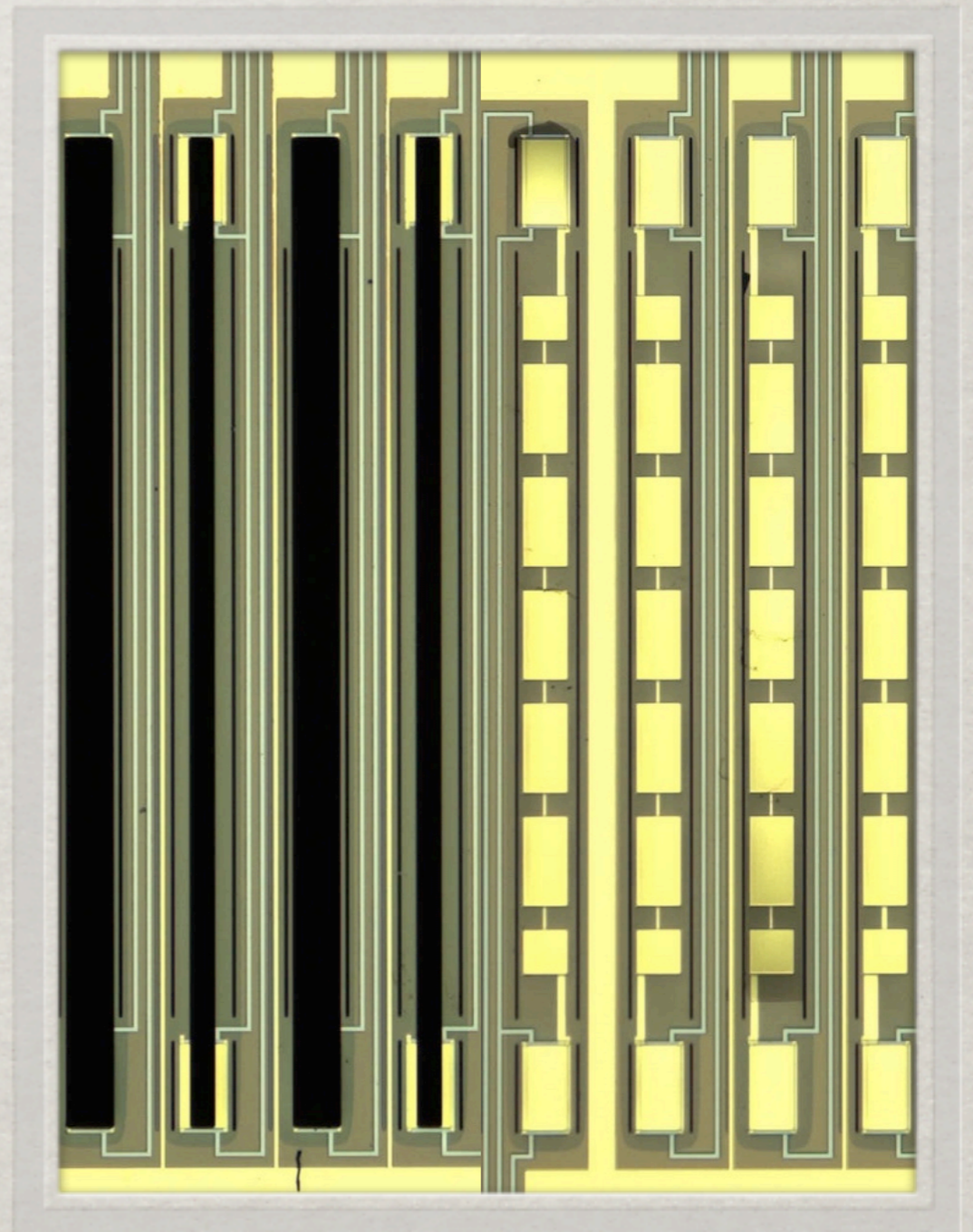
DETAILS OF DIFFUSION

Solved diffusion/detector model numerically.

- ✻ Permits arbitrary definition of device geometry and edge conditions
- ✻ Used Forward Time Centered Space differencing method
- ✻ Method allows us to solve for a given time step across the entire spacial grid simultaneously

