

Announcements

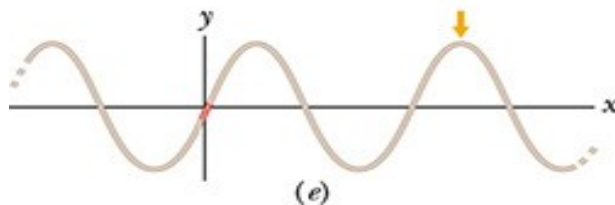
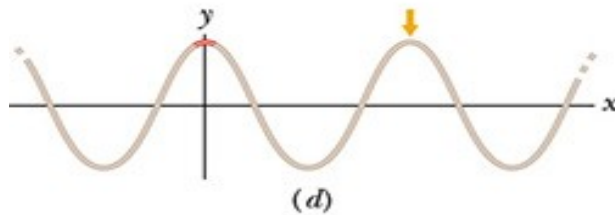
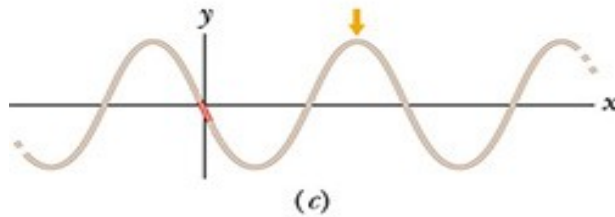
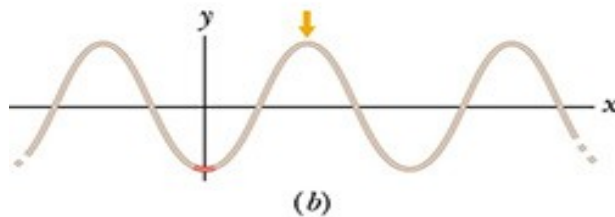
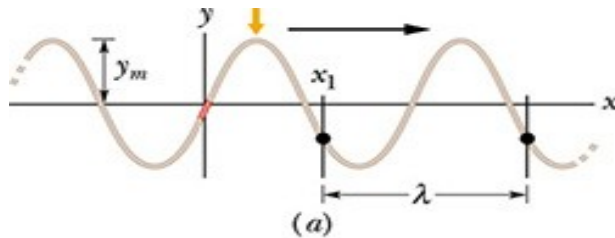
Exam 3 Review:

Sunday, December 7th 4:00pm-6:00pm in 1001 NPB (here)

Teacher Evaluations:

Wednesday, December 10th in class. In addition to the evaluations the students will be asked 3 HITT questions about the class. If they come and do the evaluations and answer the 3 questions they will get 6 HITT points! (And all possible answers are 'correct')

Chapter 16: Waves I



$$\sin(kx - \omega t)$$

Periodic function in space and time:

$$\sin(kx - \omega t) = \sin[k(x + \lambda) - \omega t]$$

$$k = 2\pi/\lambda$$

$$\sin(kx - \omega t) = \sin[kx - \omega(t + T)]$$

$$\omega = 2\pi/T$$

➡ $f = 1/T$ f : Frequency

➡ $v = \omega/k = f\lambda$ Velocity of sin wave

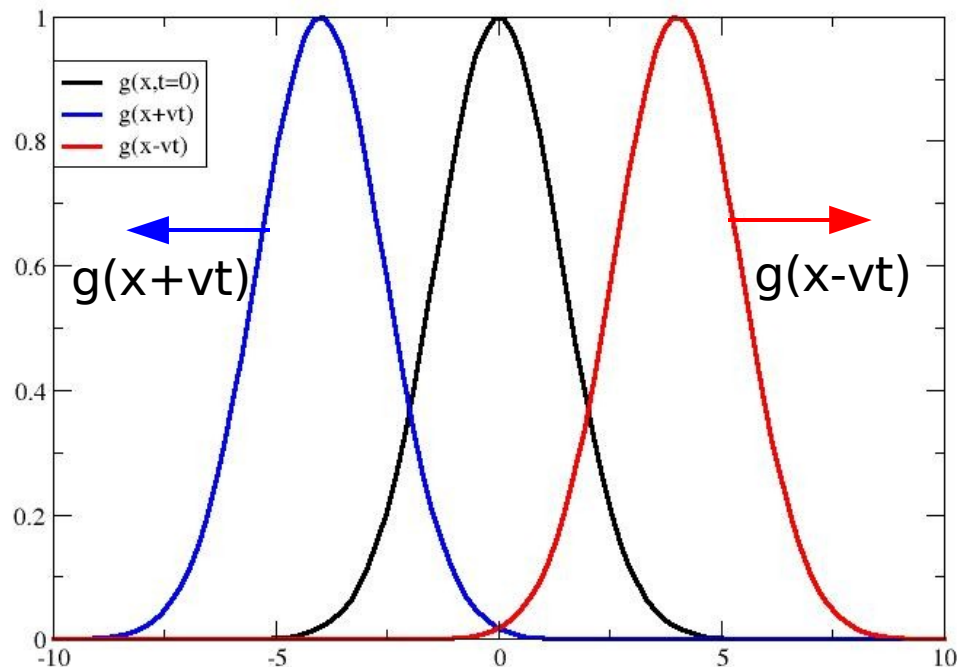
(more precise: phase velocity,
but we only discuss one velocity)

Chapter 16: Waves I

General wave:

$$y(x,t) = h(kx \pm \omega t) = g(x \pm vt)$$

$x - vt$: Propagates in the positive x-direction
 $x + vt$: Propagates in the negative x-direction



Examples for: $h(kx \mp \omega t) = g(x \mp vt)$

$$y(x,t) = 4x - 3t = 4\left(x - \frac{3}{4}t\right)$$

$$y(x,t) = 2e^{-(2x-4t)^2} = 2e^{-4(x-2t)^2}$$

$$y(x,t) = 3(x-2t)e^{-4(x-2t)^2}$$

Not a wave:

$$y(x,t) = 3(x-3t)e^{-4(x-2t)^2}$$

See plots in last lecture notes

Wave on a string:

The tension pulling to the left on a small element of length dx is equal in amplitude but has a slightly different direction than the tension pulling to the right (because of curvature, 2nd derivative $\neq 0$).

$$y - \text{component:} \quad T_{2y} + T_{1y} = a_y dm \quad T_{2y}, T_{1y} \text{ have opposite signs}$$

$$\text{using } dm = \mu dx \quad T_{2y} + T_{1y} = a_y \mu dm = \frac{d^2 y}{dt^2} \mu dx$$

gives the acceleration of the small element in the y-direction.

The difference in direction is so small that the x-components of the two tensions are assumed to be equal (change with $\cos \theta$).

The direction of the tension is tangential to the string:

$$\frac{T_{2y}(x+dx)}{T_{2x}} = \frac{dy}{dx} = \frac{y(x+dx) - y(x)}{dx} \quad \frac{T_{1y}(x-dx)}{T_{1x}} = -\frac{dy}{dx} = -\frac{y(x) - y(x-dx)}{dx}$$

