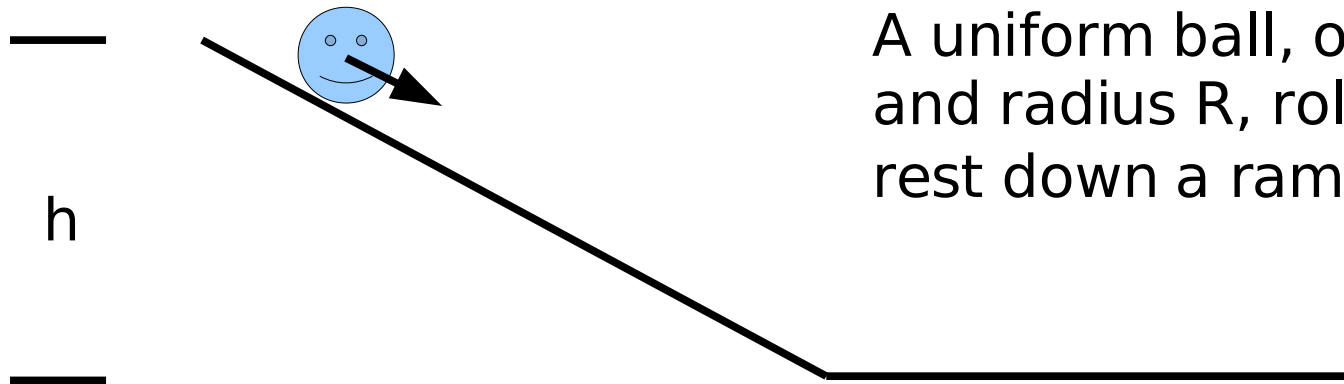


# Example

---



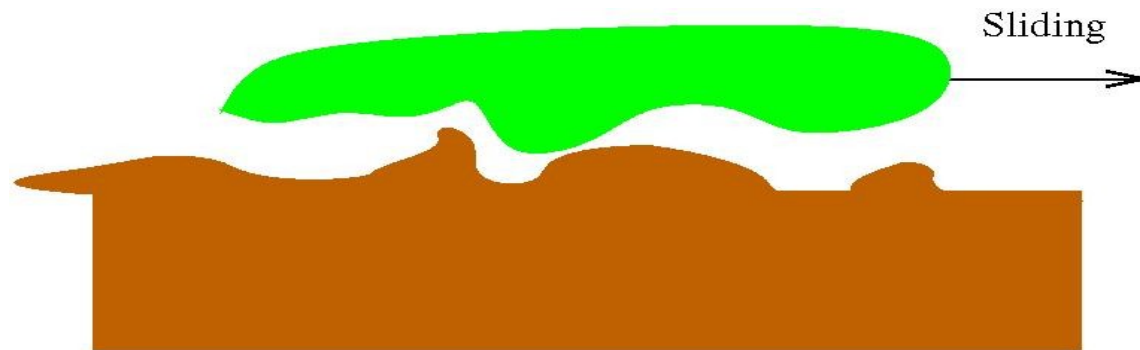
A uniform ball, of mass  $M = 6\text{kg}$  and radius  $R$ , rolls smoothly from rest down a ramp at angle  $\theta = 30^\circ$ .

- a) The ball descends a vertical height  $h = 1.2\text{m}$  to reach the bottom of the ramp. What is the speed at the bottom?

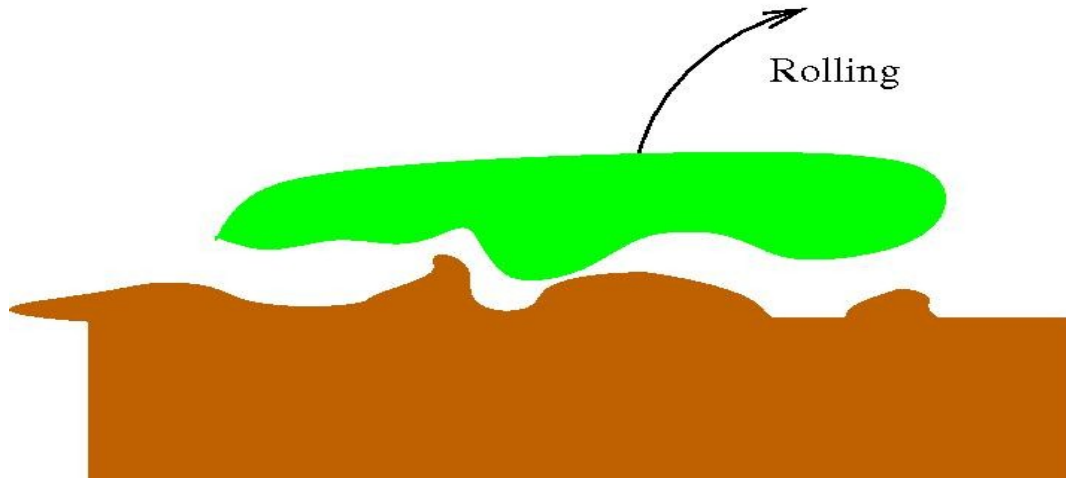
The only force doing work is gravity, a conservative force. The normal force is perpendicular to the path. The frictional force does **not transfer** any energy to thermal because the ball does **not slide** (rolls smoothly).

# Brief reminder about friction

---



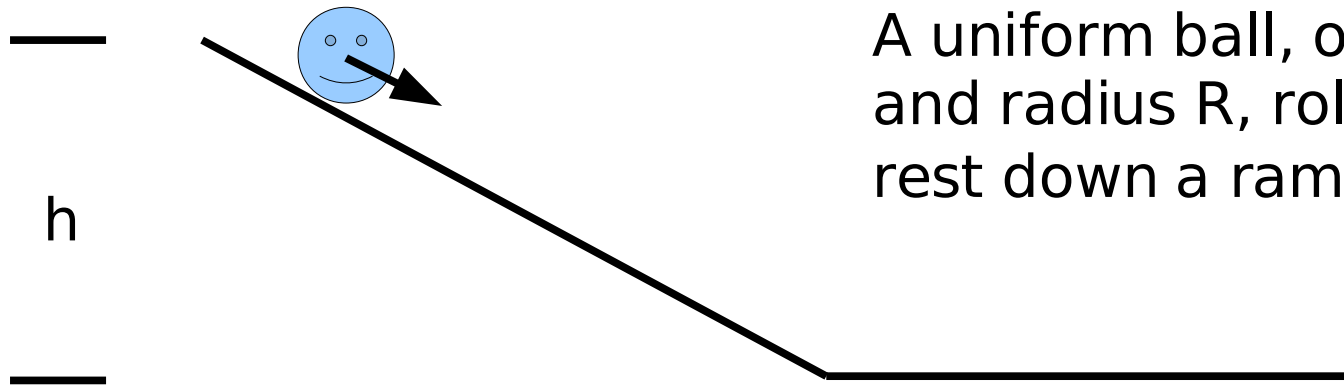
Sliding: Ripples in the surfaces get stuck and the force has to overcome this so the piece can move again -> generates heat



Smooth rolling: The 'Ripples' are moving in and out of each other without getting stuck -> Does not generate heat

Do you see the connection with static friction?