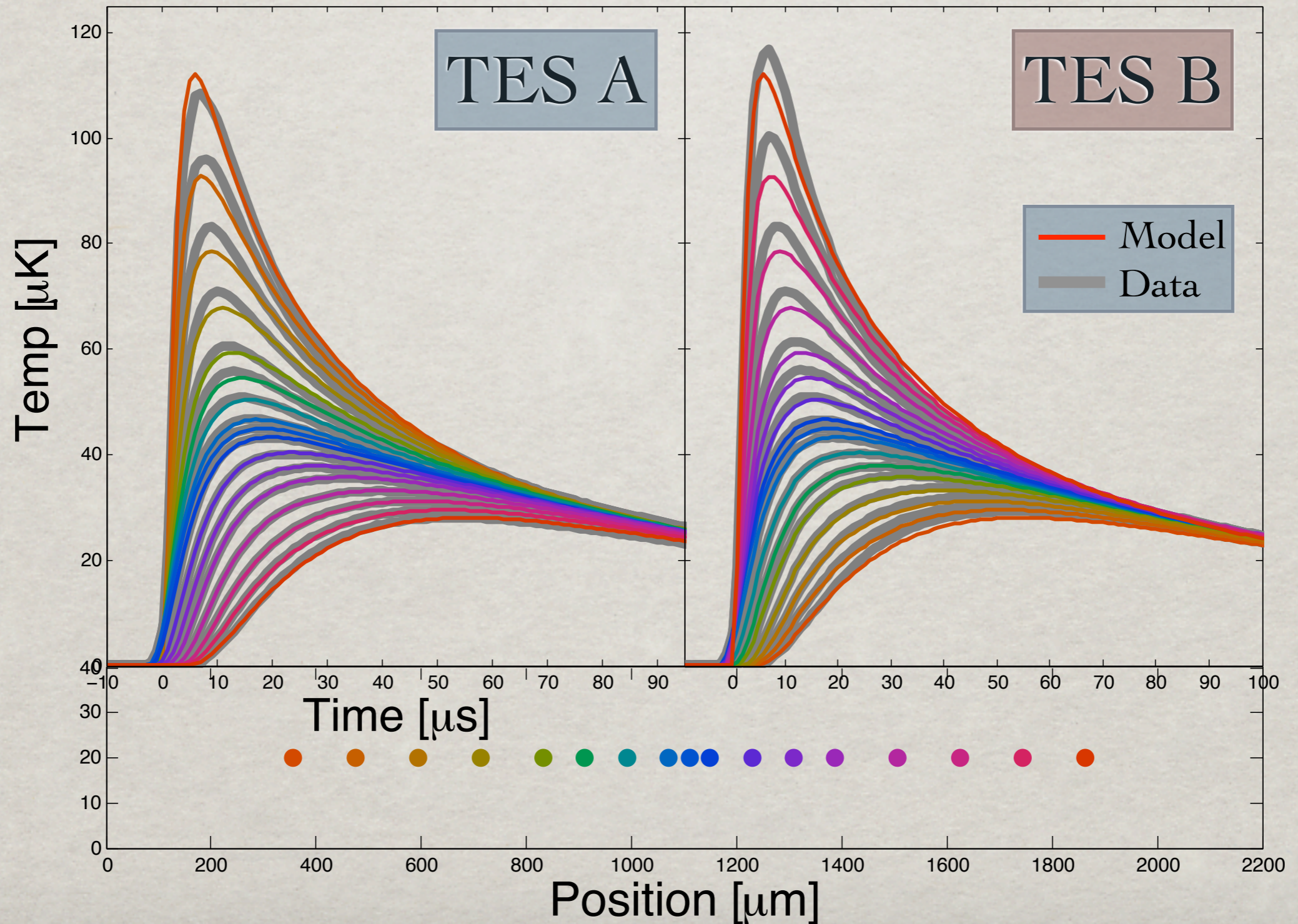
VERIFICATION OF DIFFUSION MODEL

- ✻ Applied the diffusion model to a continuous PoST geometry
- ✻ Compared to data obtained from PoST device
- ✻ Defined pixels based on equal count bins



