

11/20/15

(Work)  $\vec{F}_B = q \vec{v} \times \vec{B}$

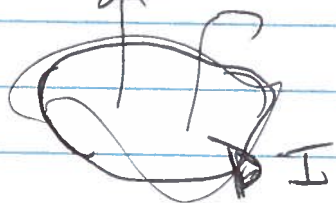
$\vec{w} \cdot \vec{F} = 0$

"magnetic fields do no work"

changing  $\vec{B}$

needs ~~less~~ work!

current flow against  $\mathcal{E}$



$$\mathcal{E} = -\frac{d\Phi}{dt}$$

$$\frac{dW}{dt} = -I\mathcal{E} \quad \left(\frac{dW}{dt}\right) \left(\frac{d\Phi}{dt}\right)$$

$$dW = +I \delta\Phi = I \cdot \int_S d^2a \hat{n} \cdot \delta\vec{B}$$

$$= I \int_S d^2a \hat{n} \cdot (\vec{\nabla} \times \vec{A}) = \oint I d\vec{l} \cdot \vec{A}$$

"many loops"  $\rightarrow$   $\boxed{dW = \int d^3x \vec{J} \cdot \vec{A}} \quad (5.144)$

[cf.  $dW = \int d^3x \vec{\phi} \cdot \vec{\phi} \quad (4.84)$ ]

$\vec{J} = \vec{\nabla} \times \vec{H}$

$$\vec{\nabla} \cdot (\vec{H} \times \vec{A}) = \epsilon_{ijk} \nabla_i H_j A_k$$

$$= \epsilon_{ijk} ((\nabla_i H_j) A_k + H_j (\nabla_i A_k))$$

$$= \underbrace{\vec{A} \cdot (\vec{\nabla} \times \vec{H})}_{\vec{J}} - \underbrace{\vec{H} \cdot (\vec{\nabla} \times \vec{A})}_{\vec{B}}$$



$$\int_V d^3x \vec{J} \cdot \vec{\delta A} = \int d^3x \left( \vec{\nabla} \cdot (\vec{H} \times \vec{\delta A}) + \vec{H} \cdot \vec{\delta B} \right) \quad \text{Surface} \rightarrow 0 \quad (2)$$

$$\delta W = \int d^3x \vec{H} \cdot \vec{\delta B} \quad (5.147)$$

(cf.  $\delta W = \int d^3x \vec{E} \cdot \vec{\delta D}$ .)

Linear material  $\vec{B} = \mu \vec{H}$

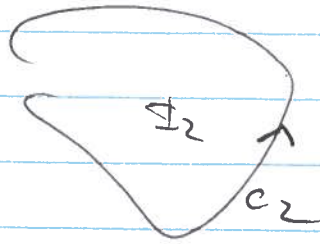
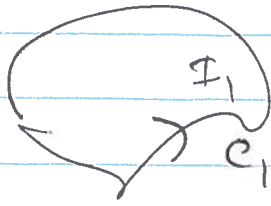
$$\delta(\vec{B} \cdot \vec{H}) = \delta\left(\frac{B^2}{\mu}\right) = \frac{2}{\mu} \vec{B} \cdot \vec{\delta B} = 2\vec{H} \cdot \vec{\delta B}$$

$$W = \int d^3x \frac{1}{2} \vec{H} \cdot \vec{B} = \frac{1}{2\mu} \int d^3x \mu H^2 = \int d^3x \frac{1}{2\mu} B^2$$

Hard magnetization  $\vec{m}_0$

$$W = \int d^3x \frac{1}{2} \mu_0 H^2 - \int d^3x \frac{1}{2} \mu_0 m_0^2$$

independent of location, orientation



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$$W = \int d^3x \frac{1}{2} \vec{J}(\vec{x}) \cdot \vec{A}(\vec{x})$$

$$= \int d^3x \frac{1}{2} \vec{J}(\vec{x}) \cdot \frac{\mu_0}{4\pi} \int d^3x' \frac{\vec{J}(\vec{x}')}{|\vec{x} - \vec{x}'|}$$

$$= \sum_{i=1}^N \int_{C_i} d^3x_i \vec{J}_i(\vec{x}_i) \cdot \frac{\mu_0}{4\pi} \int_{C_j} d^3x_j \frac{\vec{J}_j(\vec{x}_j)}{|\vec{x}_i - \vec{x}_j|}$$

$$= \frac{1}{2} \sum_{i=1}^N \sum_{j=1}^N I_i I_j \frac{\mu_0}{4\pi} \int_{C_i} d^3x_i \int_{C_j} d^3x_j \frac{\vec{J}_i(\vec{x}_i) \cdot \vec{J}_j(\vec{x}_j)}{|\vec{x}_i - \vec{x}_j|}$$

$$W = \frac{1}{2} \sum_i L_i I_i^2 + \sum_{i < j} \frac{1}{2} M_{ij} I_i I_j$$

self.  $L_i = \frac{\mu_0}{4\pi} \frac{1}{|I_i|^2} \int_{C_i} d^3x_i \int_{C_i} d^3x_i \frac{\vec{J}_i(\vec{x}_i) \cdot \vec{J}_i(\vec{x}_i)}{|\vec{x}_i - \vec{x}_i|}$