# Example



A uniform ball, of mass  $M = 6\text{kg}$  and radius  $R$ , rolls smoothly from rest down a ramp at angle  $\theta = 30^\circ$ .

- a) The ball descends a vertical height  $h = 1.2\text{m}$  to reach the bottom of the ramp. What is the speed at the bottom?

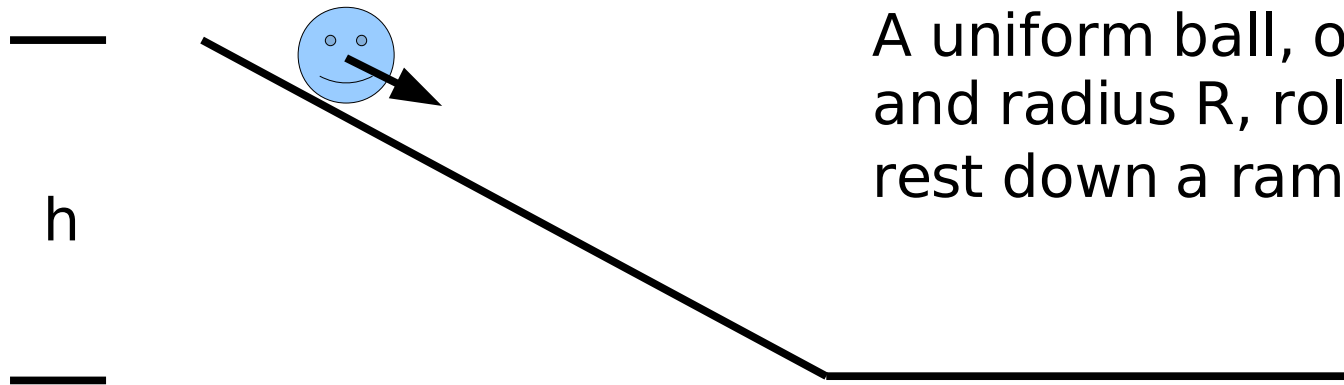
Energy is conserved:

$$\text{initial: } U_i = mgh, \quad K_i = 0 \quad \text{final: } U_f = 0, \quad K_f = 0.5I_{\text{com}}\omega^2 + 0.5mv_{\text{com}}^2$$

$$\text{Use: } v_{\text{com}} = R\omega \quad \text{and} \quad I_{\text{com}} = \frac{2MR^2}{5} \quad v_{\text{com}} = (10gh/7)^{1/2} = 4.1\text{m/s}$$

$$a_{\text{com}} = -g\sin\theta / (1 + I_{\text{com}}/MR^2) = -g\sin\theta / (1 + 2/5) = -3.5\text{m/s}^2$$

# Example



A uniform ball, of mass  $M = 6\text{kg}$  and radius  $R$ , rolls smoothly from rest down a ramp at angle  $\theta = 30^\circ$ .

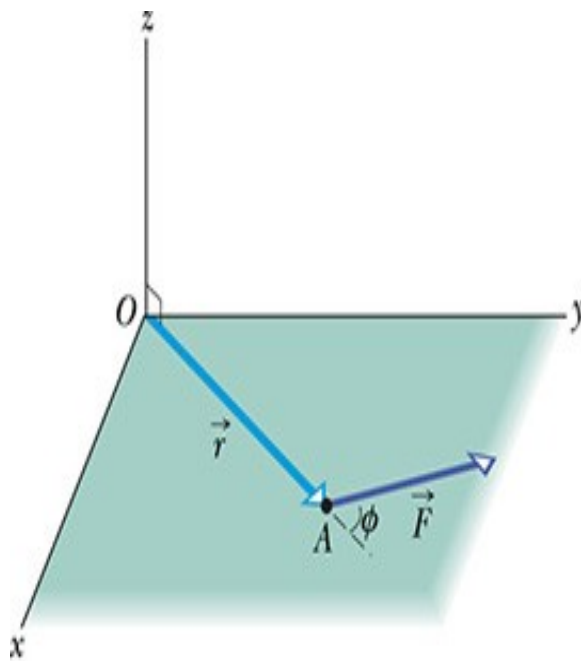
$$a_{\text{com}} = -g \sin \theta / (1 + I_{\text{com}} / MR^2) = -g \sin \theta / (1 + 2/5) = -3.5 \text{m/s}^2$$

Always independent of mass and radius if  
 $I_{\text{com}} = \text{Constant times } MR^2$

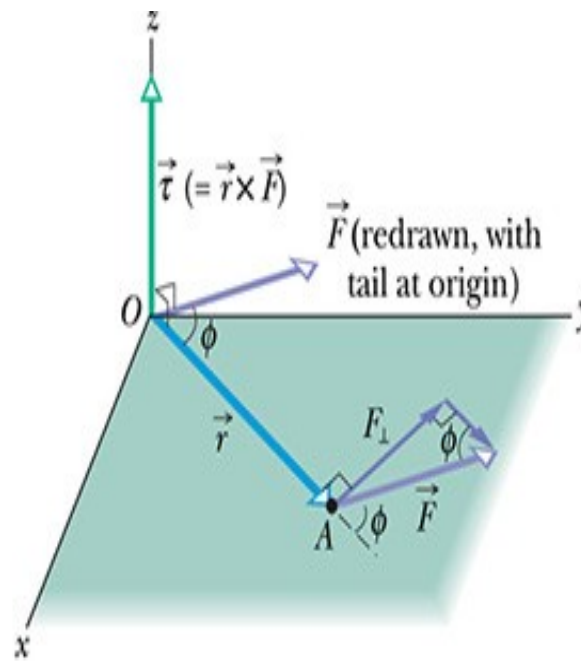
# Torque Revisited

$$\vec{\tau} = \vec{r} \times \vec{F}$$

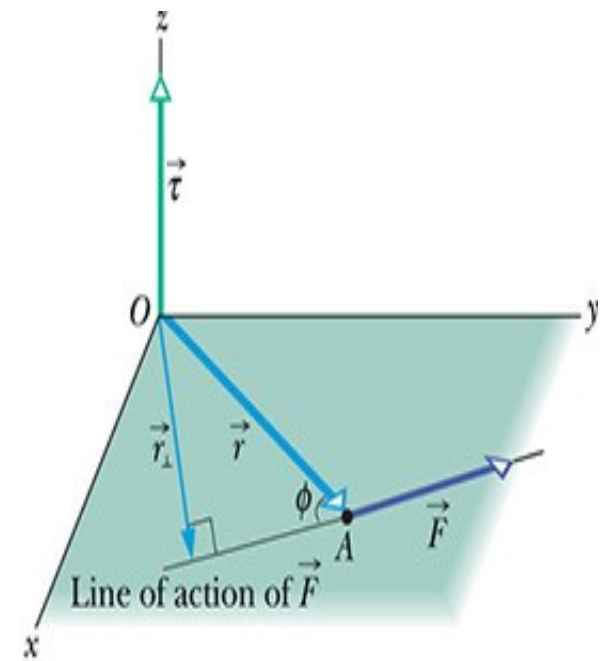
$\vec{\tau}$  starts at the origin and is perpendicular to the  $\vec{r}$ - $\vec{F}$  plane. (Right-hand rule)



