

SUSY Dark Matter?

Konstantin Matchev

University of Florida



SpacePart 2003 Conference,
Washington DC, December 10 2003

Dark matter and physics beyond the Standard Model

- Dark matter is our best evidence for new physics beyond the Standard Model
- WIMPs are motivated by both particle physics and astrophysics.
 - Predicted in many particle theories BSM.
 - Give the right order of magnitude Ω_{DM} .

$$\Omega_{DM} h^2 \sim 0.1 \left(\frac{\sigma_{EW}}{\sigma_{ann}} \right)$$

- Since they must have been able to annihilate in the Early Universe, they should also produce observable signals in direct and indirect dark matter detection experiments.
- Recipe for BSM dark matter
 - invent a new model
 - invent a symmetry which guarantees a stable particle
 - fudge parameters until the lightest new stable particle is neutral and has the correct relic density



Outline

- Three generic examples of theories BSM with DM candidates
 - supersymmetry: DM = lightest superpartner.
 - extra dimensions: DM = lightest Kaluza-Klein mode.
 - Little Higgs: DM = ? (billion, gagesino, ...)
- In each case:
 - brief introduction to the model.
 - who is the DM particle?
 - why is it stable?
 - what is the preferred mass range for DM?
 - DM discovery prospects.
- How can we identify the DM particle?
 - high-energy coliders
 - astroparticle physics experiments in space (Jonathan's talk)



Supersymmetry

- Supersymmetry is an extra dimension theory with new **anticommuting** coordinates θ_α :

$$\Phi(x^\mu, \theta) = \phi(x^\mu) + \psi^\alpha(x^\mu)\theta_\alpha + F(x^\mu)\theta^\alpha\theta_\alpha$$

- SUSY relates SM particles and their superpartners ($\phi \leftrightarrow \psi$)
 - quarks, leptons \Leftrightarrow squarks, sleptons
 - gauge bosons: $g, W^\pm, W_3^0, B^0 \Leftrightarrow$ gauginos: $\tilde{g}, \tilde{w}^\pm, \tilde{w}^0, \tilde{b}^0$
 - Higgs bosons: $h^0, H^0, A^0, H^\pm \Leftrightarrow$ higgsinos: $\tilde{h}^\pm, \tilde{h}_u^0, \tilde{h}_d^0$
- The superpartners have
 - spins differing by 1/2
 - identical couplings
 - unknown masses (model-dependent)
- The superpartners are charged under a conserved R -parity
 - SM particles: $R = 0$
 - superpartners: $R = -1 \implies$ stable LSP (DM?).
- No tree-level contributions to precision EW observables



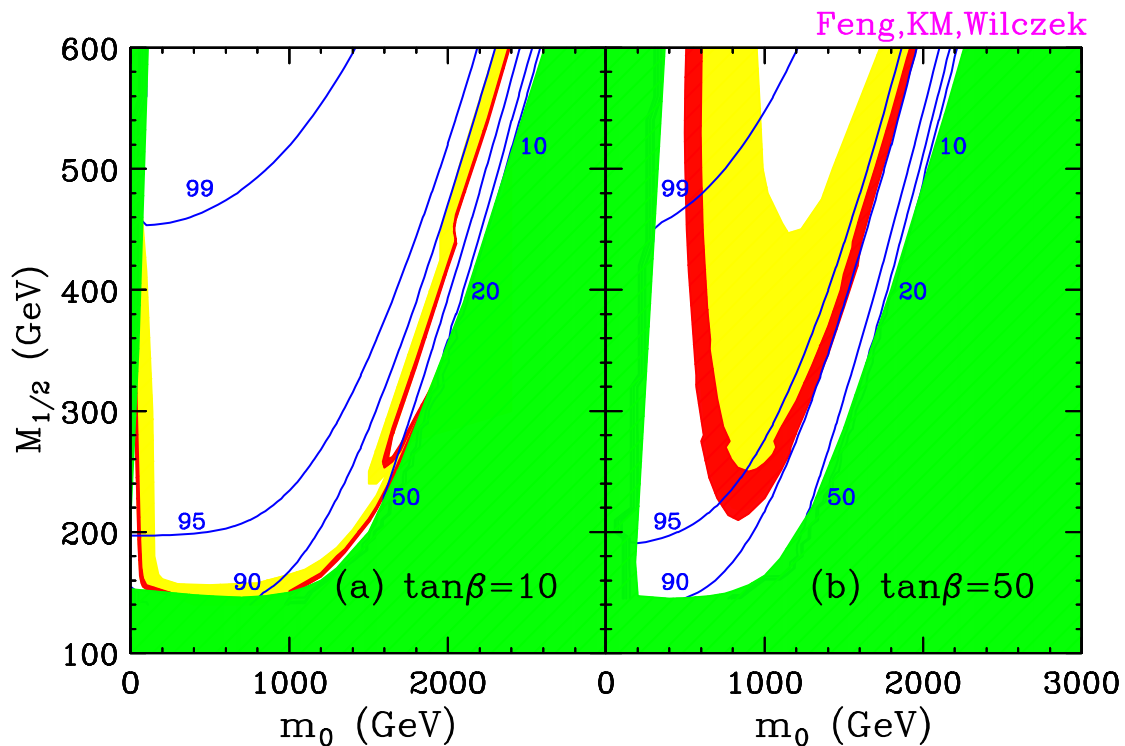
Gaugino fraction of SUSY WIMPs

- The lightest neutralino $\tilde{\chi}_1^0$ is a mixture of \tilde{b}^0 , \tilde{w}^0 , \tilde{h}_u^0 , \tilde{h}_d^0 :

$$\tilde{\chi}_1^0 = a_1 \tilde{b}^0 + a_2 \tilde{w}^0 + a_3 \tilde{h}_u^0 + a_4 \tilde{h}_d^0$$

