

### PHY 6646 Spring 2003 – Homework 3

**Due by 5 p.m. on Friday, January 31. Partial credit will be available for solutions submitted by 5 p.m. on Monday, February 3.**

*Answer both questions. To gain maximum credit you should explain your reasoning and show all working. Please write neatly and include your name on the front page of your answers.*

1. Consider an isolated point-like particle of mass  $m$  and charge  $q$ , situated in a static, uniform magnetic field  $\mathbf{B}$ , and described by the Hamiltonian

$$H = \frac{1}{2m} \left| \mathbf{P} - \frac{q}{c} \mathbf{A} \right|^2 + q\phi. \quad (1)$$

Any choice of gauge for the electromagnetic potential can be expressed as a local gauge transformation applied to the symmetric Coulomb gauge:

$$\mathbf{A}^{(\Lambda)}(\mathbf{r}, t) = \frac{1}{2} \mathbf{B} \times \mathbf{r} - \nabla \Lambda(\mathbf{r}, t), \quad \phi^{(\Lambda)}(\mathbf{r}, t) = \frac{1}{c} \frac{\partial \Lambda(\mathbf{r}, t)}{\partial t},$$

where each gauge is specified by a different real-valued function  $\Lambda(\mathbf{r}, t)$ . The real-space wave function corresponding to an abstract state  $|\psi\rangle$  is also gauge-dependent; let us denote it  $\psi^{(\Lambda)}(\mathbf{r}, t)$ .

- (a) Substitute the coordinate-representations of  $\mathbf{P}$ ,  $\mathbf{A}^{(\Lambda)}$ , and  $\phi^{(\Lambda)}$  into the Hamiltonian (1), expand all parentheses, and then write the resulting expression in the form  $H^{(\Lambda)} = H_0^{(\Lambda)} + H_{\mathbf{B}}^{(\Lambda)}$ , where  $H_0^{(\Lambda)}$  is the Hamiltonian for the case  $\mathbf{B} = \mathbf{0}$ .
- (b) Find the particle's orbital magnetic moment operator, defined (by analogy with classical electromagnetism) as  $\boldsymbol{\mu}_L^{(\Lambda)} = -\partial H^{(\Lambda)}(\mathbf{B}') / \partial \mathbf{B}' \Big|_{\mathbf{B}'=\mathbf{0}}$ . Make it resemble as closely as possible its form in the gauge  $\Lambda = 0$ , i.e.,  $\boldsymbol{\mu}_L^{(0)} = q/(2mc)\mathbf{L}$ , where  $\mathbf{L}$  is the orbital angular momentum.
- (c) Show that, even though the orbital angular momentum operator defined in part (b) is gauge-dependent, the matrix element of  $\boldsymbol{\mu}_L^{(\Lambda)}$  between any pair of states  $\psi^{(\Lambda)}(\mathbf{r}, t)$  and  $\chi^{(\Lambda)}(\mathbf{r}, t)$  is gauge-independent.
- (d) Show that if the magnetic operator is given the alternative definition  $\boldsymbol{\mu}_L^{(\Lambda)} = -\partial H^{(\Lambda)}(\mathbf{B}') / \partial \mathbf{B}' \Big|_{\mathbf{B}'=\mathbf{B}}$ , this operator can be expressed in the gauge-independent form  $\boldsymbol{\mu}_L = (q/2c)\mathbf{R} \times \mathbf{V}$ , where  $\mathbf{V}$  is the velocity operator.
- (e)  $H_0$  found in part (a) is gauge-dependent, even though it describes a situation when no electromagnetic field is present. Explain the source of this gauge dependence. Show that the matrix element of  $H_0^{(\Lambda)} - i\hbar\partial/\partial t$  between any pair of states  $\psi^{(\Lambda)}(\mathbf{r}, t)$  and  $\chi^{(\Lambda)}(\mathbf{r}, t)$  is gauge-independent.

2. Based on Ballentine Problem 12.2: The most general state operator for a spin- $\frac{1}{2}$  system has the form  $\rho = \frac{1}{2}(I + \mathbf{p} \cdot \boldsymbol{\sigma})$ , where  $\sigma_j$  ( $j = 1, 2, 3$ ) is a Pauli operator and  $\mathbf{p}$  is the (real) polarization vector, whose length cannot exceed 1. Suppose that the system has a magnetic moment  $\boldsymbol{\mu} = \frac{1}{2}\gamma\hbar\boldsymbol{\sigma}$  and is in a constant, uniform magnetic field  $\mathbf{B}$ .
- (a) Calculate the time dependence of  $\rho(t)$  in the Schrödinger picture. Describe the result geometrically in terms of the variation of the vector  $\mathbf{p}$ , assuming a general value for  $\mathbf{p}(t = 0)$ .
- (b) Now suppose that the system is in thermal equilibrium at temperature  $T$ , so that  $\rho_H \propto \exp(-H/k_B T)$ . Find  $\mathbf{p}(t)$  for this situation, and hence deduce  $\langle \boldsymbol{\mu}(t) \rangle$ .