(a)



(b)



(c)

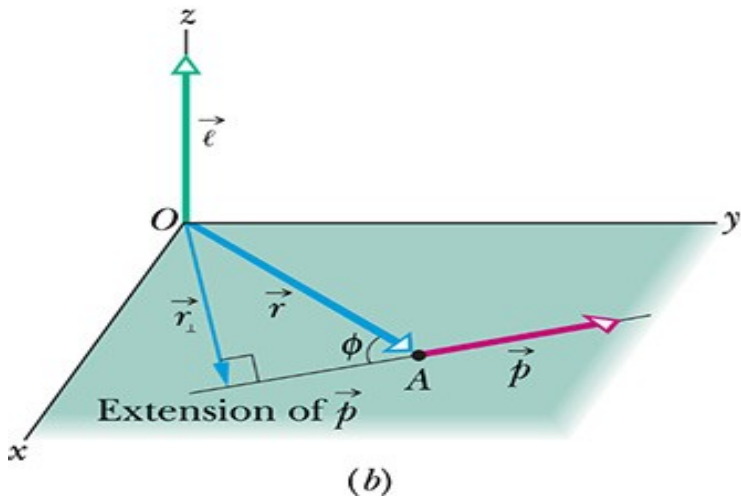
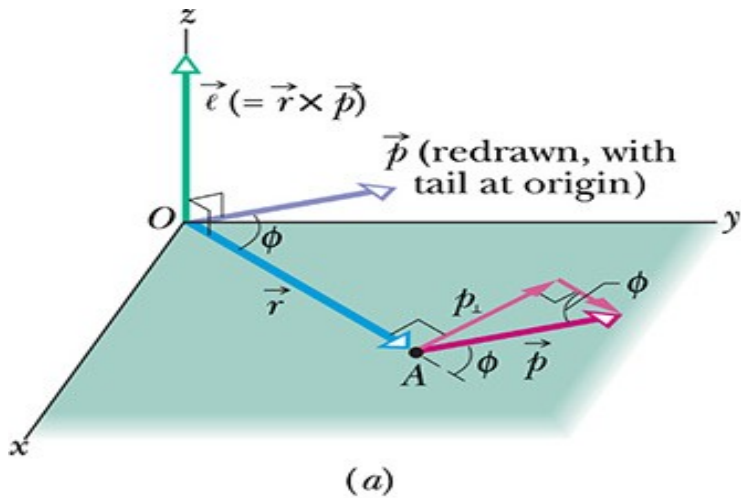
# Angular Momentum

## Angular Momentum

The angular momentum for a single particle is:

$$\vec{l} = \vec{r} \times \vec{p} = m (\vec{r} \times \vec{v})$$

$$\frac{d}{dt} \vec{l} = \vec{\tau}_{\text{net}}$$

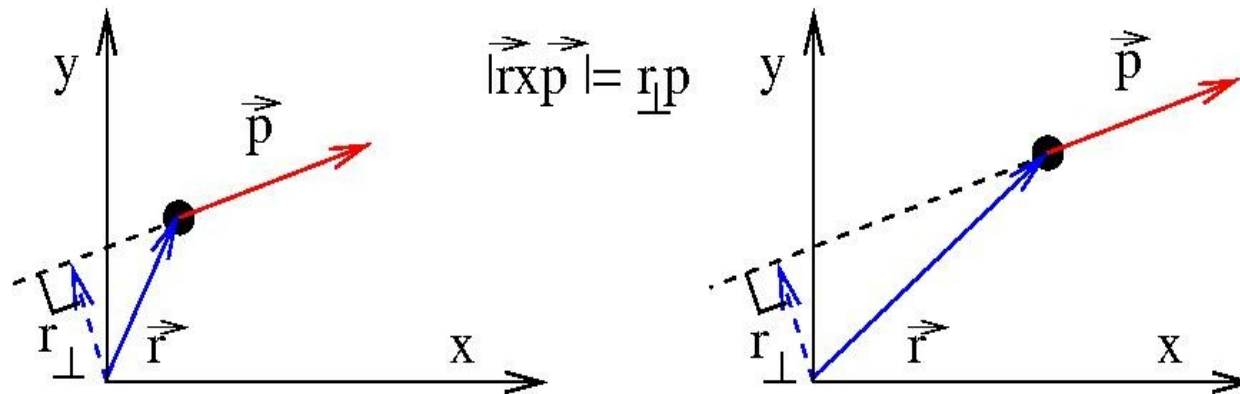


# Angular Momentum

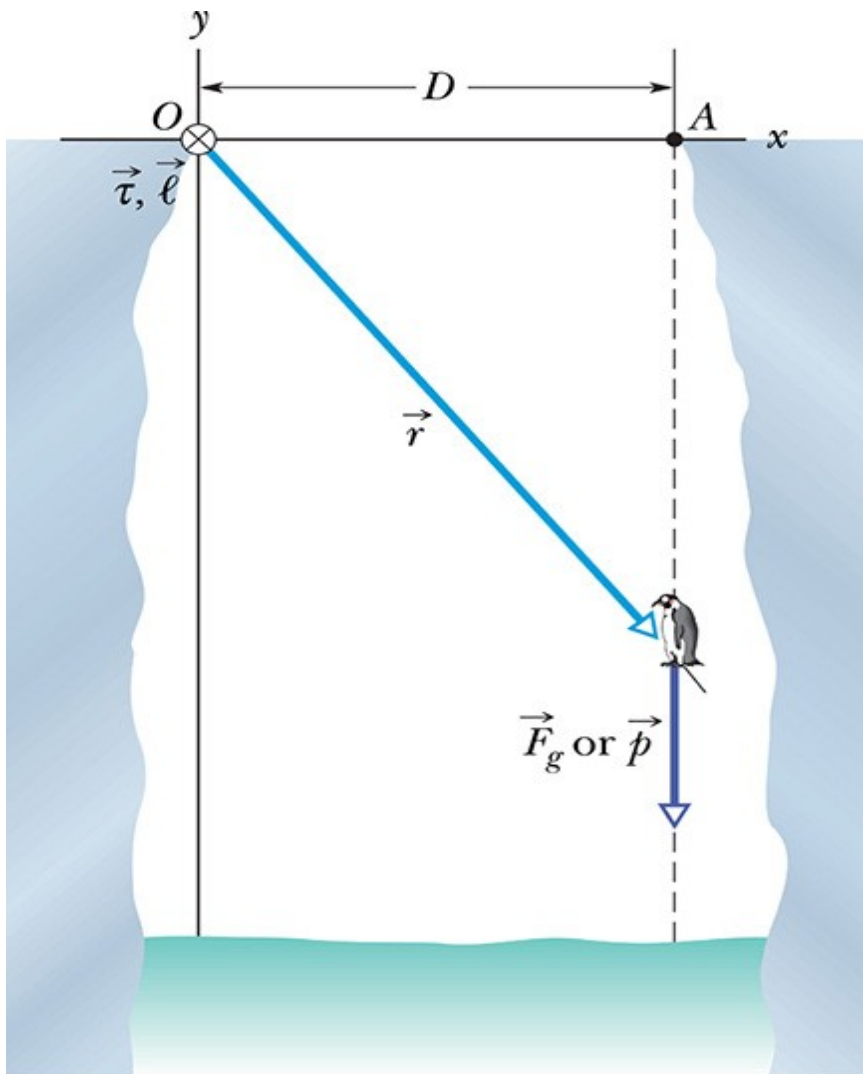
$$\frac{d}{dt} \vec{l} = \vec{\tau}_{\text{net}}$$

1. The vector sum of all torques acting on a particle is equal to the time rate of change of the angular momentum of that particle.
2. If the net torque is zero, the angular momentum is conserved. Not that this is a vector quantity and works for each component individually.