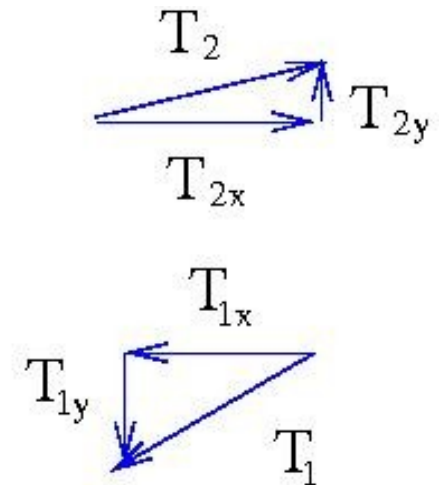
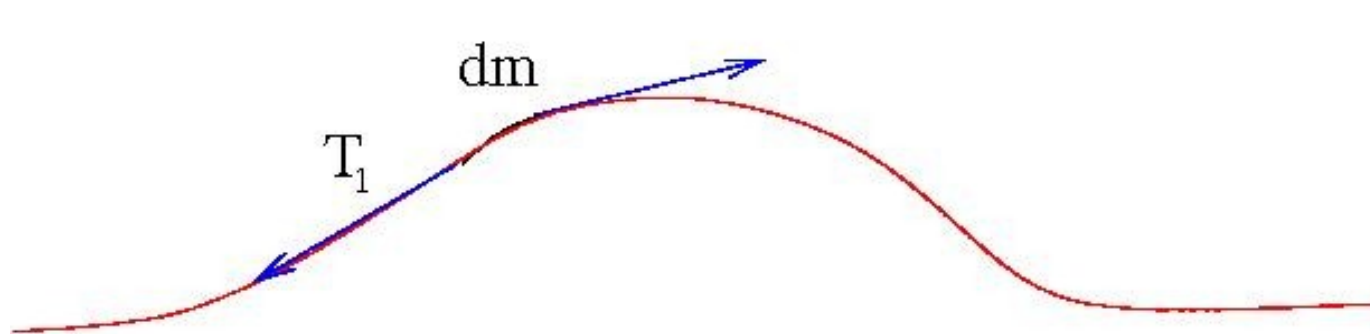
$$\Rightarrow \frac{T_{2y}(x+dx)}{T_{2x}} + \frac{T_{1y}(x-dx)}{T_{1x}} = \frac{T_{2y}(x+dx) + T_{1y}(x-dx)}{T} = \frac{y(x+dx) - 2y(x) + y(x-dx)}{dx}$$

$$\Rightarrow \mu \frac{d^2 y}{dt^2} = T \frac{y(x+dx) - 2y(x) + y(x-dx)}{dx^2} = T \frac{d^2 y}{dx^2}$$

$$\Rightarrow \frac{d^2 y}{dt^2} = \frac{T}{\mu} \frac{d^2 y}{dx^2} = v^2 \frac{d^2 y}{dx^2} \quad \text{Wave equation for small amplitudes}$$

$$v = \sqrt{\frac{T}{\mu}} \quad \text{Wave velocity} \quad \text{Solution: } y(x,t) = h(x \pm vt)$$

Chapter 16: Waves I



Key idea:

Tension is pulling in non-parallel directions on curved string as the slope changes.

For ideal longitudinal wave:

$T = T_{1x} = -T_{2x}$ (small amplitude \Leftrightarrow small angles, difference will scale with $\cos\theta$), compensate each other and no acceleration in x-direction (propagation direction)

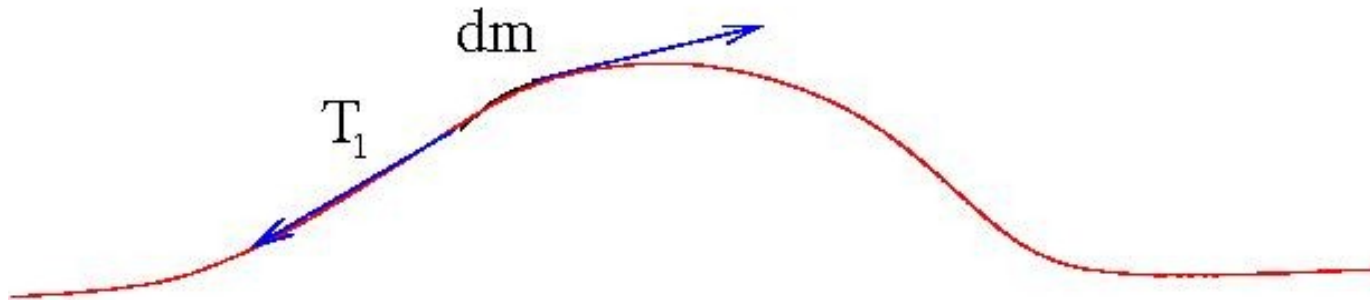
The y-components are different in first order in θ and cause an acceleration in the y-direction: $T_{1y} + T_{2y} = a_y dm = a_y \mu dx$ (plus: T_{1y} will be negative in the above case)

The y-components of the tension are also proportional to the slope of the rope on both sides of the string segment: $T_{1y} = -T dy/dx$ at (x) $T_{2y} = T dy/dx$ at $(x+dx)$

Adding the two tensions and dividing by dx gives $(T_{1y} + T_{2y})/dx = T d^2y/dx^2$

Putting this together gives wave equation.

