

4/16/2018 Jackson (11.149)

2

$$\vec{E}' = \gamma(\vec{E} + \vec{\beta} \times \vec{B}) - \frac{\gamma^2}{\gamma H} \vec{\beta} (\vec{\beta} \cdot \vec{E})$$

$$\vec{B}' = \gamma(\vec{B} - \vec{\beta} \times \vec{E}) - \frac{\gamma^2}{\gamma H} \vec{\beta} (\vec{\beta} \cdot \vec{B})$$

$\vec{E} \parallel \vec{\beta}$ $\vec{\beta} \times \vec{E} = 0$ $\vec{E}' = \gamma \vec{E} - \frac{\gamma^2 \beta^2}{\gamma H} \vec{E}$

$$\gamma - \frac{\gamma^2 \beta^2}{\gamma H} = \gamma(\gamma H) - \gamma^2 \beta^2 = \frac{\gamma^2 (1 - \beta^2)}{\gamma H} = 1$$

$$E'_\parallel = E_\parallel$$

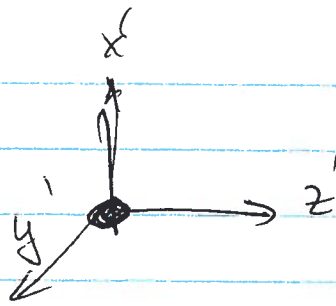
$$B'_\parallel = B_\parallel$$

$$\vec{E}'_\perp = \gamma(\vec{E}_\perp + \frac{v}{c} \times \vec{B}_\perp)$$

$$\vec{B}'_\perp = \gamma(\vec{B}_\perp - \frac{v}{c} \times \vec{E}_\perp)$$

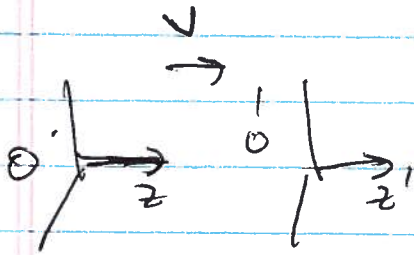
②

Point q



@ rest.

$$\vec{E} = \frac{q \vec{r}'}{r'^2} = \frac{q \vec{z}'}{(x'^2 + y'^2 + z'^2)^{3/2}}$$



$$\begin{aligned} x &= x' \\ y &= y' \\ z &= \gamma(z' + vt') \\ t &= \gamma(t' + \frac{vz'}{c^2}) \end{aligned}$$

inverse.

$$\begin{aligned} z' &= \gamma(z - vt) \\ t' &= \gamma(t - \frac{vz}{c^2}) \end{aligned}$$

origins coincide @ $t=t'=0$

$$E^z = E_{||} = E'_{||} = E'^z = \frac{q z'}{(x'^2 + y'^2 + z'^2)^{3/2}} = \frac{q \cdot \gamma(z - vt)}{[x^2 + y^2 + \gamma^2(z - vt)^2]^{3/2}}$$

$$\vec{E}_{\perp} = \gamma \left(\vec{E}'_{\perp} + \frac{v}{c^2} \times \vec{B}' \right) = \gamma \vec{E}'_{\perp} = \gamma \cdot \frac{q \vec{x}'_{\perp}}{(x'^2 + y'^2 + z'^2)^{3/2}}$$

$$E^x = \frac{\gamma q x}{(x^2 + y^2 + \gamma^2(z - vt)^2)^{3/2}}$$

$$E^y = \dots$$

(at ∞)

$$\vec{E} = \frac{\gamma q (x \hat{x} + y \hat{y} + z \hat{z})}{(x^2 + y^2 + \gamma^2 z^2)^{3/2}}$$

($z = r \cos \theta$)

$$\vec{E} = \frac{\gamma q \vec{x}}{r^3 (\sin^2 \theta \cos^2 \phi + \sin^2 \theta \sin^2 \phi + \gamma^2 \cos^2 \theta)^{3/2}}$$

$$\vec{E} = \frac{q \vec{r}}{r^2} \frac{1}{(\sin^2 \theta + \gamma^2 \cos^2 \theta)^{3/2}}$$

$$\vec{E} = q \frac{\vec{r}}{r^2} \frac{1}{r^2} \frac{1}{(1 - \frac{v^2}{c^2} \sin^2 \theta)^{3/2}}$$