Example:



# Angular Momentum

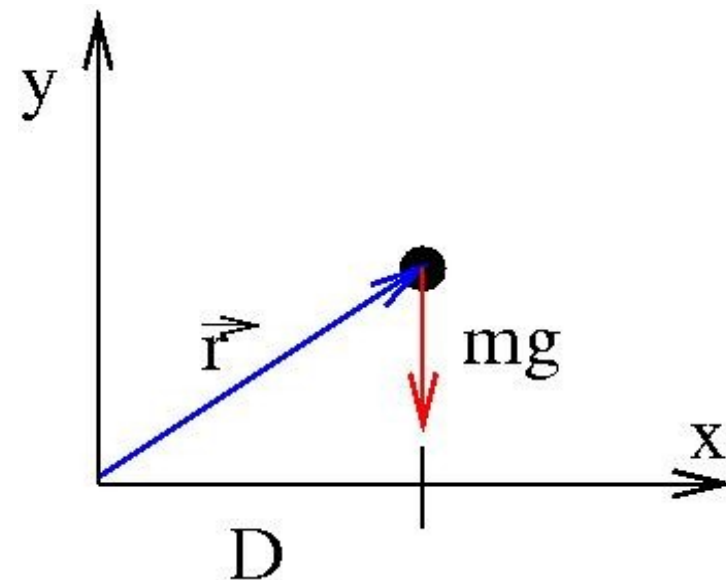


$$v = gt$$

$$L = r_{\perp} mv = Dmv = Dmgt$$

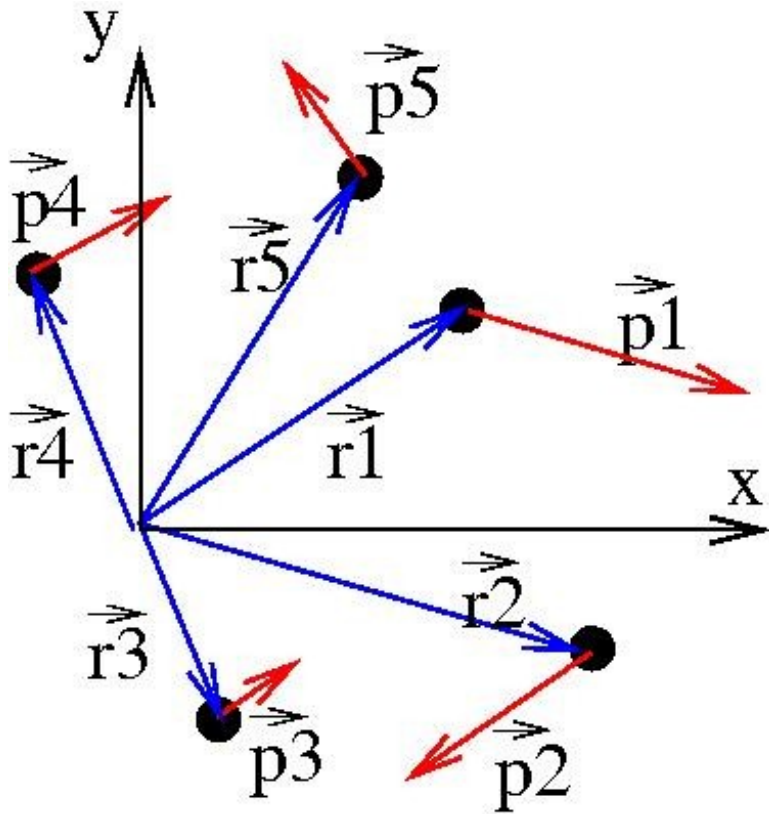
$$\tau = Dmg$$

Direction: going 'clockwise'  
-> negative z-direction





# Angular Momentum of a system of particles



$$\vec{L} = \vec{l}_1 + \vec{l}_2 + \vec{l}_3 + \vec{l}_4 + \vec{l}_5 + \dots$$

Time derivative

$$\frac{d}{dt} L = \frac{d}{dt} l_1 + \frac{d}{dt} l_2 + \dots$$

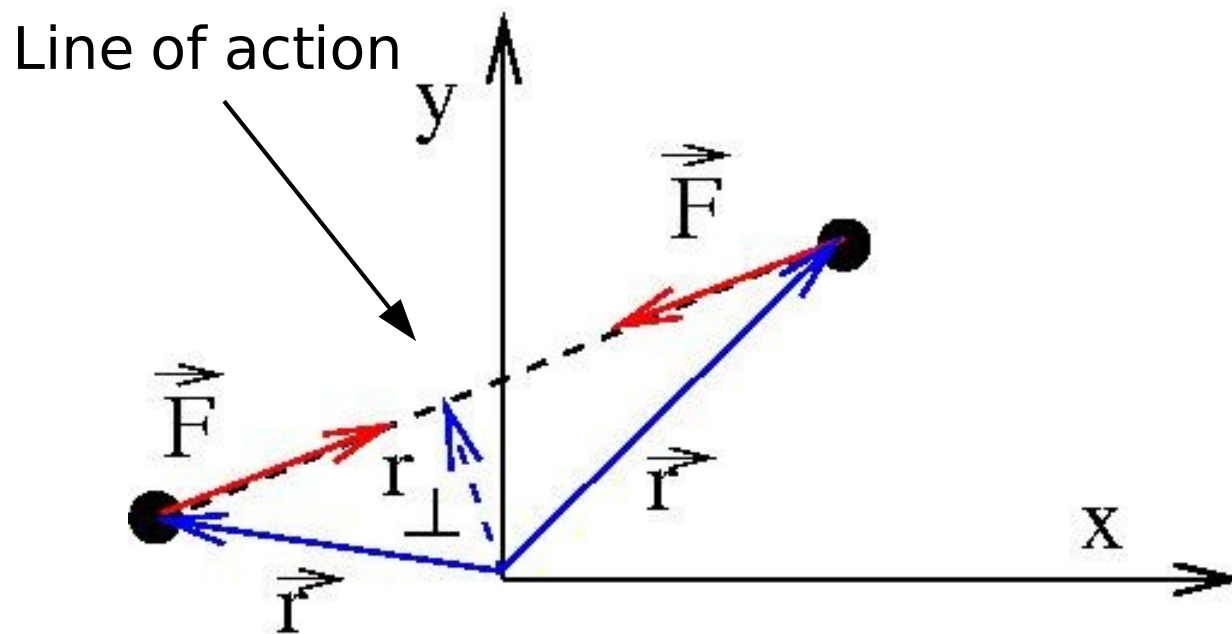
$$\frac{d}{dt} L = \sum \tau_i$$

Is the sum over all torques.