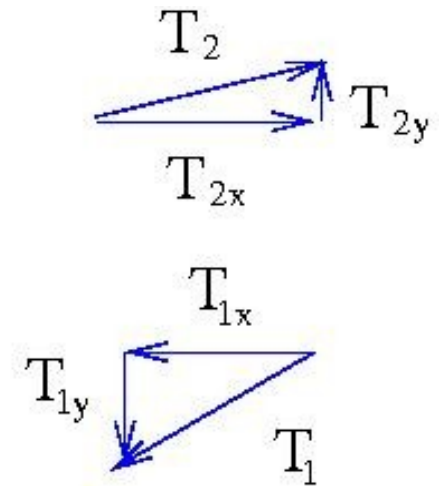
Chapter 16: Waves I



For you important is:

$$\frac{d^2y}{dt^2} = \frac{T}{\mu} \frac{d^2y}{dx^2} = v^2 \frac{d^2y}{dx^2} \quad \text{Wave equation}$$

$$v = \sqrt{\frac{T}{\mu}} \quad \text{Wave velocity} \quad \text{Solution: } y(x,t) = h(x \mp vt)$$



Check: Substitute: $x' = x \mp vt$ $\frac{dx'}{dx} = 1, \frac{dx'}{dt} = \mp v$

$$\frac{d}{dt} \left[\frac{dh(x')}{dx'} \frac{dx'}{dt} \right] = \frac{d}{dt} \left[\mp v \frac{dh(x')}{dx'} \right] = \mp v \frac{d}{dx'} \left[\frac{dh(x')}{dt} \right] = \mp v \frac{d}{dx'} \left[\frac{dh(x')}{dx'} \frac{dx'}{dt} \right] = v^2 \frac{d^2h(x')}{dx'^2}$$

$$v^2 \frac{d}{dx} \frac{dh(x \mp vt)}{dx} = v^2 \frac{d}{dx} \left[\frac{dh(x')}{dx'} \frac{dx'}{dx} \right] = v^2 \frac{d^2h(x')}{dx'^2} \quad \text{q.e.d.}$$

Chapter 16: Waves I

Problem:

Assume a string of length 200m, a mass of $m = 6\text{kg}$, stretched out by a tension of $T = 500\text{N}$. If this string is excited by an external force at one end with a frequency of $f = 100\text{Hz}$, what is the wavelength of the traveling wave?

Plan:

To get the wavelength: $\lambda = \frac{v}{f}$ we need the speed

To get the speed: $v = \sqrt{\frac{T}{\mu}}$ we need μ

General velocity
for all types of waves
(not only harmonic)

$$\mu = \frac{m}{L} = \frac{6\text{kg}}{200\text{m}} = 3 \times 10^{-2} \frac{\text{kg}}{\text{m}} \quad v = \sqrt{\frac{500\text{Nm}}{3 \times 10^{-2}\text{kg}}} = 129 \frac{\text{m}}{\text{s}}$$

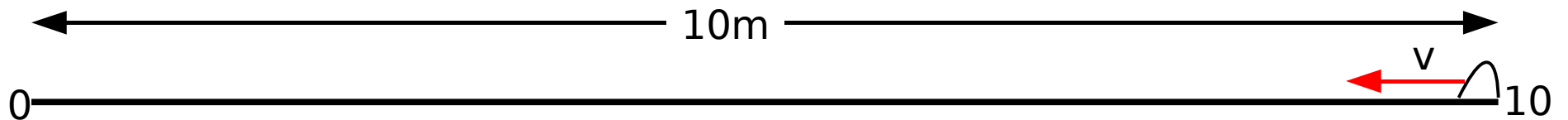
$$\lambda = \frac{129 \frac{\text{m}}{\text{s}}}{100\text{Hz}} = 1.29\text{m}$$