POST PULSES

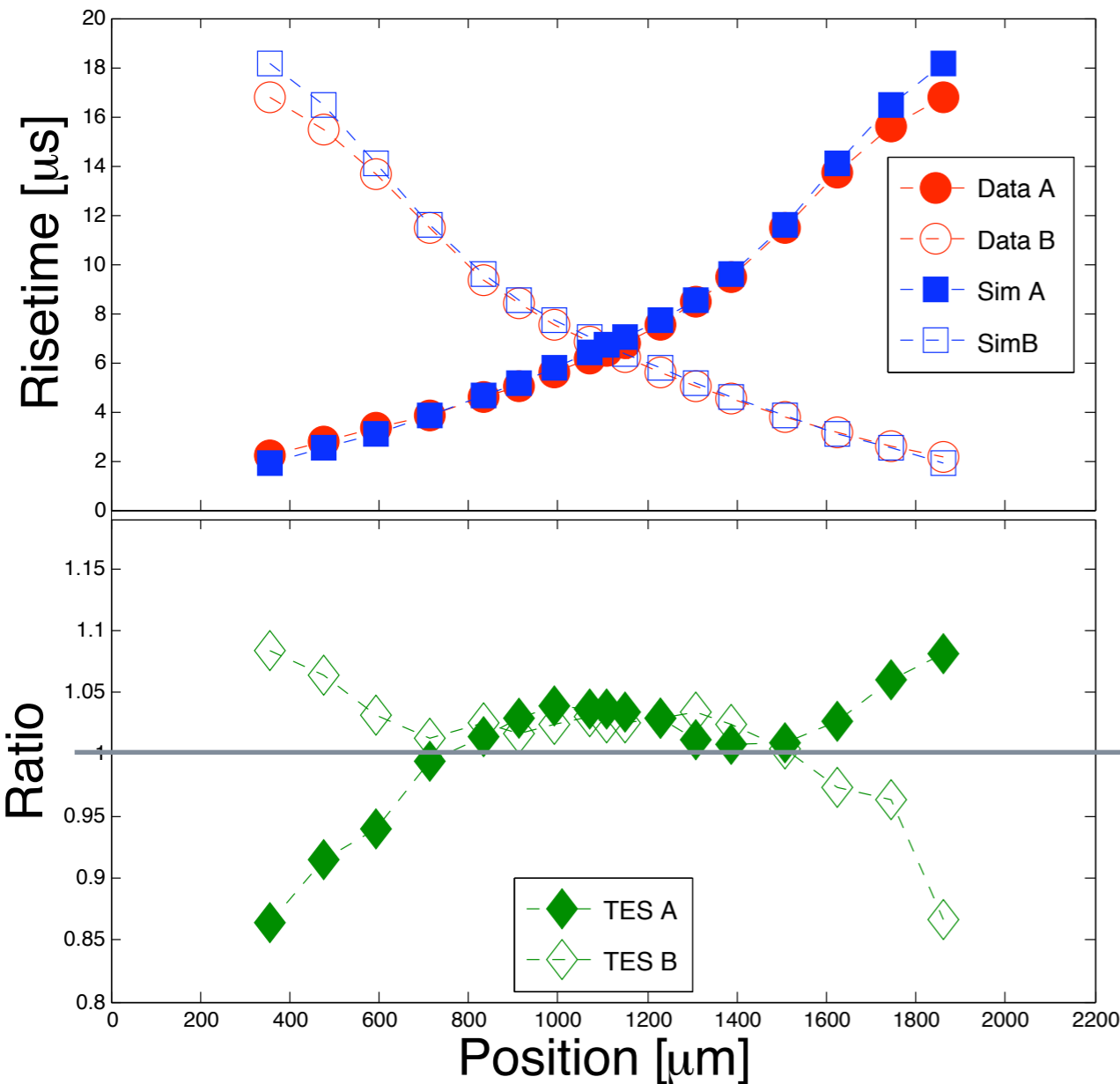
Diffusion Modeling of a PoST



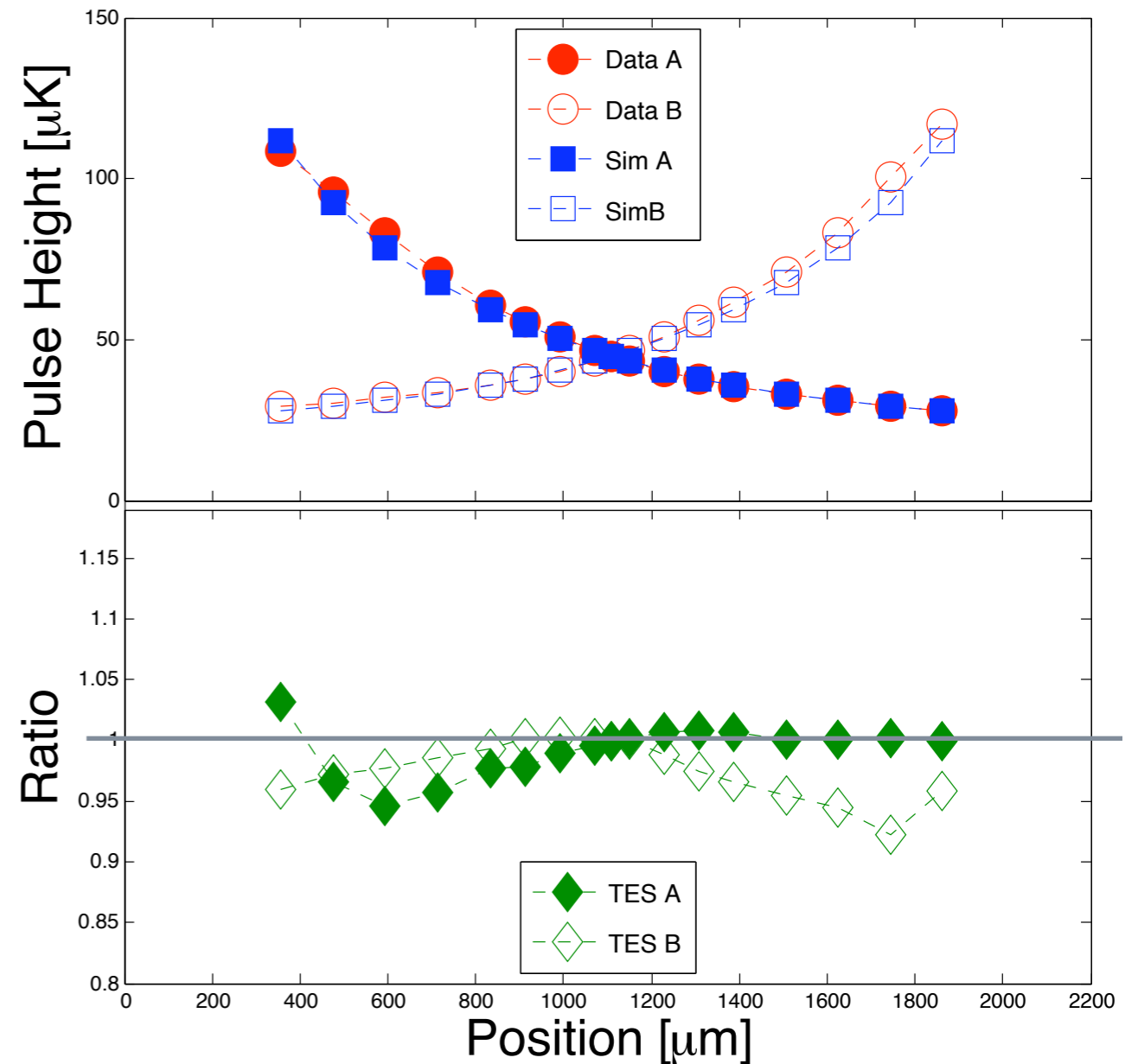
VALIDITY OF THE MODEL

- Comparison of risetime and pulse height shows good agreement between data and model and indicate $D \sim 3.2 \times 10^4 \mu\text{m}^2/\mu\text{s}$

