

11/2/2016

Lie derivative \mathbb{L} directional

$$\mathbb{L}_u f = \nabla_u f = u^\alpha \nabla_\alpha f. \quad (\text{scalar})$$

$$\mathbb{L}_u v = [u, v] = \nabla_u v - \nabla_v u. \quad (\text{vector})$$

$$= u^\alpha v^\mu_{;\alpha} - v^\alpha u^\mu_{;\alpha} = u^\alpha v^\mu_{,\alpha} - v^\alpha u^\mu_{,\alpha}$$

$$\mathbb{L}_u (fv) = \nabla_u (fv) - \nabla_{(fv)} u$$

$$= (\nabla_u f) \cdot v + f (\nabla_u v) - f (\nabla_u v)$$

$$= (\nabla_u f) v + f \cdot (\nabla_u v - \nabla_u v)$$

$$= \mathbb{L}_u f \cdot v + f \cdot \mathbb{L}_u v$$

Leibnitz rule

$$\mathbb{L}_u \langle w, v \rangle = \mathbb{L}_u (w_\alpha v^\alpha) = \langle \mathbb{L}_u w, v \rangle + \langle w, \mathbb{L}_u v \rangle$$

$$= \nabla_u \langle w, v \rangle = \langle \nabla_u w, v \rangle + \langle w, \nabla_u v \rangle$$

$$\langle \mathbb{L}_u w, v \rangle = \langle \nabla_u w, v \rangle + \langle w, \nabla_u v \rangle - \langle w, \nabla_u v - \nabla_u v \rangle$$

$$\langle \mathbb{L}_u w, v \rangle = (\nabla_u w)_\alpha v^\alpha + w_\alpha (\nabla_u v^\alpha)$$

$$\Rightarrow \mathbb{L}_u w_\mu = u^\alpha \nabla_\alpha w_\mu + w_\alpha \nabla_\mu u^\alpha$$

(2)

covariant $(L_u w)_\mu = u^\alpha w_{\mu;\alpha} + w_\alpha u^\alpha_{;\mu}$

$$= u^\alpha (w_{\mu,\alpha} - \Gamma^\lambda_{\mu\alpha} w_\lambda) + w_\alpha (u^\alpha_{;\mu} + \Gamma^\alpha_{\mu\lambda} u^\lambda)$$

$$= u^\alpha w_{\mu,\alpha} + w_\alpha u^\alpha_{;\mu} - \Gamma^\lambda_{\mu\alpha} w_\lambda u^\alpha + \Gamma^\alpha_{\mu\lambda} w_\alpha u^\lambda$$

I defined ~~this~~ w/o covariant ~~derivative~~ ∇_μ

Second rank tensor each index

$$L_u T_{\mu\nu} = u^\sigma T_{\mu\nu;\sigma} + u^\sigma_{; \mu} T_{\sigma\nu} + u^\sigma_{; \nu} T_{\mu\sigma}$$

metric

$$L_u g_{\mu\nu} = u^\sigma g_{\mu\nu;\sigma} + u^\sigma_{; \mu} g_{\sigma\nu} + u^\sigma_{; \nu} g_{\mu\sigma}$$

$$L_u g_{\mu\nu} = u_{\mu;\nu} + u_{\nu;\mu}$$

K_μ = Killing vector

$$L_{K_\mu} g_{\nu\sigma} = 0$$

$$K^\mu \nabla_\mu R = 0$$

3.

Inner workings of Lie derivative.

Change of coordinate

(A) Move to new (shifted) point.

(B) Change name of coordinate at same point.