In sinusoidal (harmonic)
waves: λ, f, v are coupled
(In non-sinusoidal waves:
we don't have a wavelength)

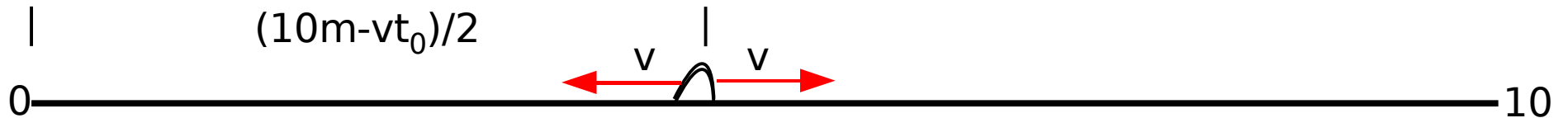
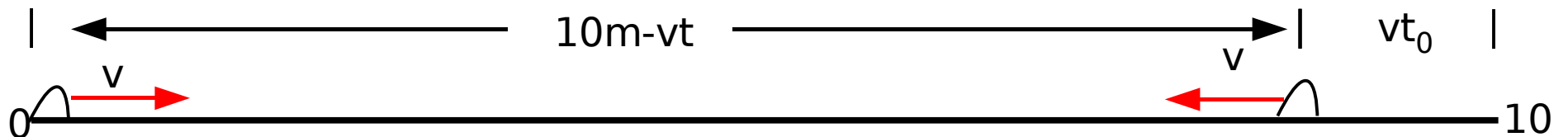
Chapter 16: Waves I

A 100g wire is held under a tension of 250N with one end at $x=0$ and the other at $x=10.0\text{m}$. At time $t=0$, pulse 1 is sent along the wire from the end at $x=10.0\text{m}$. At time $t=30\text{ms}$, pulse 2 is sent along the wire from the end at $x=0$. Where do the pulses meet?

$t=0$



$t_0=30\text{ms}$



$$v = \sqrt{\frac{T}{\mu}} = \sqrt{\frac{TL}{m}} = \sqrt{\frac{250\text{N} \cdot 10.0\text{m}}{0.1\text{kg}}} = 158 \frac{\text{m}}{\text{s}}$$

$$vt_0 = 158 \frac{\text{m}}{\text{s}} \cdot 0.03\text{s} = 4.74\text{m}$$

$$d = \frac{10\text{m} - 4.74\text{m}}{2} = 2.63\text{m}$$

Chapter 16: Waves I

Energy of a wave traveling along a string:

➤ Kinetic energy

- as the mass elements dm move up and down
- max when string element passes through $y=0$
- zero when mass element reaches the max displacement $y=y_m$

➤ Elastic potential energy

- as the wave stretches the string
- max at $y=0$ where the string is stretched the most
- zero when mass element reached the max displacement $y=y_m$

➤ Energy transport

- both energy forms are max when $y=0$
- no energy at locations where $y=y_m$
- the tension in the string where $y=0$ pulls the parts where $y\sim y_m$ back, causing them to accelerate.
- the parts where $y=0$ continue to move up (or down) with their initial velocity until the tension which is now building up in the other areas stops them.

Chapter 16: Waves I

To get the potential energy of string element, we start with:

$$U = \frac{1}{2}ks^2 \quad \frac{dU}{ds} = ks = -T \quad s : \text{additional length after stretching string}$$

$$\Rightarrow dU = -Tds \quad ds = \sqrt{[y(x+dx) - y(x)]^2 + dx^2} - dx \quad \text{additional length of string}$$

$$\sqrt{[y(x+dx) - y(x)]^2 + dx^2} = dx \sqrt{1 - \left(\frac{dy}{dx}\right)^2} \approx dx \left[1 - \frac{1}{2} \left(\frac{dy}{dx}\right)^2\right] \quad \text{for small derivatives}$$

$$\Rightarrow ds = -\frac{1}{2} \left(\frac{dy}{dx}\right)^2 dx \quad \Rightarrow dU = T \frac{1}{2} \left(\frac{dy}{dx}\right)^2 dx$$

$$\text{Recall: } T = v^2\mu \quad \Rightarrow \quad \frac{dU}{dt} = \frac{1}{2}v^3\mu \left(\frac{dy}{dx}\right)^2 \left[= \frac{1}{2}\mu v \omega^2 y_m^2 \cos^2(kx - \omega t) \text{ for harmonic w.} \right]$$