Risetime vs Position

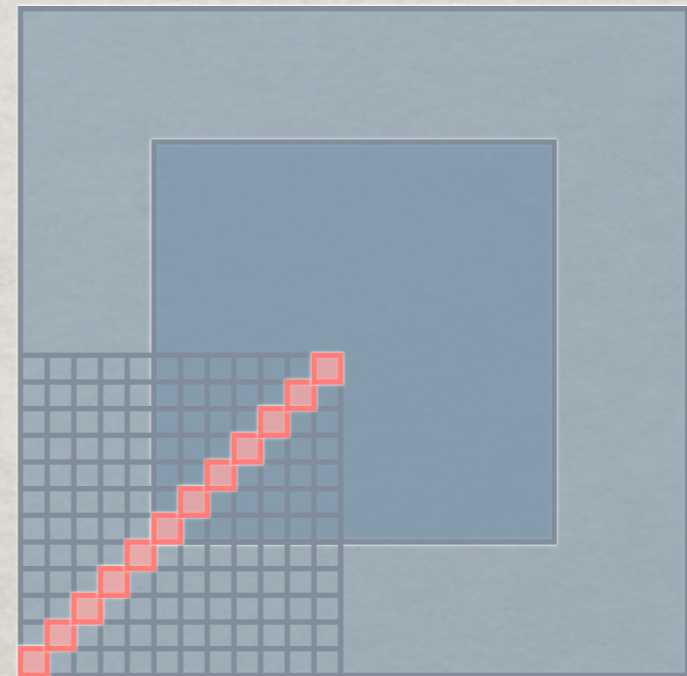


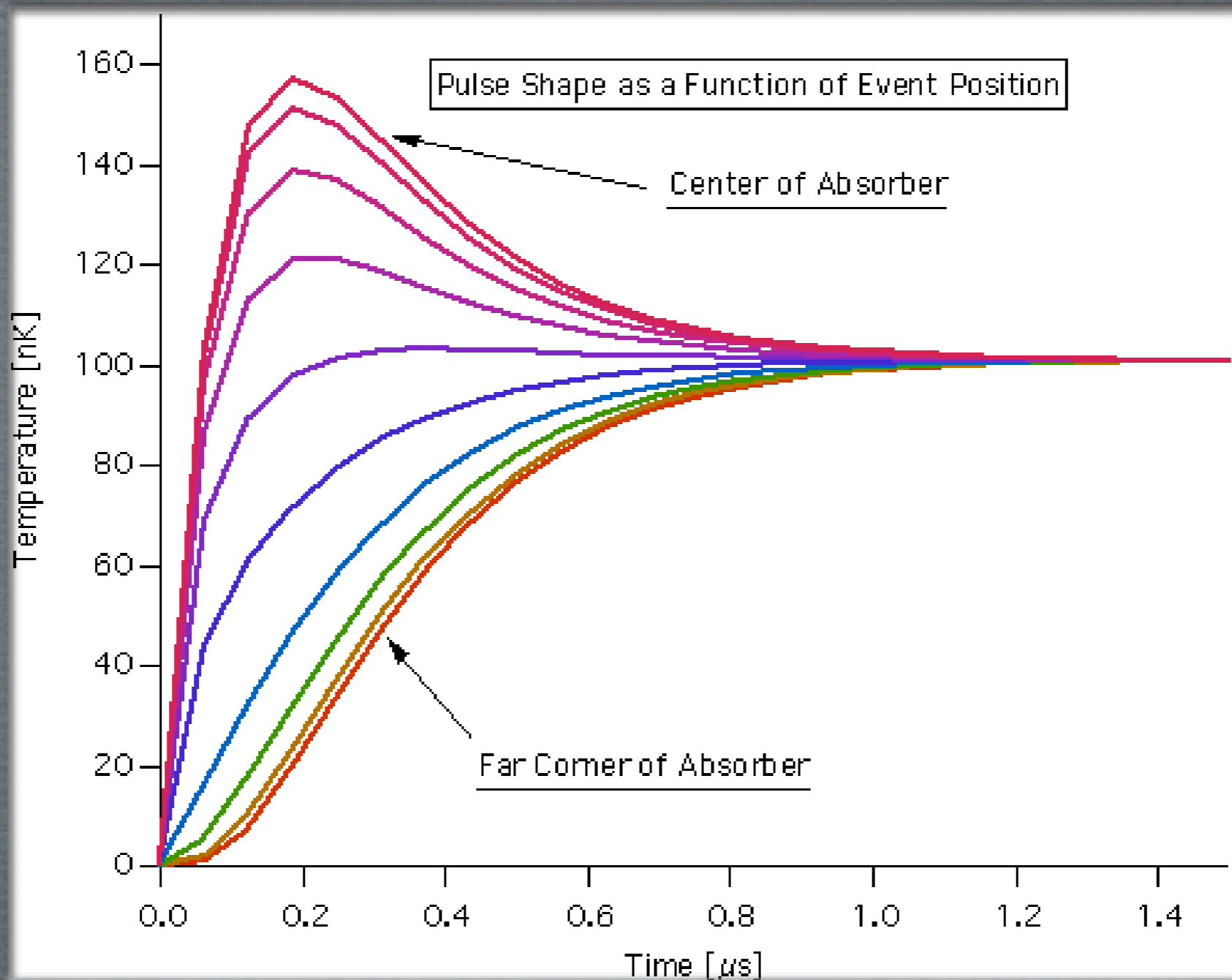
Pulse Height vs Position



APPLICATION TO A MUSHROOM UCAL

- ☼ Simulated pulses across the surface of a $250 \times 250 \mu\text{m}$ absorber, attached to a $150 \times 150 \mu\text{m}$ TES.
- ☼ Device parameters were chosen to give a nominal energy resolution of 2 eV.





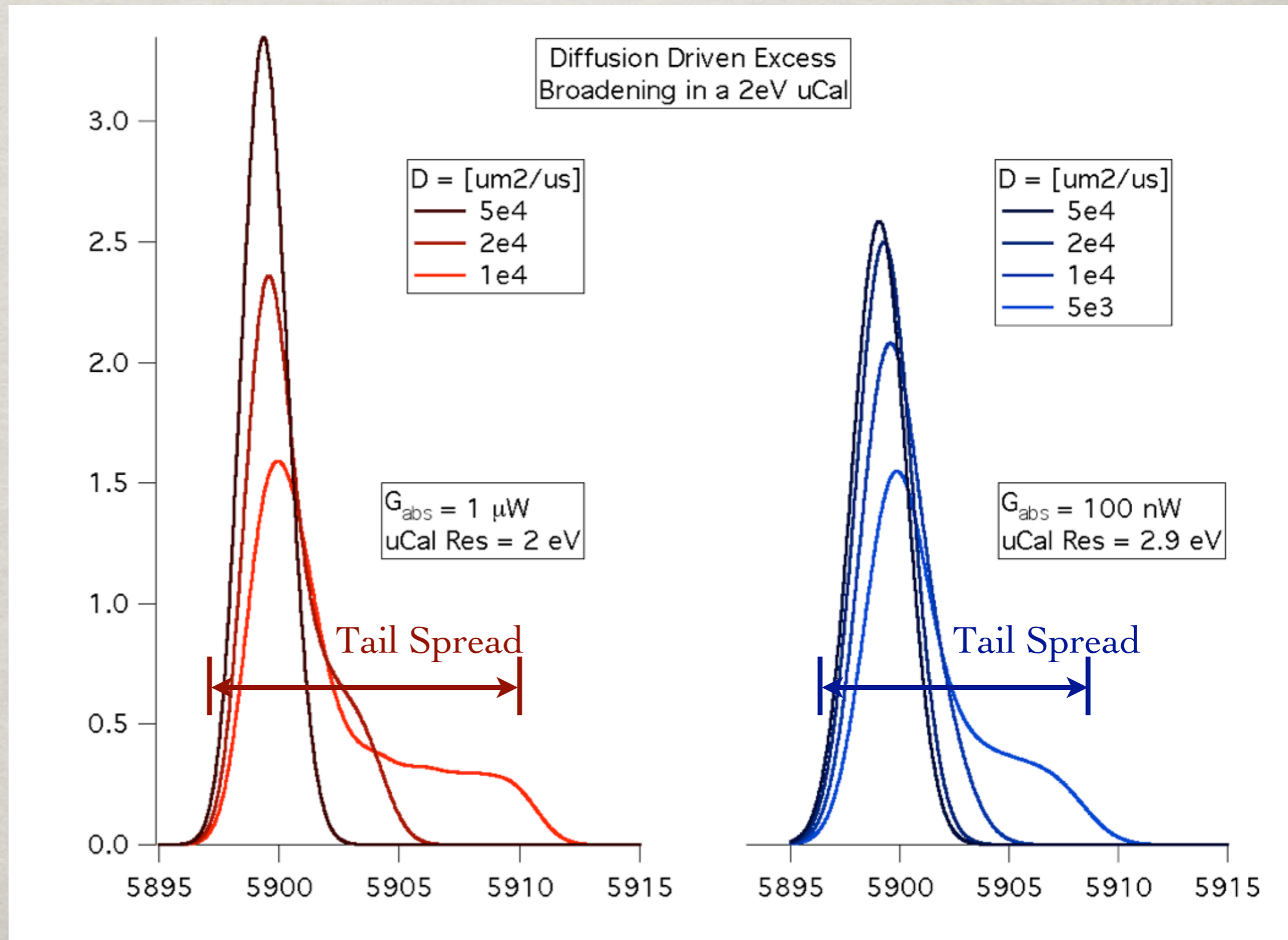
Pulse shape variation across surface of absorber shows significant “peakiness” near TES contact area