- Gaugino fraction R_χ of the LSP:

$$R_\chi \equiv |a_1|^2 + |a_2|^2 \approx |a_1|^2.$$

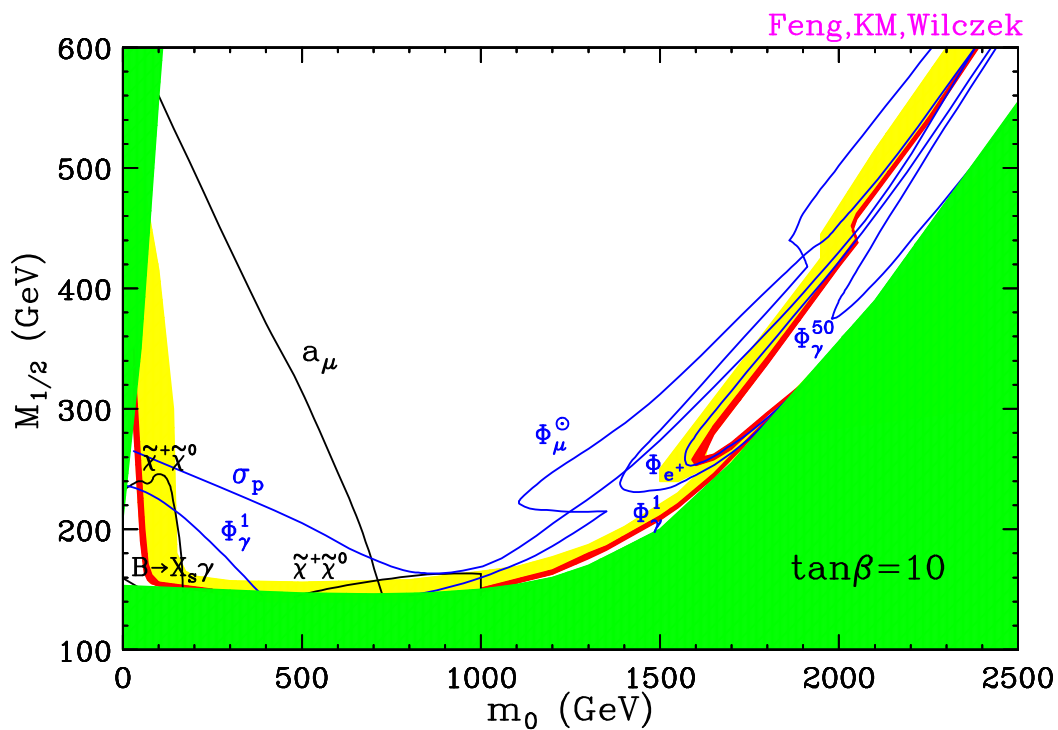


- Focus point region: large m_0 , mixed LSP.
- Coannihilation region: small m_0 , $\tilde{\tau} - \tilde{\chi}_1^0$ degeneracy.



SUSY WIMP detection

- Combination of “all” pre-LHC experiments
 - Direct SUSY searches: Tevatron
 - Indirect SUSY searches: E827, B-factories
 - Direct WIMP searches: CDMS, CRESST, GENIUS
 - Indirect WIMP searches: Amanda, AMS, GLAST



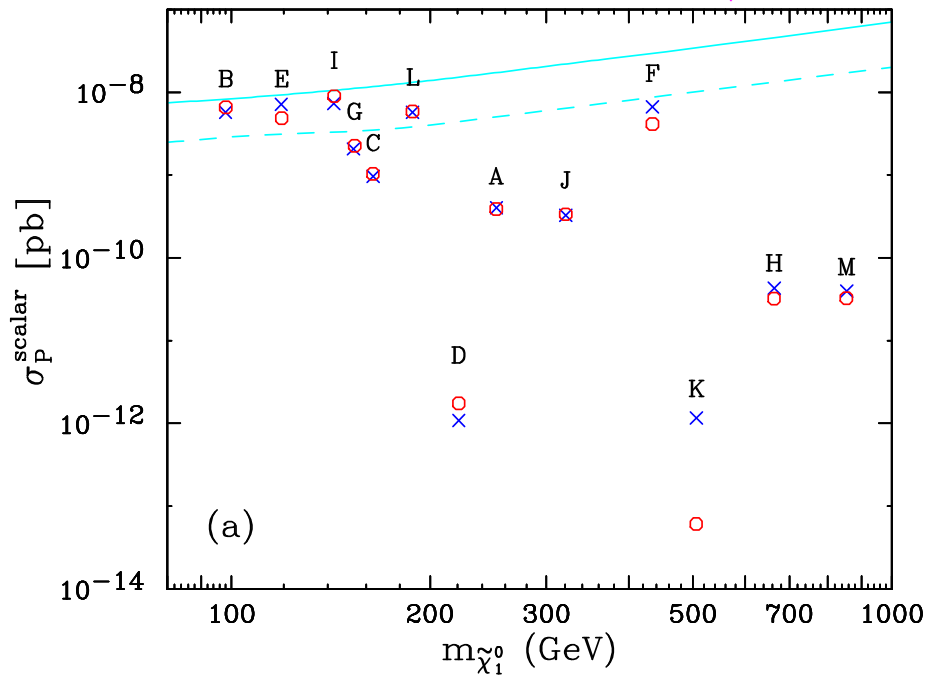
- Many possible DM signals before 2007-08.
- Particle physics and astrophysics probes are highly complementary.