Chapter 16: Waves I

kinetic energy of line element: $dK = \frac{1}{2}u^2 dm = \frac{1}{2}u^2 \mu dx$ u : transverse speed

General: $u = \frac{dy(x-vt)}{dt} = -v \frac{dy}{dx} [= -\omega y_m \cos(kx - \omega t)]$ for harmonic waves]

transverse velocity = wave speed \times slope of pulse form

$$dK = \frac{1}{2}v^2 \left(\frac{dy}{dx} \right)^2 \mu dx \left[= \frac{1}{2}\omega^2 y_m^2 \cos^2(kx - \omega t) \mu dx \quad \text{for harmonic waves} \right]$$

The rate at which kinetic energy passes through a line element:

$$\frac{dK}{dt} = \frac{1}{2}v^2 \left(\frac{dy}{dx} \right)^2 \mu \frac{dx}{dt} = \frac{1}{2}v^3 \mu \left(\frac{dy}{dx} \right)^2 \left[= \frac{1}{2}\mu v \omega^2 y_m^2 \cos^2(kx - \omega t) \text{ for harmonic w.} \right]$$

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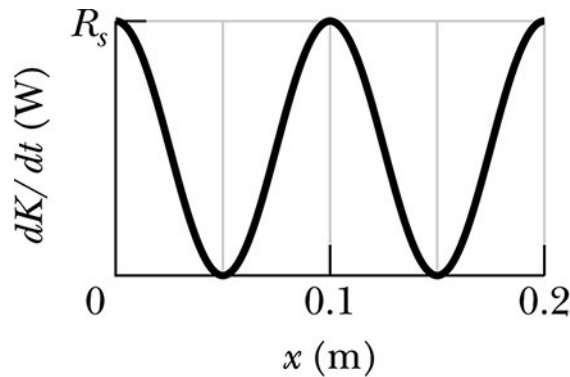
$$\frac{dK}{dt} + \frac{dU}{dt} = v^3 \mu \left(\frac{dy}{dx} \right)^2 \quad [= \mu v \omega^2 y_m^2 \cos^2(kx - \omega t) \text{ for harmonic w.}]$$

The average power (averaged over several wave lengths) is the average rate with which energy propagates through the string

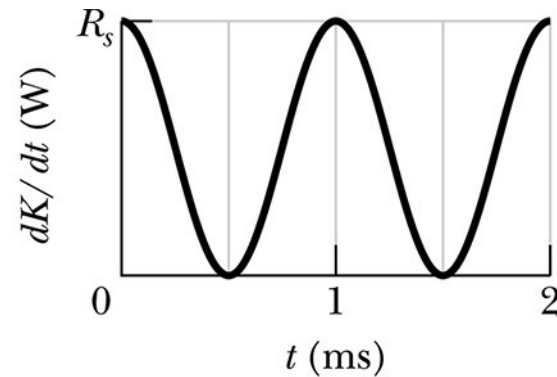
$$P_{avg} = v^3 \mu \left(\frac{dy}{dx} \right)^2 \quad \left[= \frac{1}{2} \mu v \omega^2 y_m^2 \text{ for harmonic waves} \right]$$

As the book only discusses energy in harmonic waves, you will only need the expressions for harmonic waves for exams and homeworks.