THE EFFECT OF VARYING PULSE SHAPE

To determine the effect of diffusion on detector performance we :

- ✻ Constructed an average pulse to be used as a template for an optimal filter
- ✻ Used the appropriate noise PSD for the device parameters
- ✻ Applied the optimal filter to all simulated pulses
- ✻ Convolved results with a 2 eV (FWHM) gaussian

δ -FN RESPONSE



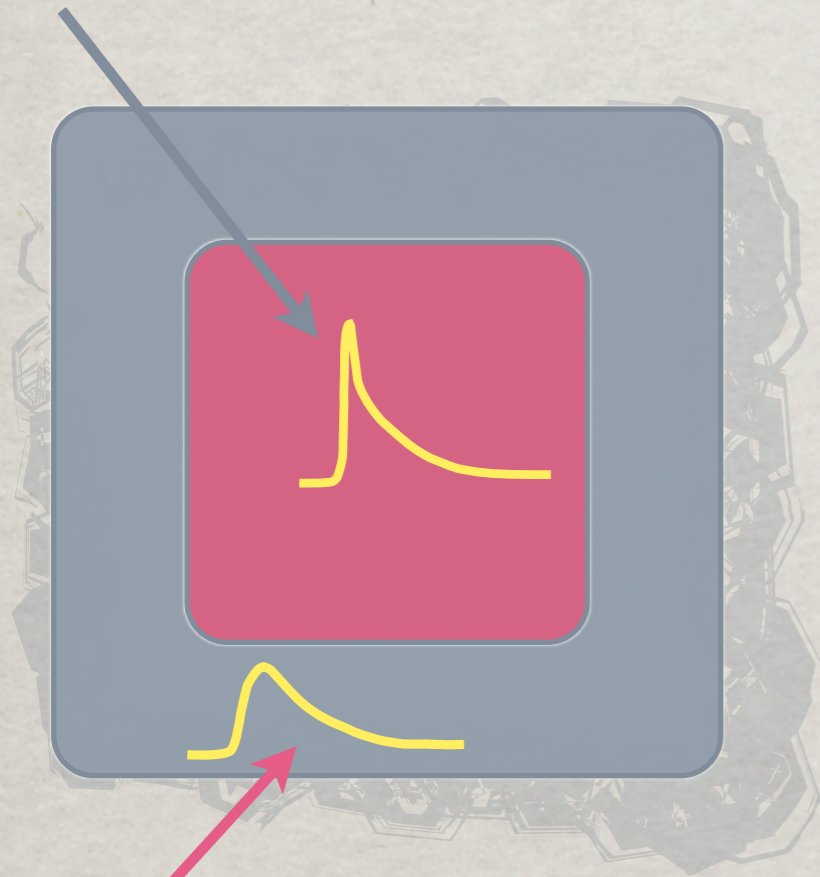
THE HIGH ENERGY TAIL

- ✻ For this mushroom geometry, the presence of a finite diffusion time in the absorber leads to a high energy tail in the δ -fn response
- ✻ No energy is being lost in the absorber. It all goes through the TES eventually
- ✻ The tail can be understood based on a geometric areas argument



