## Report on Projects on CTCs and Black Holes

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Quantum Black Holes Meeting, OSU, Columbus, Ohio I 7th September 2004



#### Overview

Two projects:

What is the Stringy View of Closed Timelike Curves?

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Discover useful relation to recent proposal of S. Mathur?

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Pick a Theory of QG and learn what it has to say

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#### Taub-NUT in GR

$$ds^{2} = -f_{1}(dt - l\cos\theta d\phi)^{2} + f_{1}^{-1}dr^{2} + (r^{2} + l^{2})(d\theta^{2} + \sin^{2}\theta d\phi^{2})$$
$$f_{1} = 1 - 2\frac{Mr + l^{2}}{r^{2} + l^{2}}$$

*t* has period  $4\pi l$  (S<sup>3</sup> topology of constant *r* slices.)

$$f_1 < 0 \implies r \text{ is time}$$
  
 $f_1 > 0 \implies t \text{ is time}$ 

NUT: region with CTCs

Taub: Cosmological region bounded by Big Bang and Big Crunch

#### Taub-NUT in GR

Neighbourhood of Bang/Crunch is "Misner Spacetime":  $f_1(r) = 0$  is located at  $r = r_{\pm}$  For small  $\tau = r - r_{-}$ :

$$ds^2 = -(c\tau)^{-1}d\tau^2 + c\tau d\xi^2$$

Classic computation of stress tensor by Hiscock and Konkowski showed that it diverges there.

$$d\xi = dt - l\cos\theta d\phi$$
  
t fibred over S<sup>2</sup>

Suggests that (if one could do the QG computation), high energy effects destabilize the junction region.

Basis of e.g. Hawking's "Chronology Protection Conjecture"

Wish:

Embed this problem into string theory and retain computational control over the relevant questions:

Is there a mechanism by which string theory removes CTCs, or makes them harmless?

Is the geometry modified by high energy effects, as suggested in GR?

Could there be other mechanisms that we can discover to achieve this?

And/or are there some contexts in which CTCs have a physical role?

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Luckily, there is a non-trivial embedding of Taub-NUT into string theory... (C.V. Johnson, hep-th/9403192)



Furthermore... the full conformal field theory is known!

$$ds^{2} = k \left\{ d\sigma^{2} - \frac{\cosh^{2} \sigma - 1}{(\cosh \sigma + \delta)^{2}} (dt - \lambda \cos \theta d\phi)^{2} + d\theta^{2} + \sin^{2} \theta d\phi^{2} \right\}$$
$$\Phi - \Phi_{0} = -\frac{1}{2} \ln(\cosh \sigma + \delta) \qquad B_{t\phi}, A_{\phi}, A_{t}, \text{ also non-zero.}$$

Leading order in large  $k \sim 1/\alpha'$ Heterotic string solution.

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 $\boldsymbol{L}$ 



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$$\Phi - \Phi_0 = -\frac{1}{2}\ln(\cosh\sigma + \delta)$$

 $B_{t\phi}, A_{\phi}, A_t$ , also non-zero.

This is the near-horizon geometry of a larger solution



$$ds^{2} = k \left( \frac{dx^{2}}{x^{2} - 1} - \frac{x^{2} - 1}{(x + \delta)^{2}} (dt - \lambda \cos \theta d\phi)^{2} + d\theta^{2} + \sin^{2} \theta d\phi^{2} \right) ,$$
  
$$\Phi - \Phi_{0} = -\frac{1}{2} \ln(x + \delta) .$$

 $0 \le x \le +\infty \qquad x = \cosh \sigma$  $-\infty \le x \le 0 \qquad -\infty \le x \le 0$  $-1 \le x \le +1 \qquad x = -\cos \tau$ 

Yes, can extend to the full three regions. The CFT naturally does this!