SUSY DM: direct detection

- Spin-independent cross-sections for the 13 benchmark points of Battaglia et al. hep-ph/0106204.

Ellis, Feng, Ferstl, KM, Olive hep-ph/0110225



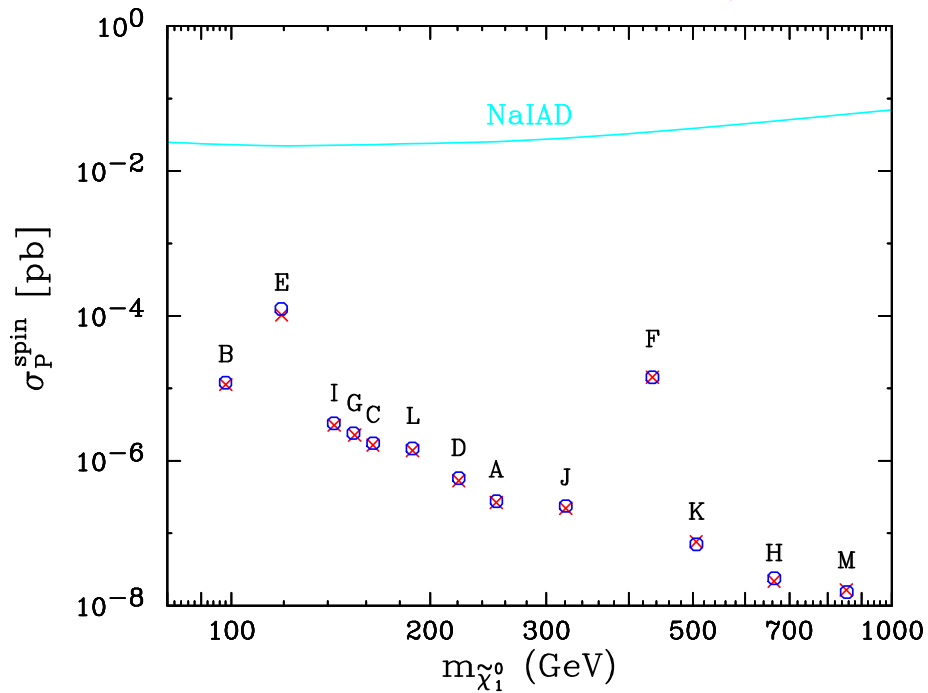
- No lower limit: cancellations are possible.



SUSY DM: direct detection

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- Far below sensitivity of near-term future experiments.



Universal Extra Dimensions

Appelquist, Cheng, Dobrescu, hep-ph/0012100

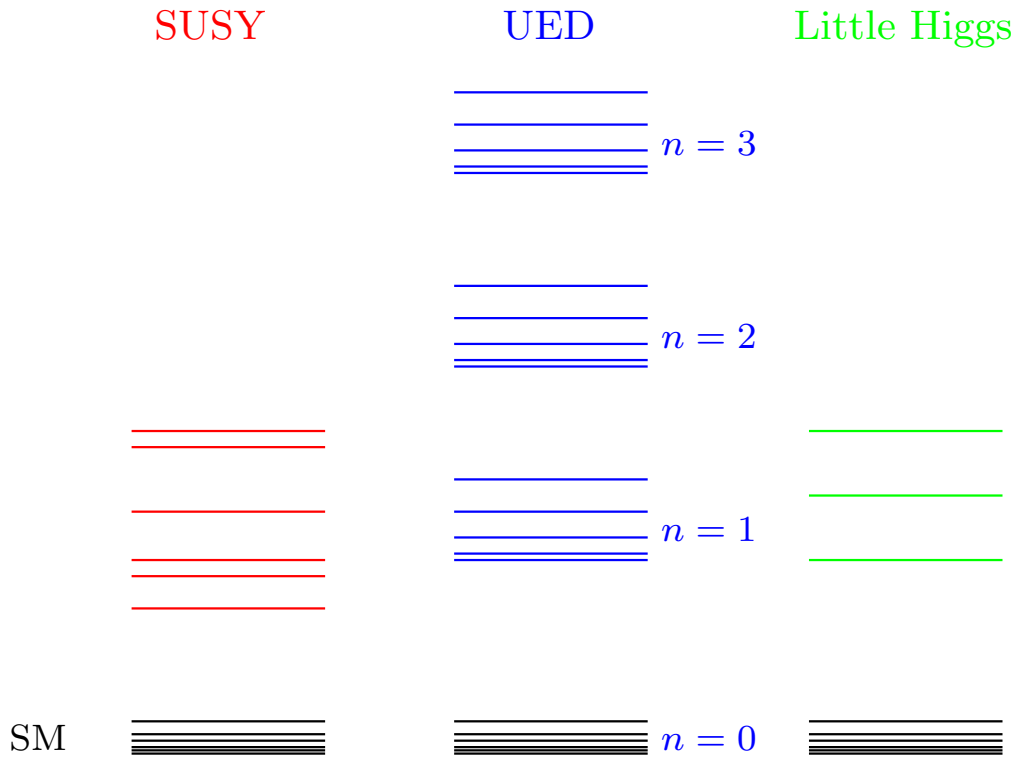
- Universal Extra Dimensions is an extra dimension theory with new **bosonic** coordinates y (spanning a circle of radius R):