Linear momentum: All internal forces compensate each other  
Time derivative of linear momentum = sum over all net Forces  
Can we ignore the internal torques here too?

# Angular Momentum of a system of particles

Linear momentum: All internal forces compensate each other  
Time derivative of linear momentum = sum over all net Forces  
Can we ignore the internal torques here too?



Actio = Reactio

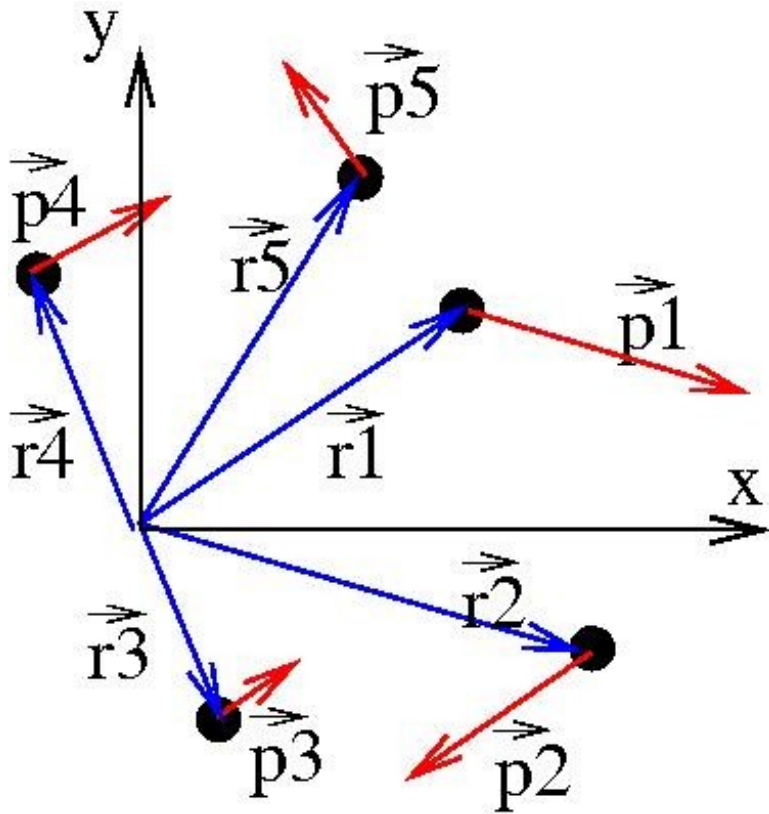
Internal Forces  
are equal and opposite  
in sign.

Lever arm  $r_{\perp}$  = equal

Torques are equal  
and opposite in sign  
(composite each other  
when summed up)

$$\tau_{\text{net}} = \sum \tau_{\text{ext}}$$

# Angular Momentum of a system of particles



$$\vec{L} = \vec{l}_1 + \vec{l}_2 + \vec{l}_3 + \vec{l}_4 + \vec{l}_5 + \dots$$

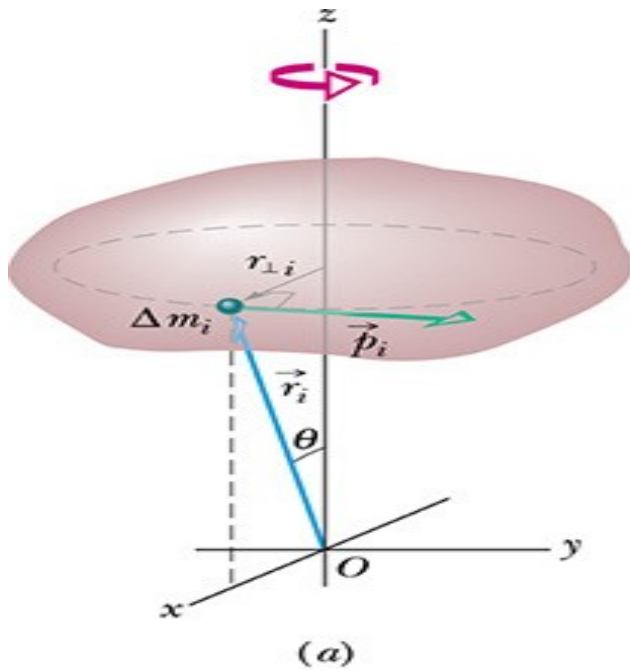
Time derivative

$$\frac{d}{dt} L = \frac{d}{dt} l_1 + \frac{d}{dt} l_2 + \dots$$

$$\frac{d}{dt} L = \sum \tau_i = \tau_{\text{net}} = \sum \tau_{\text{ext}}$$

The net torque acting on a system of particles is equal to the time rate of change of the system's total angular momentum.

# Angular Momentum of a Rigid body rotating around a fixed axis

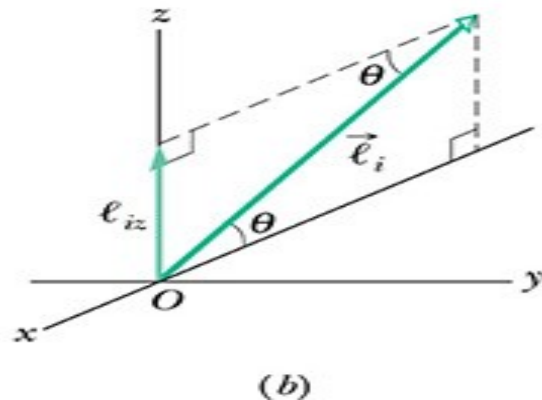


Angular momentum of mass element:

$$dl_i = r_i v_i dm_i = \omega r_i^2 dm$$

$r_i$  is the distance to  $dm$  and is perpendicular to  $v_i$

$$L = \int dl_i = \omega \int r_i^2 dm = I\omega$$



Rigid Body, Fixed Axis!

and obviously:

$$\frac{d}{dt} L = \tau_{\text{net}}$$

# Conservation of Angular Momentum

---

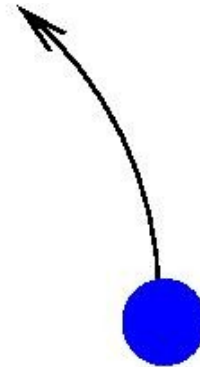
If the **component** of the **net external torque** on a system along a certain axis **is zero**, then the **component** of the **angular momentum** along that axis is **conserved**.

- It is a conservation law for each individual component of the angular momentum, not the magnitude.
- Similar to the conservation law for the components of the linear momentum.