Scalar: $x'^M = x^M + \epsilon u^M$

$$\phi' = \phi(x') = \phi(x + \epsilon u) = \phi(x) + \epsilon u^\alpha \frac{\partial \phi}{\partial x^\alpha}$$

$$\lim_{\epsilon \rightarrow 0} \frac{\phi' - \phi}{\epsilon} = u^\alpha \frac{\partial \phi}{\partial x^\alpha} = \mathcal{L}_u \phi$$

Vector $v'^M = v^M + \frac{\partial v^M}{\partial x^\alpha} \epsilon u^\alpha$

change coordinate value. $x''^\mu = x^\mu - \epsilon u^\mu$

$$v''^\mu = (v'^\alpha) \left(\frac{\partial x''^\mu}{\partial x'^\alpha} \right) = v'^\alpha \left(\delta^\mu_\alpha - \epsilon \frac{\partial u^\mu}{\partial x^\alpha} \right)$$

$$= \left(v^\alpha + \frac{\partial v^\alpha}{\partial x^\beta} \epsilon u^\beta \right) \left(\delta^\mu_\alpha - \epsilon \frac{\partial u^\mu}{\partial x^\alpha} \right)$$

$$= v^\mu + \frac{\partial v^\mu}{\partial x^\beta} \epsilon u^\beta - v^\alpha \epsilon \frac{\partial u^\mu}{\partial x^\alpha}$$

$$\lim_{\epsilon \rightarrow 0} \frac{v''^\mu - v^\mu}{\epsilon} = u^\beta \frac{\partial v^\mu}{\partial x^\beta} - v^\alpha \frac{\partial u^\mu}{\partial x^\alpha} = [\mathcal{L}_u v]^\mu$$

(4)

Schwarzschild revisited. $\left\{ \begin{array}{l} \text{unique mess.} \\ \text{sources} \end{array} \right.$

spherical symmetry: (θ, ϕ)

pitch, yaw, roll

$$R = \frac{d}{dt}$$

$$S = \omega \sin \theta \frac{d}{d\theta} - \omega \sin \theta \frac{d}{d\phi}$$

$$T = -\sin \theta \frac{d}{d\theta} - \omega \cos \theta \frac{d}{d\phi}$$

$$[k_i, k_j] = \epsilon_{ij} k_k$$

Integral curves

$$\frac{dx^i}{dt} = k^i$$

trace out spheres: ("submanifold")

4D spacetime "subfoliated" into spheres.

$$ds^2 = \underline{dt^2(a,b)} + r^2(a,b) (d\theta^2 + \sin^2 \theta d\phi^2)$$

(indep. of θ, ϕ
(else depends on
direction))

\uparrow $\underline{dt^2(a,b)} \equiv \underline{dt^2(a_0, b_0)}$
else "lumpy"

$$ds^2 = g_{aa} da^2 + g_{ab} (da db + db da) + g_{bb} db^2 + \underline{r^2 d\Omega^2}$$

$$= g_{tt} dt^2 + g_{tr} (dt dr + dr dt) + g_{rr} dr^2$$

$r(a,b)$ inverted to $b(a,r)$: trade $b \leftrightarrow r$
coord.

(5)

Can choose (t^*) so that there is no $dr dt^*$.

$$dt^* = \frac{\partial t^*}{\partial r} dr + \frac{\partial t^*}{\partial t} dt.$$

$$g_{tt}^* dt^{*2} + \cancel{2g_{tr}^* dt^* dr} + g_{rr}^* dr^2 \quad (\text{want}).$$

$$\Rightarrow g_{tt}^* \left[\left(\frac{\partial t^*}{\partial t} \right)^2 dt^2 + \left(\frac{\partial t^*}{\partial t} \right) \left(\frac{\partial t^*}{\partial r} \right) (dr dt + dt dr) + \left(\frac{\partial t^*}{\partial r} \right)^2 dr^2 \right] + g_{rr}^* dr^2$$

$$= \underline{g_{tt}^*} dt^2 + \underline{g_{tr}^*} (dr dt + dt dr) + \underline{g_{rr}^*} dr^2$$

$$\underline{g_{tt}^*} \left(\frac{\partial t^*}{\partial t} \right)^2 = \underline{g_{tt}}$$

3 conditions

$$\underline{g_{tt}^*} \left(\frac{\partial t^*}{\partial t} \right) \left(\frac{\partial t^*}{\partial r} \right) = \underline{g_{tr}}$$

3 functions

$$\underline{g_{tt}^*} \left(\frac{\partial t^*}{\partial r} \right)^2 + \underline{g_{rr}^*} = \underline{g_{rr}}$$

$\frac{g_{tt}^*}{g_{rr}^*}$
 $t(a, r)$