$$\Phi(x^\mu, y) = \phi(x^\mu) + \sum_{i=1}^{\infty} \phi^n(x^\mu) \cos(ny/R) + \chi^n(x^\mu) \sin(ny/R)$$

- Each SM field ϕ ($n = 0$) has an infinite tower of Kaluza-Klein (KK) partners ϕ^n and χ^n with
 - identical spins
 - identical couplings
 - unknown masses of order n/R
- Remnant of p_5 conservation: KK -parity $(-1)^n$
 - $KK = +1$ for even n and $KK = -1$ for odd n .
 - lightest KK partner at level 1 (LKP) is stable.
 - $P_3 \rightarrow P'_3 P_0, P_2 P_1, P_1 P_0;$
 - $P_2 \rightarrow P'_2 P_0, P_1 P_1, P_0 P_0;$
 - $P_1 \rightarrow P'_1 P_0.$
- No tree-level contributions to precision EW observables



Model summary



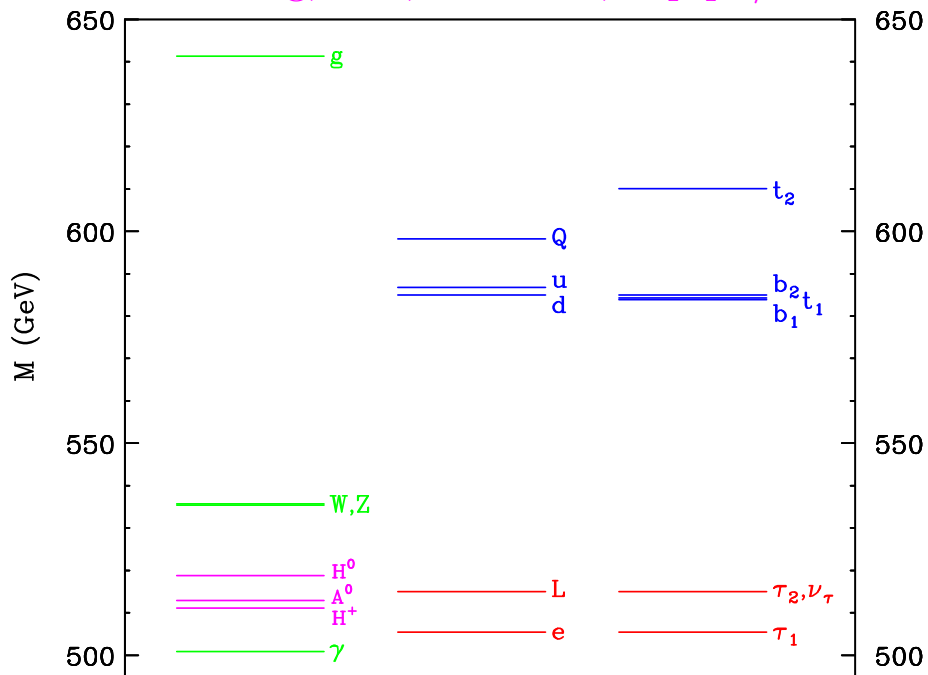
	SUSY	UED	Little Higgs
DM particle	LSP	LKP	LTP
Spin	1/2	1	0
Symmetry	<i>R</i> -parity	KK-parity	<i>T</i> -parity
Mass range	50-200 GeV	600-800 GeV	700-800 GeV



UED spectrum at level 1

- Including radiative corrections, the mass spectrum of level 1 KK modes looks something like this:

Cheng, KM, Schmaltz, hep-ph/0204342



- Mimics (fermionic) supersymmetry!
- Seems difficult to discover at the LHC, but...
- W_1^\pm, Z_1 have pure leptonic branchings!
- $\sin^2 \theta_W^1 \approx 0 \implies \gamma^1 \approx B^1$, similar to \tilde{B} in SUSY.