# Conservation of Angular Momentum

---

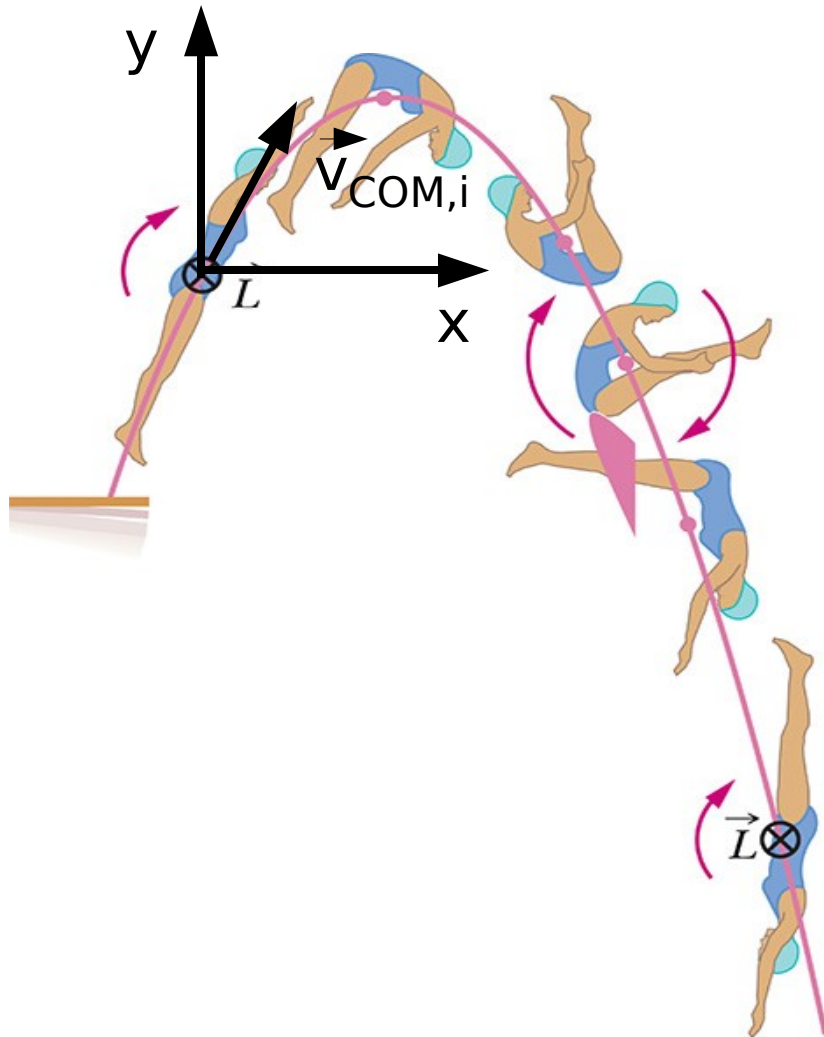
Example:



Sun-Earth system  
rotating around  
their Center of mass.

No torques:  
Angular momentum  
is conserved

# Conservation of Angular Momentum



Motion can be split into:

- Center of mass motion
- Rotation around center of mass

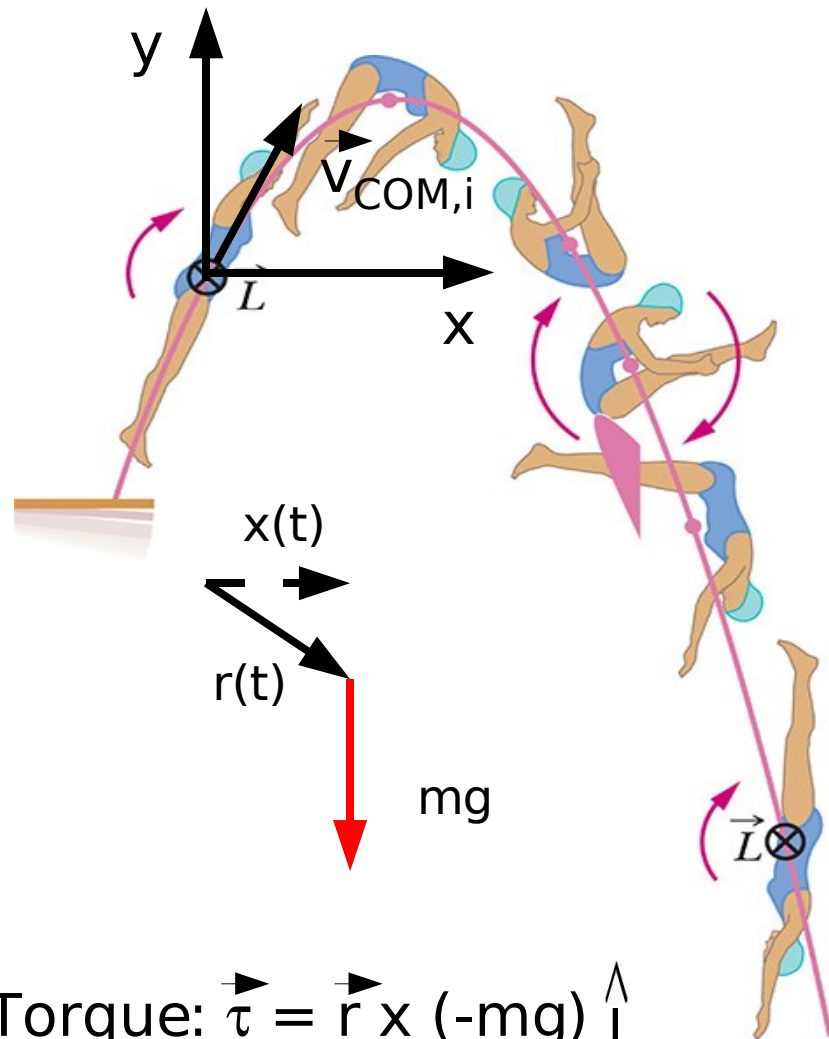
Initial conditions:

- Initial velocity  $\vec{v}_{COM,i}$
- Initial rotation  $\omega_i$

Initial momenta:

- Linear momenta  $p_x, p_y$
- Angular momenta (in z-axis):
  - Rotation around COM  $L_{int}$
  - COM motion  $L_{ext}$