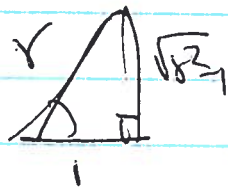
$$\vec{E} = q \frac{\vec{r}}{r^2} \cdot \frac{1}{r^2} \frac{1}{(1 - \frac{v^2}{c^2} \sin^2 \theta)^{3/2}} \quad (11.154)$$

radial inverse square squeezed in direction of motion

$$\oint_{\text{sphere}} d\vec{a} \cdot \vec{E} \cdot \hat{n} = \int r^2 d\Omega \cdot q \frac{1}{r^2} \frac{1}{(1 - \frac{v^2}{c^2} \sin^2 \theta)^{3/2}}$$

$$= q \int d\Omega \frac{1}{(1 + \gamma^2 - \gamma^2 \cos^2 \theta)^{3/2}} = 2\pi q \int_{-1}^1 dp \frac{1}{(1 + (\gamma^2 - 1)p^2)^{3/2}}$$

$$(\gamma^2 - 1)p^2 = \tan^2 \phi \rightarrow 2\pi q \int_{-\phi_0}^{\phi_0} \frac{\sec^2 \phi d\phi}{\sqrt{\gamma^2 - 1}} \cdot \frac{1}{(1 + \tan^2 \phi)^{3/2}}$$



$$= \frac{2\pi q}{\sqrt{\gamma^2 - 1}} \int_{-\phi_0}^{\phi_0} \cos \phi d\phi$$

$$= 2\pi q \cdot 2 \sin \phi_0$$

$$= 2\pi q \cdot 2 \left(\frac{\sqrt{\gamma^2 - 1}}{\gamma} \right) \rightarrow \underline{4\pi q}$$

(4)

Jumping ahead. §12.11 covariant wave equation point charge source.

$$\square^2 A^\mu = \frac{4\pi}{c} J^\mu$$

(1.133)
(12.124)

point particle. $\rho = e \delta^{(3)}(\vec{x} - \vec{r}(t))$
 $\vec{J} = e \vec{v} \cdot \delta^{(3)}(\vec{x} - \vec{r}(t))$

$$J^\mu(x) = ec \int d\tau U^\mu(\tau) \cdot \delta^{(4)}(x^\alpha - r^\alpha(\tau)) \quad (12.139)$$

(14.2)

$$\delta(x^0 - r^0(\tau)) \delta^{(3)}(\vec{x} - \vec{r}(\tau))$$

use $\delta(x^0 - r^0(\tau))$ to do τ -integral

$$\delta(x^0 - r^0(\tau)) = \frac{\delta(\tau - \tau_0)}{\left| \frac{dr^0}{d\tau} \right|_0} = \frac{\delta(\tau - \tau_0)}{\left| \frac{d}{d\tau}(c\tau) \right|}$$

$$= \frac{1}{rc} \cdot \delta(\tau - \tau_0)$$

$$J^\mu(x) = (ec) \left(\frac{1}{rc} \right) \left(\frac{dc}{d\tau} \right) \delta^{(3)}(\vec{x} - \vec{r}(t_0))$$

$$J^0 = ec \cdot \delta^{(3)}(\vec{x} - \vec{r}(t))$$

$$\vec{J} = e \vec{v} \cdot \delta^{(3)}(\vec{x} - \vec{r}(t))$$



③

$$\square^2 A^\mu = \frac{4\pi}{c} J^\mu$$

↓ c

$$A^\mu(\vec{x}, t) = A^\mu(x) = \int d^3x' dt' \underset{c}{\uparrow} \frac{4\pi}{c} J^\mu(x') \frac{1}{4\pi} \frac{1}{|\vec{x} - \vec{x}'|} \delta(t' - t - \frac{1}{c}R)$$

$$\delta[(x-x')^2] = \delta[(x^0-x'^0)^2 - |\vec{x}-\vec{x}'|^2]$$

$$= \delta[(x^0-x'^0+R)(x^0-x'^0-R)]$$

$$= \frac{1}{2R} \delta(x^0-x'^0+R) + \frac{1}{2R} \delta(x^0-x'^0-R)$$

$$\frac{x'^0 = x^0 + R}{\text{advanced}}$$

$$\frac{x'^0 = x^0 - R}{\text{retarded}}$$

$$F_{\text{ret}} = D_r = \frac{1}{2\pi} \delta[(x-x')^2] \Theta(x^0-x'^0) \quad \underline{ct' > ct}$$

$$F_{\text{adv.}} = D_a = \frac{1}{2\pi} \delta[(x-x')^2] \Theta(x'^0-x^0) \quad \underline{ct' > ct}$$

$$A^\mu(x) = \int d^4x' D_r(x-x') \frac{4\pi}{c} J^\mu(x')$$

(12.134)

(14.1)