$M_{ij}$   $\left( \begin{matrix} i, j \\ , \end{matrix} \right)$

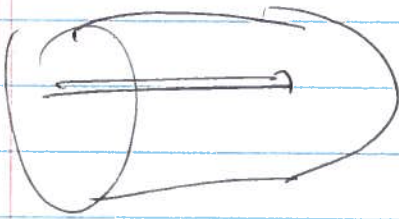
(4)

Thin wires

$$W_{ij} = \frac{1}{4\pi\epsilon_0} \int_{c_i} \vec{\Delta}_i d\vec{l}_i \cdot \frac{\mu_0}{4\pi} \int_{c_j} \frac{\vec{\Delta}_j d\vec{l}_j}{|\vec{r}_{ij}|}$$

$$= \frac{1}{4\pi\epsilon_0} \int_{c_i} d\vec{l}_i \cdot \vec{A}_j(\vec{r}_i)$$

$$= \frac{1}{4\pi\epsilon_0} \int_{S_i} d\vec{a} \cdot \hat{n} \cdot (\vec{\nabla} \times \vec{A}_j) = \frac{1}{4\pi\epsilon_0} \int_{S_i} \leftarrow \begin{matrix} \text{flux} \\ \text{through } i \end{matrix}$$



$$\text{Cocx. } \vec{B} = \frac{\mu_0 I}{2\pi r} \hat{\phi}$$

$$W = \int 2\pi r dr \cdot l \cdot \frac{1}{2\mu_0} \left( \frac{\mu_0 I}{2\pi r} \right)^2$$

$$= \frac{1}{2} \frac{\mu_0}{2\pi} I^2 \cdot l \cdot \int_a^b \frac{r dr}{r^2} = \frac{1}{2} \frac{\mu_0 I^2}{2\pi} \ln \frac{b}{a} \cdot l$$

$$= \frac{1}{2} L I^2 \quad \left| \frac{L}{l} = \frac{\mu_0}{2\pi} \ln \frac{b}{a} \right.$$

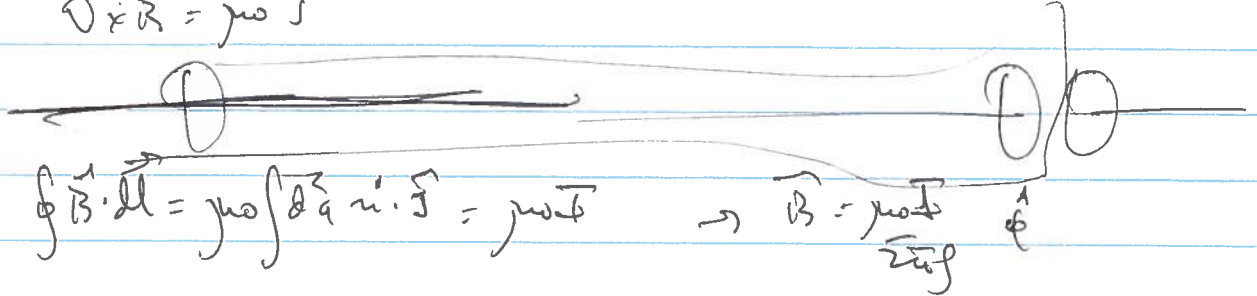
$$W_e = \int 2\pi r dr \cdot l \cdot \frac{1}{2} \epsilon_0 \left( \frac{\lambda}{2\pi\epsilon_0 r} \right)^2 = \frac{1}{2} \frac{1}{2\pi\epsilon_0} \frac{(\lambda l)^2}{l} \ln \frac{b}{a}$$

$$\stackrel{=Q^2}{2\epsilon} \left| \frac{C}{l} = \frac{2\pi\epsilon_0}{\ln \frac{b}{a}} \right. \left. \left| \frac{C}{l} \cdot \frac{l}{l} = \mu_0 \epsilon_0 = \frac{1}{c^2} \right. \right)$$

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Need another piece to Maxwell equations

$$\vec{\nabla} \times \vec{B} = \mu_0 \vec{J}$$



$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \int d\vec{a} \cdot \vec{J} = \mu_0 I$$

$$\vec{B} = \frac{\mu_0 I}{2\pi r} \hat{\phi}$$

$$\vec{E} = \frac{\sigma \hat{n}}{\epsilon_0}$$

$$\vec{J}_e = \pi a^2 \frac{\sigma}{\epsilon_0} = \frac{Q}{\epsilon_0}$$

$$\frac{dJ_e}{dt} = \frac{I}{\epsilon_0} \quad \vec{J}_E = \epsilon_0 \frac{dJ}{dt}$$

$$\vec{J}_E = \epsilon_0 \frac{\partial \vec{E}}{\partial t}$$

$$\vec{\nabla} \times \vec{B} = \mu_0 \left( \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right)$$

↑ units.

Displacement current

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{B}) = 0 = \mu_0 \left( \vec{\nabla} \cdot \vec{J} + \vec{\nabla} \cdot \left( \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right) \right)$$

$$= \mu_0 \left( \vec{\nabla} \cdot \vec{J} + \epsilon_0 \frac{\partial (\vec{\nabla} \cdot \vec{E})}{\partial t} \right)$$

$$\frac{\partial \rho}{\partial t} + \vec{\nabla} \cdot \vec{J} = 0$$

automatic

EM/waves

relativity

Gauge invariance  
→ renormalizability