THE HIGH ENERGY TAIL

Fast “peaky” pulses



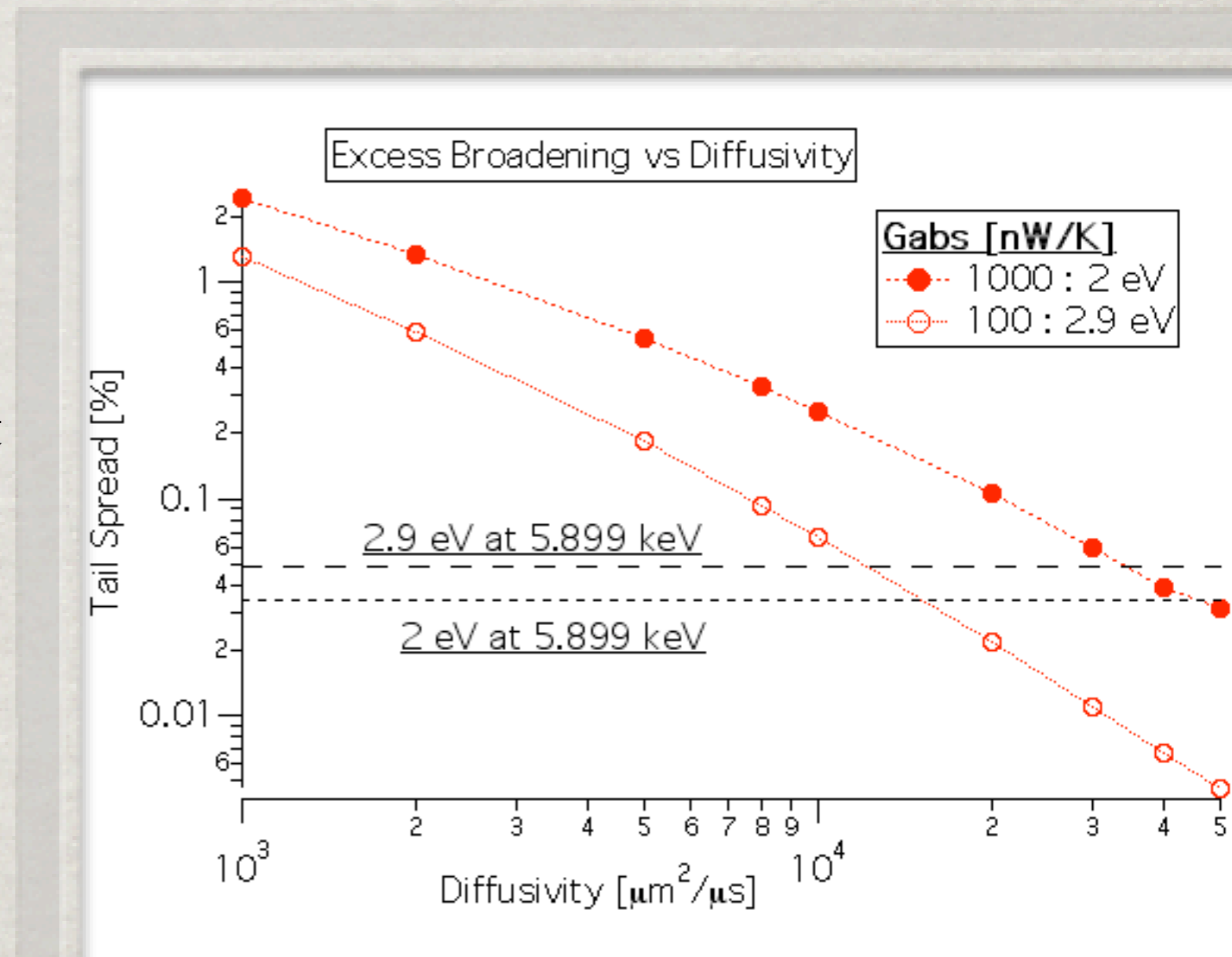
Slow pulses

- ✱ Area of mushroom overhang
~2x area of TES overlap.
- ✱ Template pulse is dominated
by slower overhang pulse
shapes
- ✱ Optimal filter operating on the
central “peaked” pulses results
in a larger reconstructed
energy than for the slower
pulses

QUANTIFYING THE EXCESS BROADENING

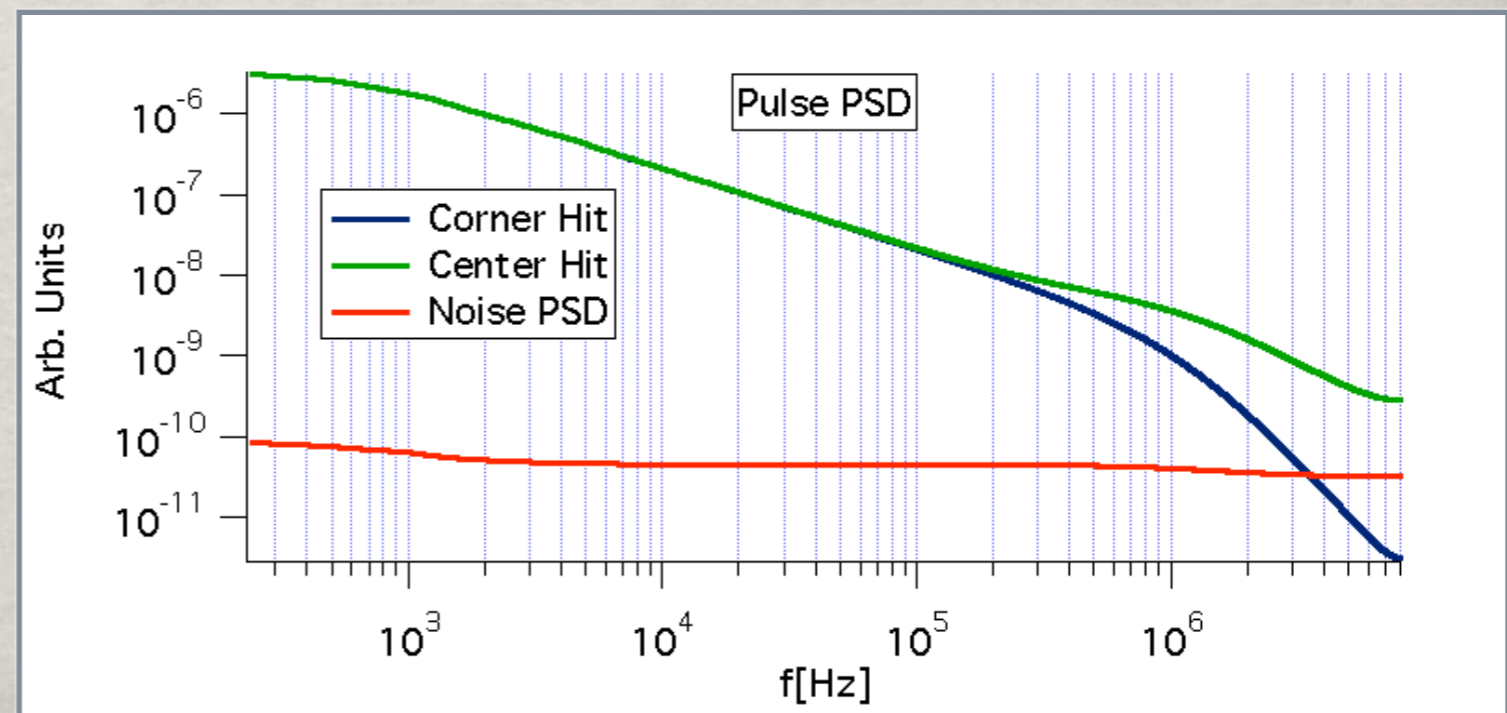
☼ Diffusion broadening leads to a non-gaussian high energy tail. We characterize it by the % spread with respect to the input energy

☼ Simulations spanning a factor of 10x in G_{abs} and 50x in D show that $D > 10^4 \mu\text{m}^2/\mu\text{s}$ is needed in order not to degrade the desired performance