The KK Weinberg angles

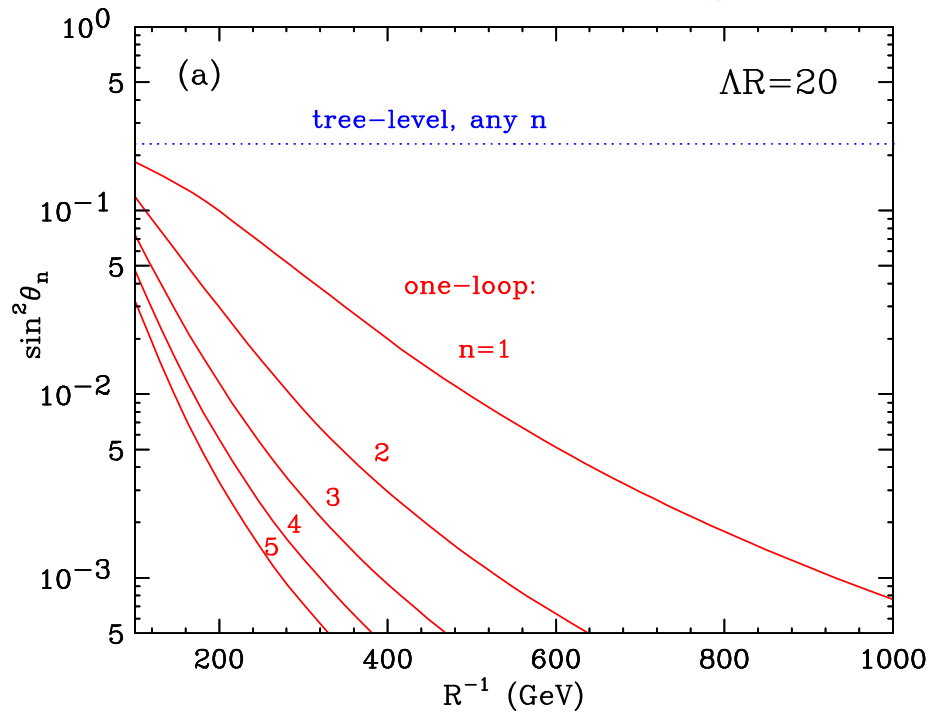
- Mass matrix for the neutral gauge bosons

$$\begin{pmatrix} \frac{n^2}{R^2} + \frac{1}{4}g_1^2v^2 + \hat{\delta}m_{B_n}^2 & \frac{1}{4}g_1g_2v^2 \\ \frac{1}{4}g_1g_2v^2 & \frac{n^2}{R^2} + \frac{1}{4}g_2^2v^2 + \hat{\delta}m_{W_n}^2 \end{pmatrix}$$

- The Weinberg angle θ_n at KK level n

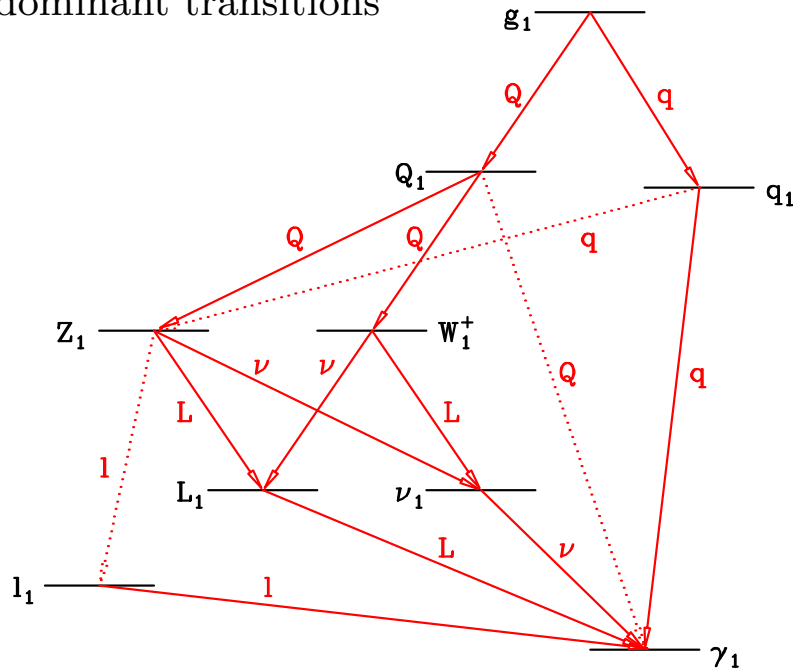
$$\gamma_n = \cos \theta_n B_n^0 + \sin \theta_n W_n^0 \approx B_n^0.$$

Cheng, KM, Schmaltz, hep-ph/0204342



Level 1 Spectroscopy

- Allowed dominant transitions



- KK gluon: $B(g_1 \rightarrow Q_1 Q_0) \simeq B(g_1 \rightarrow q_1 q_0) \simeq 0.5$.
- Singlet KK quarks: preferentially $q \rightarrow \gamma_1 q_0$
- Doublet KK quarks:

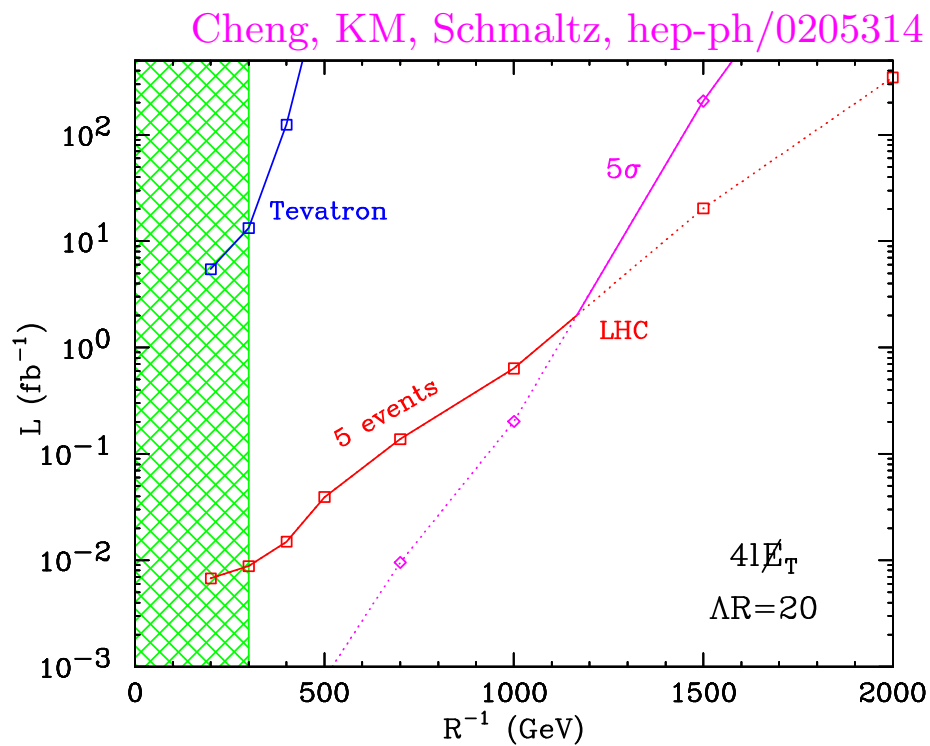
$$B(Q_1 \rightarrow W_1^\pm Q'_0) \sim 65\% \quad B(Q_1 \rightarrow Z_1 Q_0) \sim 33\%$$

- KK W - and Z -bosons: only leptonic decays!
- KK leptons: 100% directly to the LKP.
- At hadron colliders we want: **strong** production, **weak** decays!



UED discovery reach at the Tevatron and LHC

- Discovery reach in the $Q_1 Q_1 \rightarrow 4\ell \cancel{E}_T$ channel.

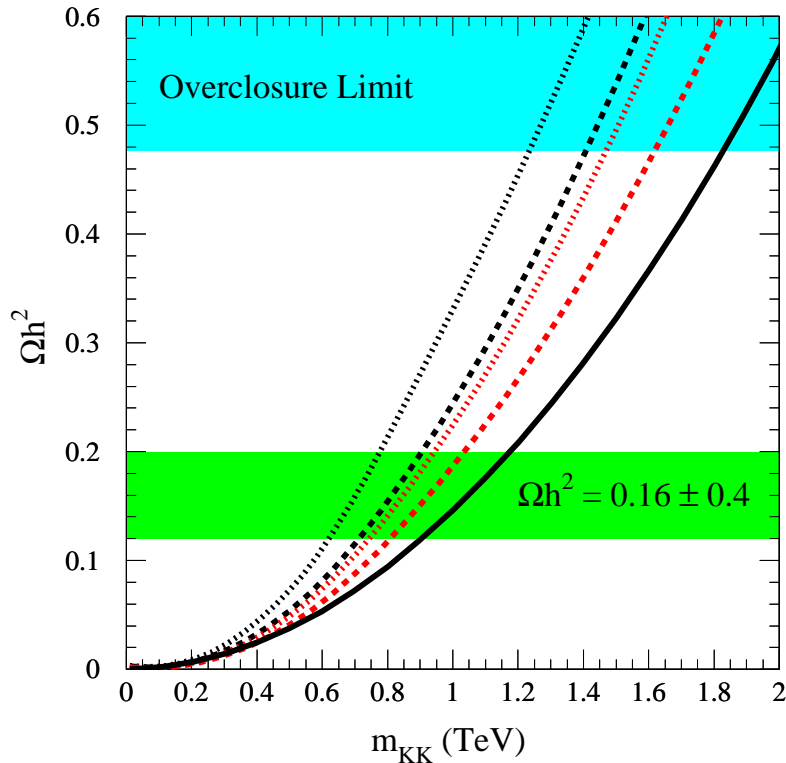


- Typical signatures include:
 - soft leptons, soft jets, not a lot of \cancel{E}_T
 - a lot of missing mass (LHC can't measure it)



Kaluza-Klein dark matter

- Relic density: G.Servant, T.Tait, hep-ph/0206071



- Unlike supersymmetry: no helicity suppression

$$\Omega h^2 = \frac{1.04 \cdot 10^9 \text{ GeV}^{-1}}{M_P \sqrt{g_*}} \frac{x_F}{a + 3b/x_F}; \quad x_F = \frac{M_{KK}}{T_F}$$