C.V. Johnson, and H.G. Svendsen, hep-th/0405141



#### CFT is "Heterotic coset model":

 $\frac{SL(2,\mathbb{R})\times SU(2)}{U(1)\times U(1)}$ 

o Asymmetric (anomalous) gauging of WZNW.
o Anomalies cured by heterotic fermions.
o Gives larger class of geometries than usual.
o Series of techniques for extracting geometry.

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## High Energy Corrections

Possible to get all corrections!

o Adapt a technique of Tseytlin, and Bars and Sfetsos. o Uses results for exact effective action of WZNW o ....and gauged WZNW

o This allows for extraction of full geometry.

A.Tseytlin, hep-th/9301015, I. Bars and K. Sfetsos, hep-th/9301047.

$$ds^{2} = (k-2) \left( \frac{dx^{2}}{x^{2}-1} + F(x)(dt - \lambda\cos\theta d\phi)^{2} + d\theta^{2} + \sin^{2}\theta d\phi^{2} \right) ,$$
$$\hat{\Phi} - \hat{\Phi}_{0} = -\frac{1}{4} \ln D(x)$$
$$F(x) = -\frac{x^{2}-1}{D(x)} = -\left(\frac{(x+\delta)^{2}}{x^{2}-1} - \frac{4}{k+2}\right)^{-1}$$

C.V. Johnson, and H.G. Svendsen, hep-th/0405141



Significant corrections, but the transition neighbourhoods,  $x = \pm 1$ , are unchanged!

The transition regions allow the CTCs to persist, contrary to expectations of chronology protection, etc!

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The transition regions allow the CTCs to persist, contrary to expectations of chronology protection, etc!

Maybe a sign that strings say that CTCs are ok? Tantalizing possibility. Especially in view of the adjoining cosmological region.

CTCs in "quantum geometric" description of pre-BB? Do not confuse with flat space result! Curvature is non-zero.

#### A Special Geometry



#### To-do List

Not done yet.

Examine quantum effects? String Loop corrections?



#### To-do List



Consider the DI-D5-P black hole, made using K3



 $x^{0}, x^{1}, x^{2}, x^{3}, x^{4}, x^{5}, x^{6}, x^{7}, x^{8}, x^{9}$  $N_1 D1$  $N_5 D5 -$ 

 $Q_P$  units of momentum in  $\chi^5$ 

Supergravity solution:

 $ds^{2} = f_{1}^{-3/4} f_{5}^{-1/4} \left( -dt^{2} + dz^{2} + k \left( dt - dz \right)^{2} \right) + f_{1}^{1/4} f_{5}^{-1/4} ds^{2} + f_{1}^{1/4} f_{5}^{3/4} \left( dr^{2} + r^{2} d\Omega_{3}^{2} \right)$ 

 $x^0 - x^4 \Rightarrow t, r, \theta, \phi, \chi$ .  $x^5 \equiv z$   $d\Omega_3^2 = d\theta^2 + \sin^2\theta (d\phi^2 + \sin^2\phi d\chi^2)$ 

Supergravity solution:

$$\begin{split} ds^2 &= f_1^{-3/4} f_5^{-1/4} \left( -dt^2 + dz^2 + k \left( dt - dz \right)^2 \right) + f_1^{1/4} f_5^{-1/4} ds^2 + f_1^{1/4} f_5^{3/4} \left( dr^2 + r^2 d\Omega_3^2 \right) \\ x^0 - x^4 &\Rightarrow t, r, \theta, \phi, \chi \, . \qquad x^5 \equiv z \qquad d\Omega_3^2 = d\theta^2 + \sin^2 \theta (d\phi^2 + \sin^2 \phi d\chi^2) \\ e^{2\Phi} &= \frac{f_1}{f_5} \, , \qquad F_{rtz}^{(3)} = \partial_r f_1^{-1} \, , \qquad F_{\theta\phi\chi}^{(3)} = 2 \, r_5^2 \sin^2 \theta \sin \phi \, . \\ f_1 &= 1 + \frac{r_1^2}{r^2} \, , \qquad f_5 = 1 + \frac{r_5^2}{r^2} \, , \qquad k = \frac{r_P^2}{r^2} \end{split}$$