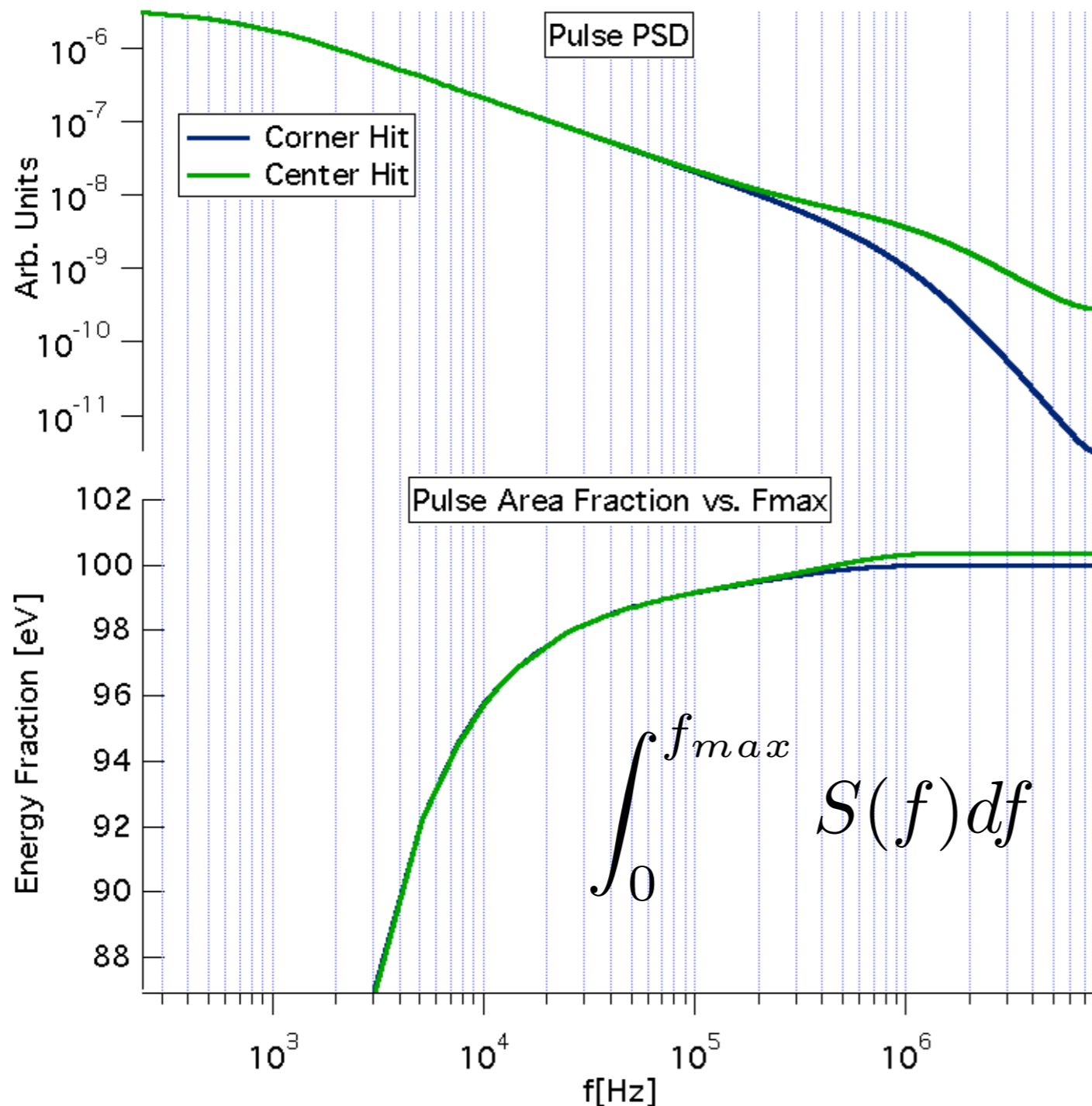


A FILTERING APPROACH TO MINIMIZING BROADENING

- ✿ The effects of position dependence in pulse shapes lies in the short time scales / high frequency bins
- ✿ Applying a low pass filter to the data stream does not help, however, since the frequency content of all signals are changed in the same way