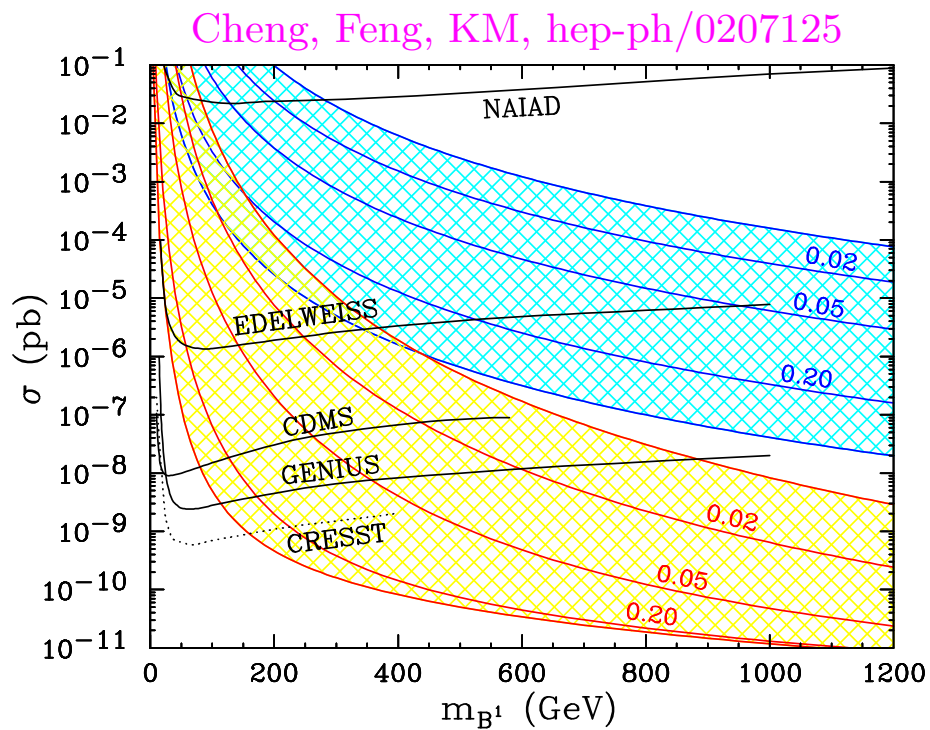
$$a = \frac{\alpha_1^2}{M_{KK}^2} \frac{380\pi}{81}; \quad b = -\frac{\alpha_1^2}{M_{KK}^2} \frac{95\pi}{162}.$$

- Unlike supersymmetry: coannihilation lowers the bound



KK DM: direct detection

- As usual, spin-dependent and spin-independent cross-sections.



- The signals are enhanced near the s-channel resonance:
 $\sigma \sim (m_{q^1} - m_{B^1})^{-2}$. Unnatural in SUSY, guaranteed here.

Cheng, Feng, KM, hep-ph/0207125

Servant, Tait, hep-ph/0209262

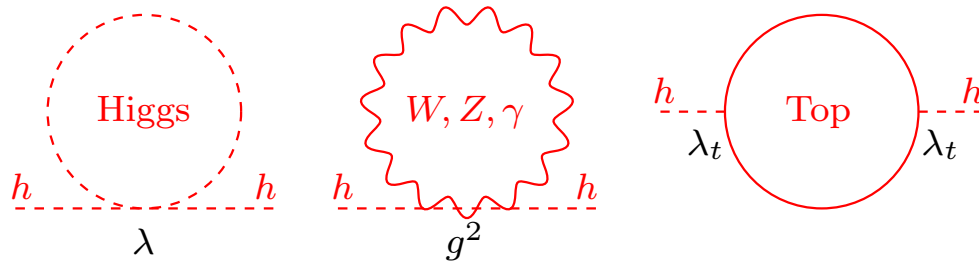
Majumdar, hep-ph/0209277

- Constructive interference: lower bound!

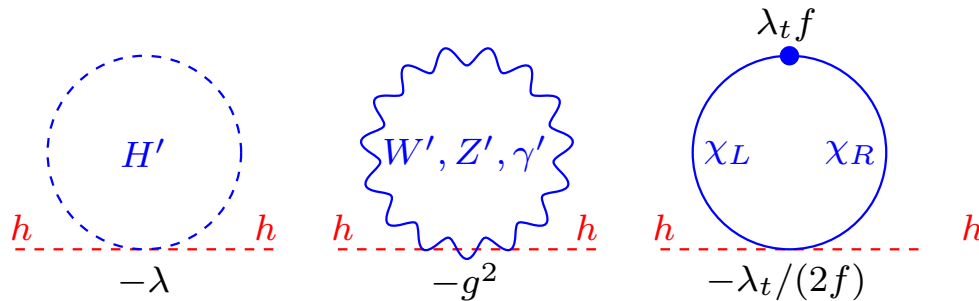


Little Higgs models

- The hierarchy problem in the SM



- Introduce new particles at TeV scale to cancel the one-loop quadratic divergences



- Conserved T -parity (Cheng, Low hep-ph/0308199)
 - $T = 0$ for **SM particles**, $T = -1$ for **new particles**.
 - the lightest T -odd particle is stable.
- No tree-level contributions to precision EW observables