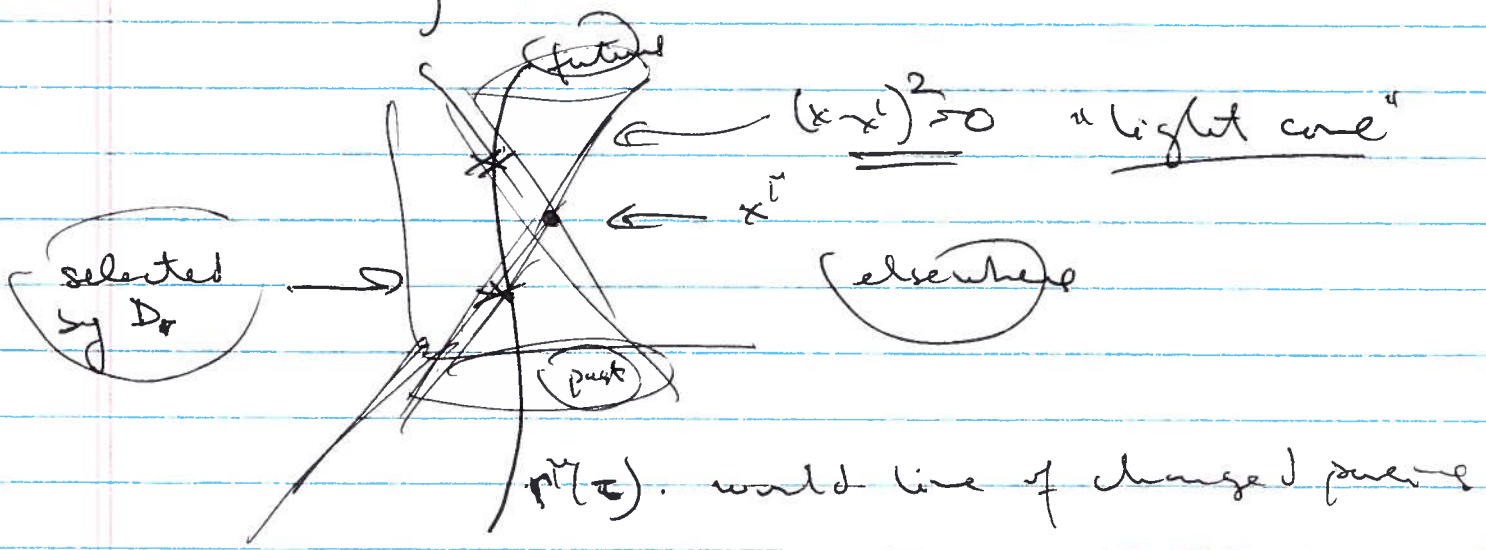
Chapter 14. Radiation by moving charges

$$A^M(x) = \int dx' \frac{1}{2a} \Theta(x^0 - x'^0) \delta[(x-x')^2]$$

≡

$$\times \text{e.c.} \frac{4\pi}{c} \int dt V^M(t) \delta^{(4)}(x - r(t))$$

$$A^M = 2e \int dt V^M(t) \delta[(x - r(t))^2] \cdot \Theta(ct - r^0(t))$$



~~$$\delta[(x - r(t))^2]$$~~

$$\frac{d}{dt} [(x - r(t))^2] = \frac{d}{dt} \left[(x^0 - r^0(t))^2 - (\vec{x} - \vec{r}(t))^2 \right]$$

$$= 2(x^0 - r^0) \left(-\frac{dr^0}{dt} \right) - 2(\vec{x} - \vec{r}(t)) \cdot \left(-\frac{d\vec{r}}{dt} \right)$$

$$= -2(x - r) \cdot \frac{dr}{dt} = -2V \cdot (x - r)$$

$$A^M(x) = \frac{e V^M(t)}{V \cdot (x - r)} \quad (14.6)$$

$$\vec{r} = \left(c(t-t') \right)$$

$$\vec{V} = \left(\frac{dc}{dt} \right) = dc \cdot \left(\frac{1}{\beta} \right)$$

$$\left. \begin{aligned} \Phi(\vec{x}, t) &= \frac{q}{R(1-\hat{u} \cdot \vec{\beta})} \\ t' &= t - \frac{1}{c}R \end{aligned} \right\} \text{ret}$$

$$\left. \begin{aligned} \vec{A}(\vec{x}, t) &= \frac{q \vec{\beta}}{R(1-\hat{u} \cdot \vec{\beta})} \end{aligned} \right\} \text{ret}$$

Famous. Liénard - Wiechert potentials

1898

1901.

Engineer.

Geophysicist.

École des mines

1892-1907

Göttingen

de St Étienne

Chaire d'électricité industrielle

École des mines de Paris 1908-1929.