# Conservation of Angular Momentum



The center of mass travels along a parabola

$$x(t) = v_{ix} t$$

$$y(t) = v_{iy} t - 0.5gt^2$$

$$\mathbf{v}(t) = v_{ix} \hat{i} + (v_{iy} - gt) \hat{j}$$

$$\vec{L}_{\text{ext}} = m(\vec{r} \times \vec{v}) = -0.5mgt^2 v_{ix} \hat{k}$$

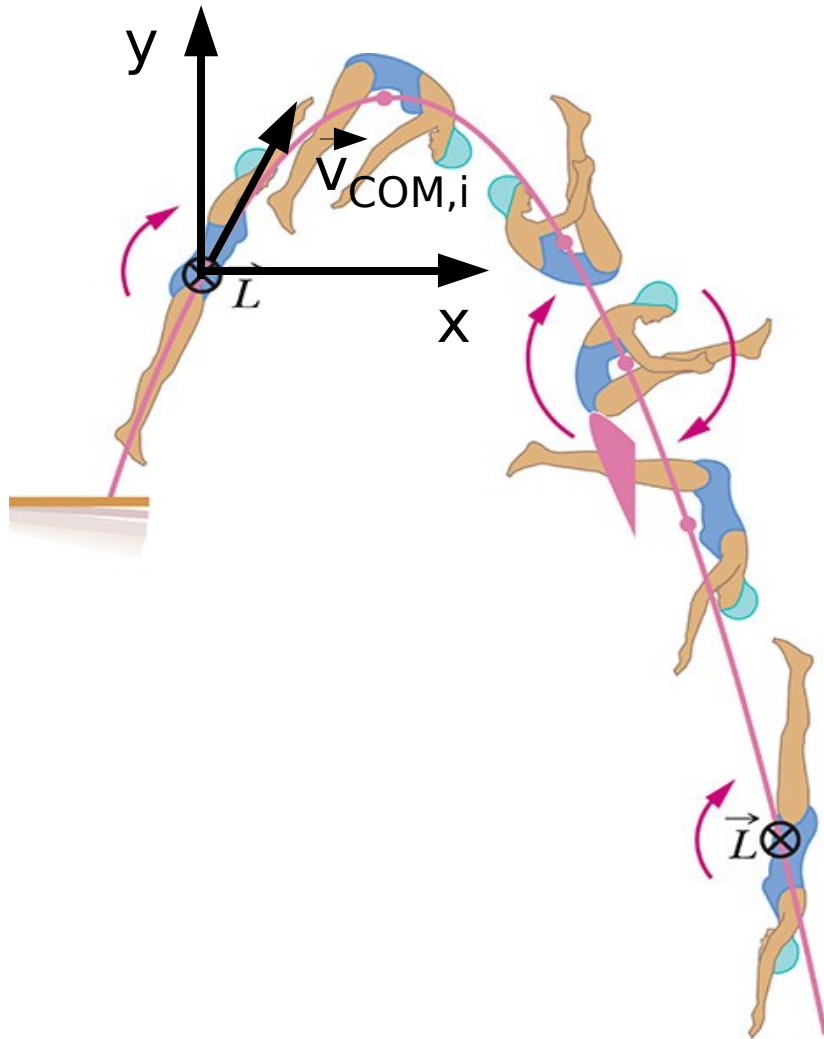
$$\text{Torque: } \vec{\tau} = \vec{r} \times (-mg) \hat{j}$$

$$\tau = -mg x(t) \hat{k} = -mgv_{ix} t \hat{k} = \frac{d}{dt} \vec{L}_{\text{ext}}$$

The external force (gravity) changes the external angular momentum



# Conservation of Angular Momentum



In addition: The jumper is rotating around its center of mass.

->  $L_{\text{int}}(t=0) < 0$  (clock wise)

Only force (gravity) works on center of mass.

-> No lever arm to change internal rotation

->  $L_{\text{int}}$  conserved

-> Reduce  $I$ ,  $\omega$  has to go up.

# Hitt

---

A beetle rides the rim of a small rotating disk.  
If the beetle crawls towards the center of the disk, do the following decrease, increase, or stay the same:  
rotational inertia, angular momentum, angular speed

A: same, increase, increase

B: decrease, same, increase

C: decrease, increase, same

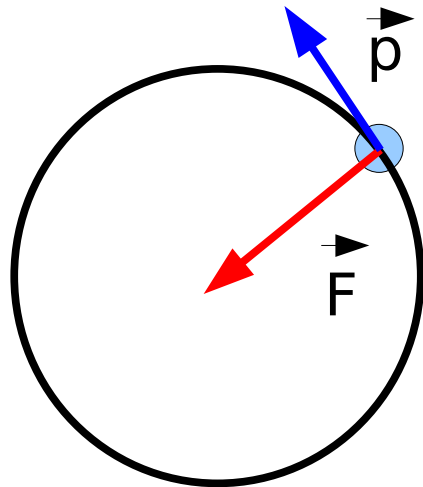
D: same, decrease, increase

E: increase, same, decrease

# Precession of a Gyroscope

---

Recall:



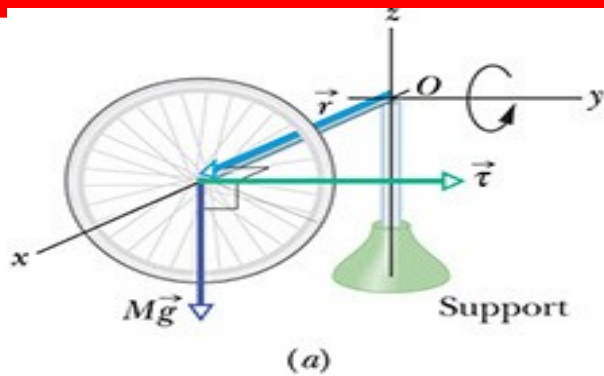
$$\frac{d}{dt} \vec{p} = \vec{F}_{\text{net}}$$

when  $\vec{F}$  is orthogonal to  $\vec{p}$   
it only changes the direction  
of  $p$  and not its magnitude

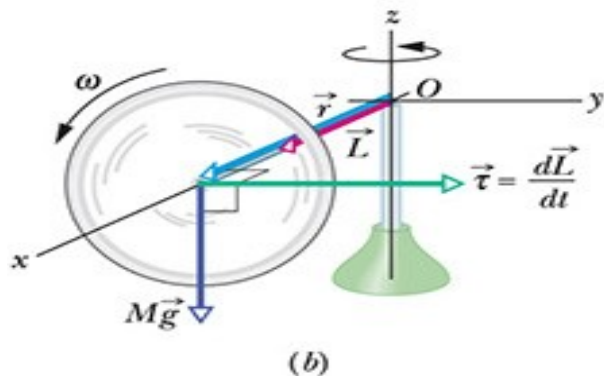
$$\frac{d}{dt} \vec{L} = \vec{\tau}_{\text{net}}$$

Something similar should happen  
when  $\vec{\tau}_{\text{net}}$  is perpendicular to  $\vec{L}$