Supergravity solution:

$$ds^{2} = f_{1}^{-3/4} f_{5}^{-1/4} \left( -dt^{2} + dz^{2} + k \left( dt - dz \right)^{2} \right) + f_{1}^{1/4} f_{5}^{-1/4} ds^{2} + f_{1}^{1/4} f_{5}^{3/4} \left( dr^{2} + r^{2} d\Omega_{3}^{2} \right)$$

$$x^{0} - x^{4} \Rightarrow t, r, \theta, \phi, \chi, \qquad x^{5} \equiv z \qquad d\Omega_{3}^{2} = d\theta^{2} + \sin^{2} \theta \left( d\phi^{2} + \sin^{2} \phi d\chi^{2} \right)$$

$$e^{2\Phi} = \frac{f_{1}}{f_{5}}, \qquad F_{rtz}^{(3)} = \partial_{r} f_{1}^{-1}, \qquad F_{\theta\phi\chi}^{(3)} = 2r_{5}^{2} \sin^{2} \theta \sin \phi.$$

$$f_{1} = 1 + \frac{r_{1}^{2}}{r^{2}}, \qquad f_{5} = 1 + \frac{r_{5}^{2}}{r^{2}}, \qquad k = \frac{r_{P}^{2}}{r^{2}} \qquad V^{*} \equiv (2\pi\ell_{s})^{4}$$

$$r_{5}^{2} = g_{s}\ell_{s}^{2}Q_{5}, \qquad r_{1}^{2} = g_{s}\ell_{s}^{2}\frac{V^{*}}{V}Q_{1}, \qquad r_{P}^{2} = g_{s}^{2}\ell_{s}^{2}\frac{V^{*}}{V}\frac{\ell_{s}^{2}}{R_{z}^{2}}Q_{P} \qquad r_{H}^{2} = 0$$

#### Being Careful with Charge

The K3 induces negative D1 charge:

 $Q_5 = N_5$  The number of D5-branes  $Q_1 = N_1 - N_5$  Correction due to induced charge.

#### Area and Entropy

# $S = \frac{A}{4G} = 2\pi \sqrt{Q_1 Q_5 Q_P} = 2\pi \sqrt{(N_1 - N_5)N_5 Q_P}$

There is something disturbing about this formula.

#### Area and Entropy



You can reduce the entropy by adding more D-branes!!

#### Violation of Second Law?

Add brane with charges  $(n_1, n_5)$ :

$$\delta S^2 = 4\pi^2 Q_P (N_5 n_1 + (N_1 - 2N_5)n_5)$$



#### Saving the Second Law

Introduce a probe brane with charges  $(n_1, n_5)$ .

Find that it is only physical for motion satisfying:

$$r^{2} > r_{e}^{2} \equiv g_{s}\ell_{s}^{2}V^{*}\frac{(2N_{5}-N_{1})n_{5}-N_{5}n_{1}}{(V-V^{*})n_{5}+V^{*}n_{1}}$$

This is the Enhançon Radius.

C.V. Johnson and R. C. Myers, hep-th/0105159

#### A Microscopic Filtering Mechanism

The bound state branes cannot approach the horizon *precisely* when they would decrease its area!

 $r_{H}^{2} = 0$ 

$$r^{2} > r_{e}^{2} \equiv g_{s}\ell_{s}^{2}V^{*}\frac{(2N_{5}-N_{1})n_{5}-N_{5}n_{1}}{(V-V^{*})n_{5}+V^{*}n_{1}}$$
$$\delta S^{2} = 4\pi^{2}Q_{P}(N_{5}n_{1}+(N_{1}-2N_{5})n_{5})$$

They *can* approach when there is no danger.

#### Where to go from here...?

D-branes are the smallest charge carriers

The mechanism works as a *microscopic* filter to preserve a *macroscopic* law.

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Deserves to be better understood.

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Clues in the context of Mathur's proposal? Can the new proposal supply a natural description for D5 emission?