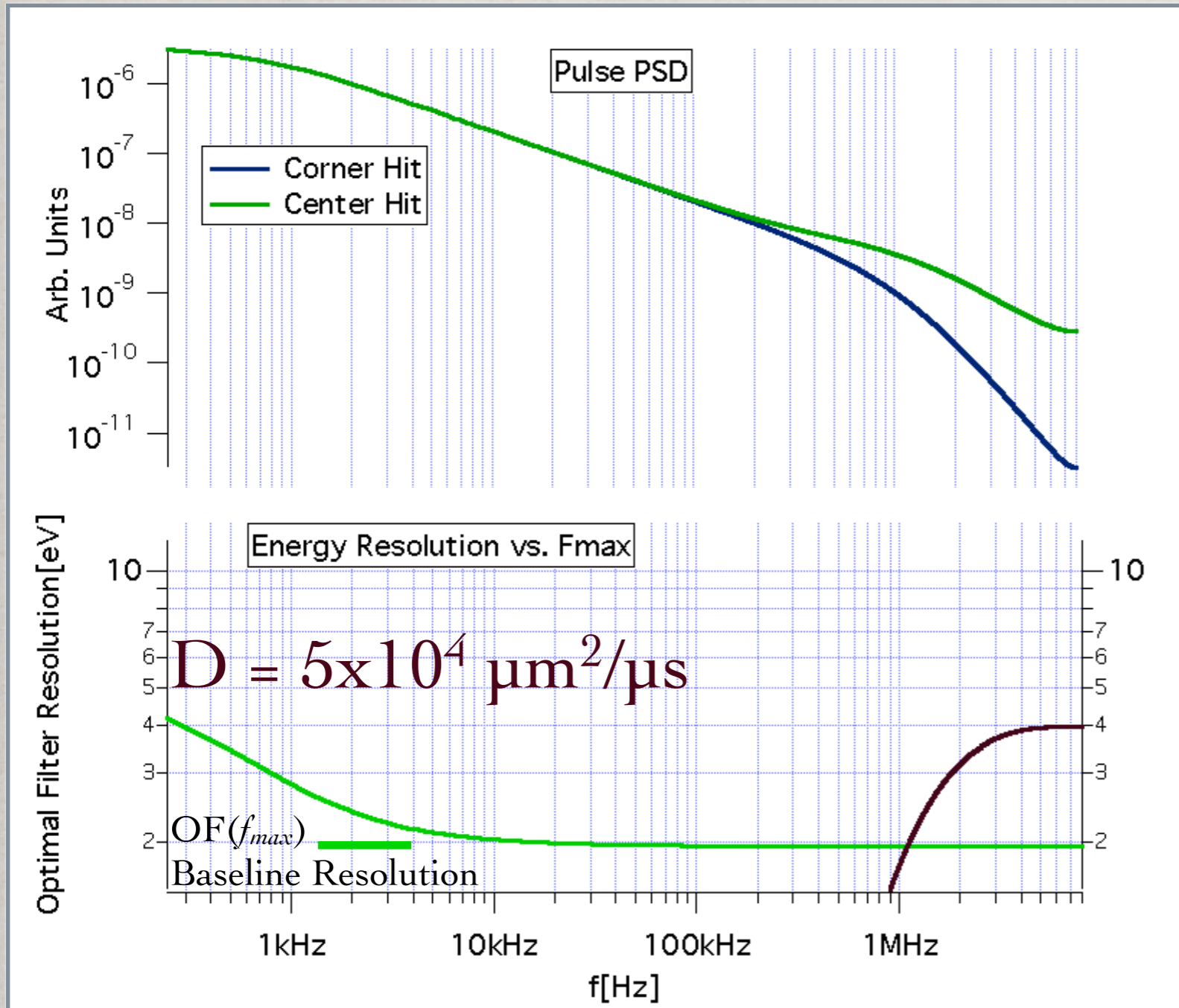


A MODIFIED OPTIMAL FILTER



- ✿ The majority of a pulses area is contained within the lower frequency bins
- ✿ We construct a modified optimal filter that only considers the signal up to a frequency f_{max} :
 $OF(f_{max})$

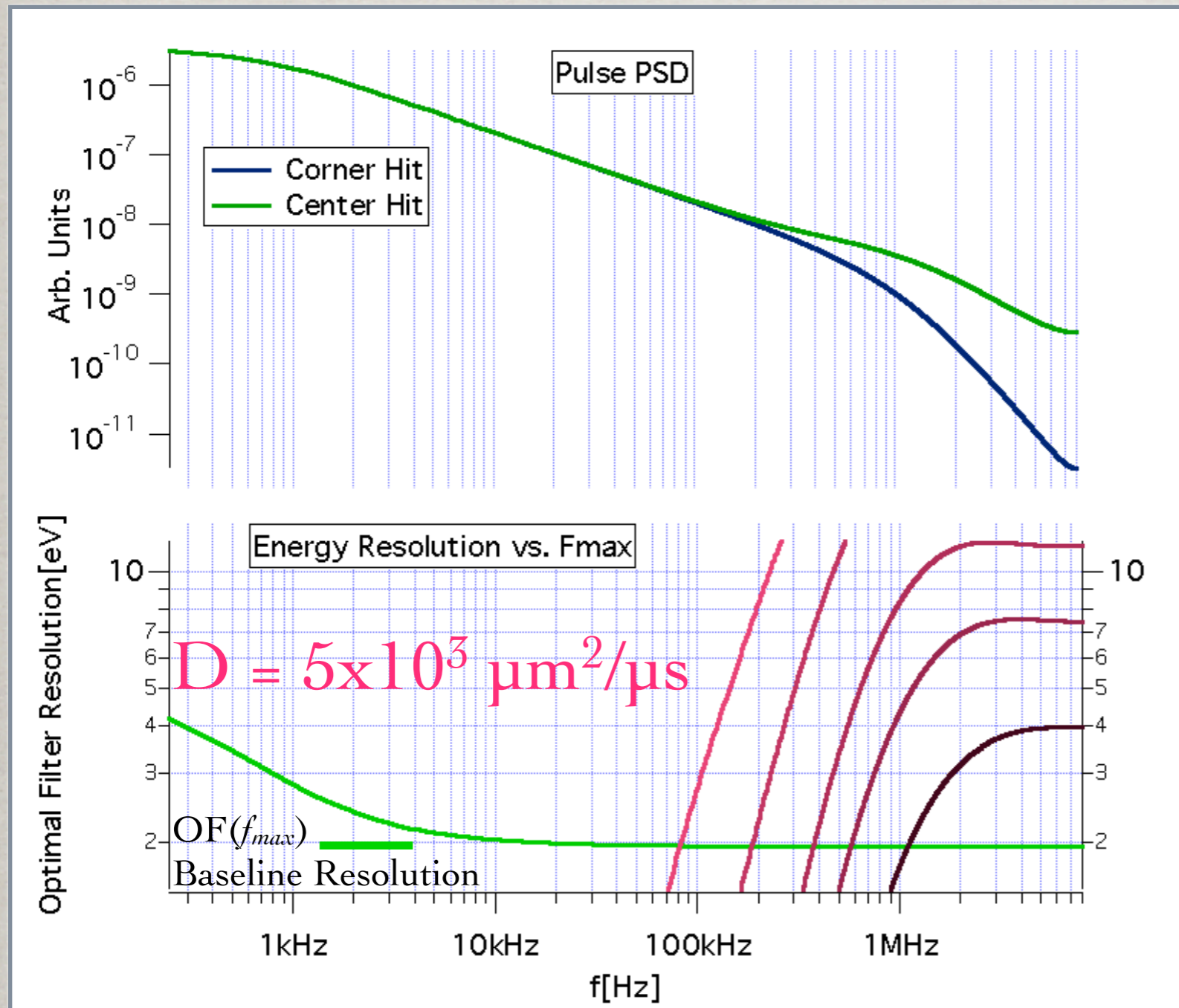
EFFECT OF $OF(f_{max})$ ON BASELINE RESOLUTION



✿ Plotting :

- ✿ Baseline resolution as a function of f_{max}
- ✿ Diffusion tail spread as a function of f_{max}

EFFECT OF $OF(f_{max})$ ON BASELINE RESOLUTION



☀ Plotting :

- ☀ Baseline resolution as a function of f_{max}
- ☀ Diffusion tail spread as a function of f_{max}
- ☀ For $f_{max} \sim 100$ kHz tail spread practically eliminated, AND the baseline resolution is unchanged

CONCLUSION

- ✻ Diffusivity value in the absorber of $\sim < 1 \times 10^4 \mu\text{m}^2/\mu\text{s}$ leads to significant pulse shape variation for mushroom shaped devices
- ✻ The use of a modified optimal filter that only considers frequencies below a certain f_{max} can eliminate excess broadening for a modest cost

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