# Precession of a Gyroscope

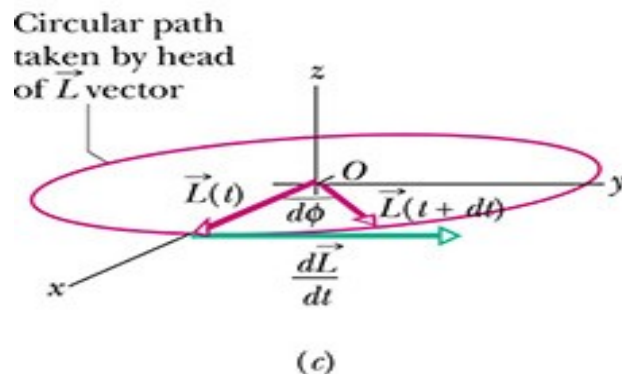


$$\frac{d}{dt} \vec{L} = \vec{\tau}_{\text{net}}$$

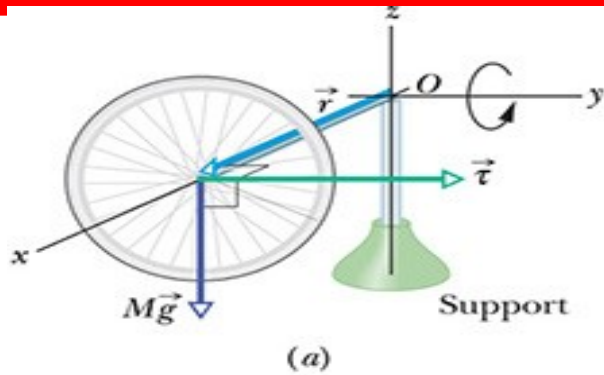


Spinning wheel supported at the end of a stick

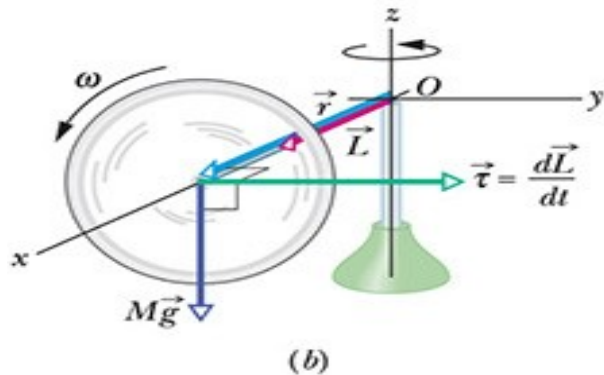
Gravity points down  
Torque is going to the side perpendicular to the angular momentum  
-> Precession



# Precession of a Gyroscope



$$\frac{d}{dt} \vec{L} = \vec{\tau}_{\text{net}}$$

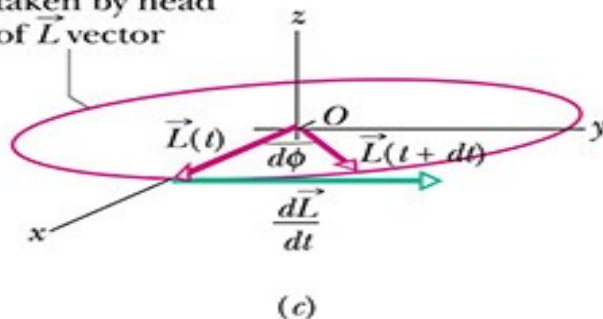


$$\tau = Mgr$$

$$dL = \tau dt = Ld\phi$$

$$d\phi = dL/L = Mgr dt / I\omega$$

Circular path  
taken by head  
of  $\vec{L}$  vector

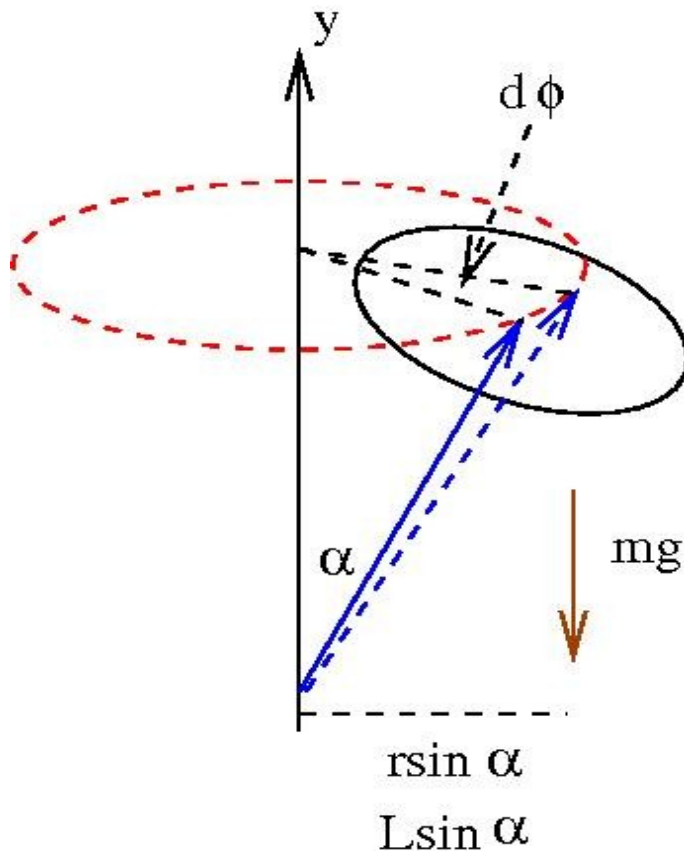


$$\frac{d}{dt} \phi = \Omega = \frac{Mgr}{I\omega}$$

precession  
rate

# Precession of a Gyroscope

What happens if  $\vec{L}$  is not horizontal?



The lever arm is now only  $r \sin \alpha$  :

$$\tau = Mg r \sin \alpha$$

But the change  $dL$  is also only scaling with  $L \sin \alpha$  :

$$dL = \tau dt = L \sin \alpha d\phi$$

$$\frac{d}{dt} \phi = \Omega = \frac{Mgr \sin \alpha}{I \omega \sin \alpha} \quad \text{precession rate}$$

Independent of  $\alpha$

# Example

---

Problem 68: A top spins at 30rev/s about an axis that makes an angle of 30 deg with the vertical. The mass of the top is 0.5kg, the rotational inertia is 0.0005kgm<sup>2</sup>, and its center of mass is 4cm from the pivot point.

If the spin is clockwise from an overhead view, what are the

- a) precession rate?
- b) direction of the precession as viewed from overhead?

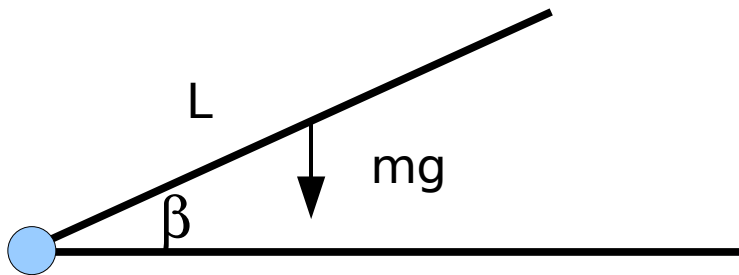
$$\Omega = \frac{Mgr}{I\omega}$$

$$\omega = (2\pi)30 \text{ rad/s}$$



# Example Problem

---



Two boards are connected such that they can rotate around each a common pivot point.

One board lies flat on the ground. What is the acceleration of the end of the second board as a function of  $\beta$  ?

$$\tau = 0.5Lmg \cos\beta$$

$$\tau = I\alpha \quad I = I_{\text{com}} + ML^2/4 = ML^2/3$$

$$\alpha = 3MgL\cos\beta / 2ML^2 = 1.5 g \cos\beta / L$$

$$a = \alpha L = 1.5 g \cos\beta \quad \text{Larger than } g \text{ for small } \beta$$