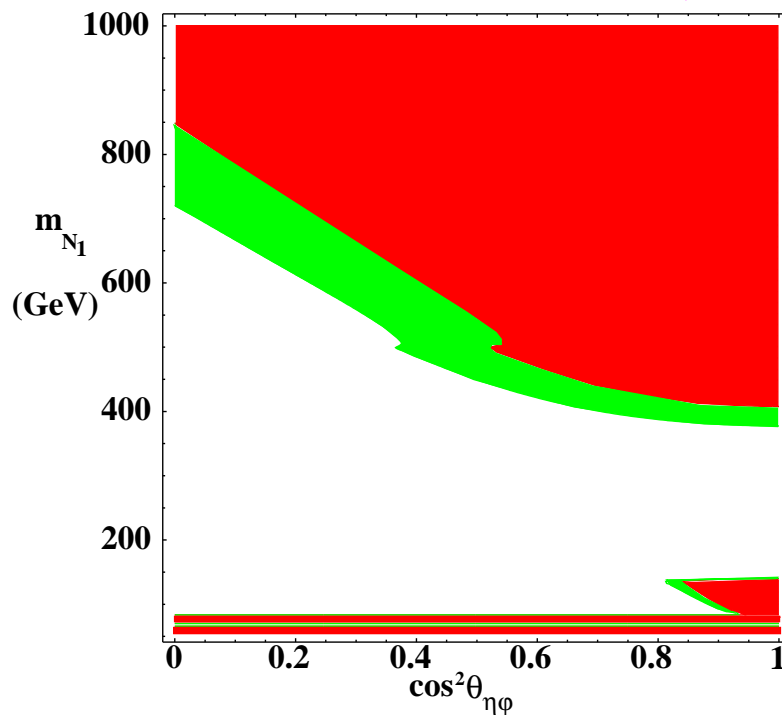


Relic density in LH models

- If the lightest T -odd particle is a scalar, it can be a mixture of an SU(2)-triplet ϕ_3^0 and an SU(2)-singlet η_1^0 :

$$N_1 = \cos \theta_{\eta\phi} \eta_1^0 + \sin \theta_{\eta\phi} \phi_3^0$$

Birkedal-Hansen, Wacker hep-ph/0306161



- The absence of helicity suppression requires large masses for the WIMP case.
- For $150 \text{ GeV} < m_{N_1} < 350 \text{ GeV}$, annihilation into $t\bar{t}$ and hh is very efficient.



SUSY versus UED at a LC

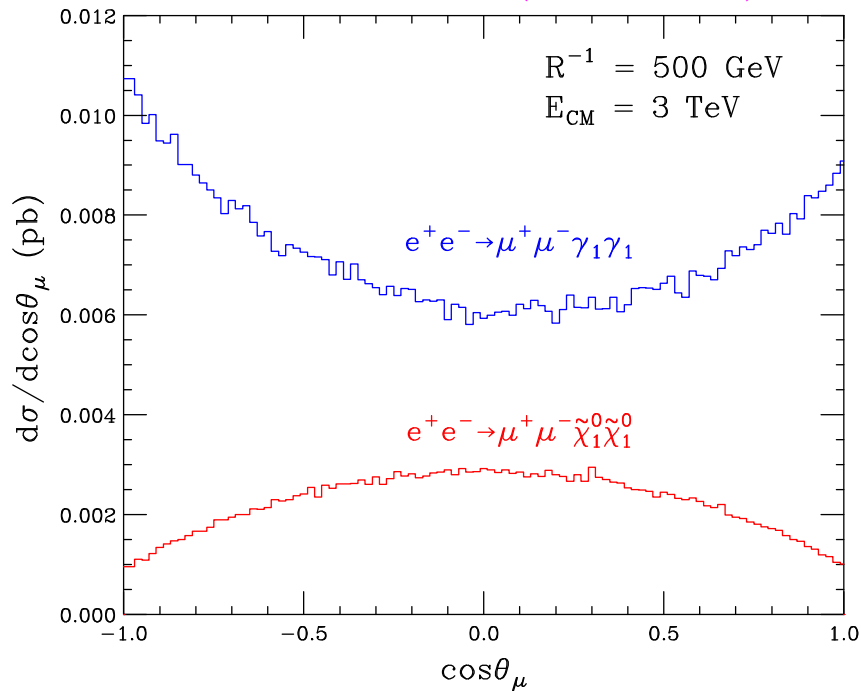
- The spin information is encoded in the angular distributions!



$$\frac{d\sigma}{d\cos\theta} \sim 1 - \cos\theta^2$$

$$\frac{d\sigma}{d\cos\theta} \sim 1 + \cos\theta^2$$

Datta, Kong, KM (preliminary)



- Significant difference in the total cross-section as well!
- The masses can be extracted from the E_μ distribution.
- Threshold scan would confirm the spins.



The Message

- Recent new ideas in particle physics lead to novel alternatives for dark matter candidates. SUSY DM? Not so fast...
- Dark matter detection experiments should be prepared for surprises, avoid theory bias.
- Elucidating the nature of the dark matter can be challenging in either direct detection experiments or hadron colliders.
- Astroparticle physics experiments in space may provide important clues to the identity of the dark matter particle.