Chapter 16: Waves I



(a)



(b)

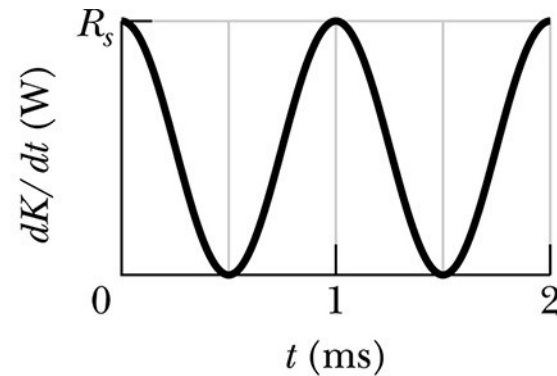
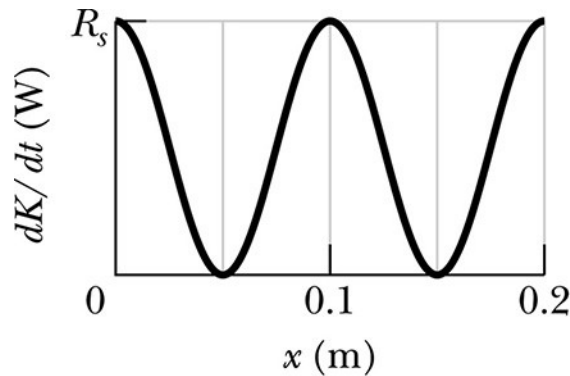
The graphs show the rate dK/dt at which kinetic energy passes through string elements:

- a) at a particular time as a function of distance x
- b) at a particular location as a function of time

The amplitude $R_s=10\text{W}$ in both graphs. The linear mass density is $\mu=2\text{g/m}$

What is the amplitude of the wave?

Chapter 16: Waves I



The rate of change in kinetic energy of a harmonic wave is:

$$\frac{dK}{dt}_{max} = \frac{1}{2} \mu v \omega^2 y_m^2 = 10W \quad \text{with: } \mu = 2 \times 10^{-3} \frac{\text{kg}}{\text{m}}$$

From the graphs we see:

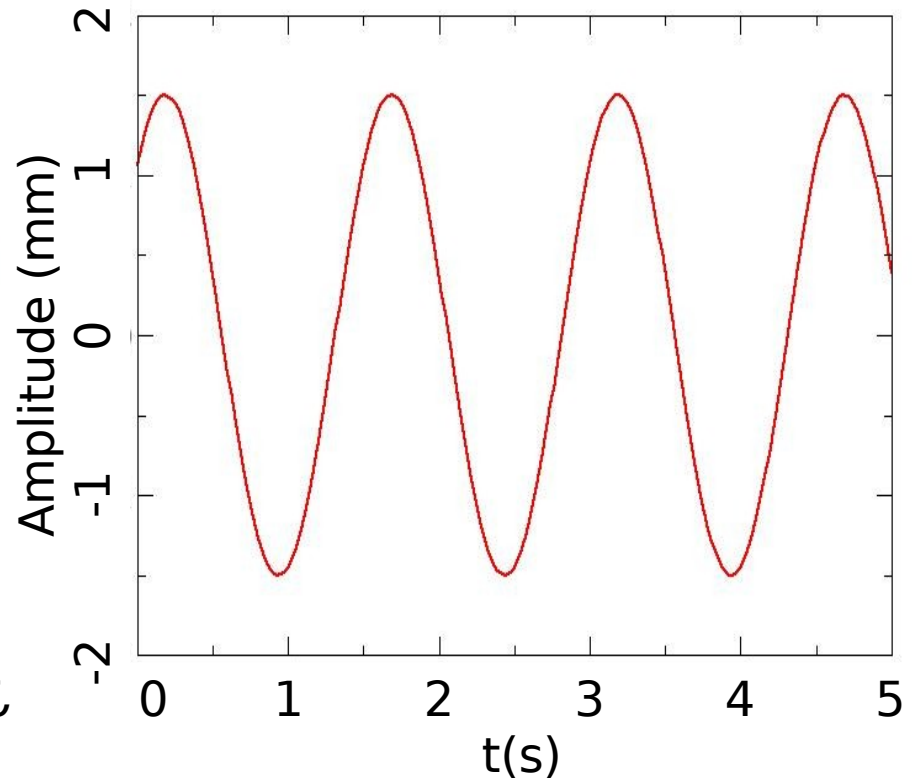
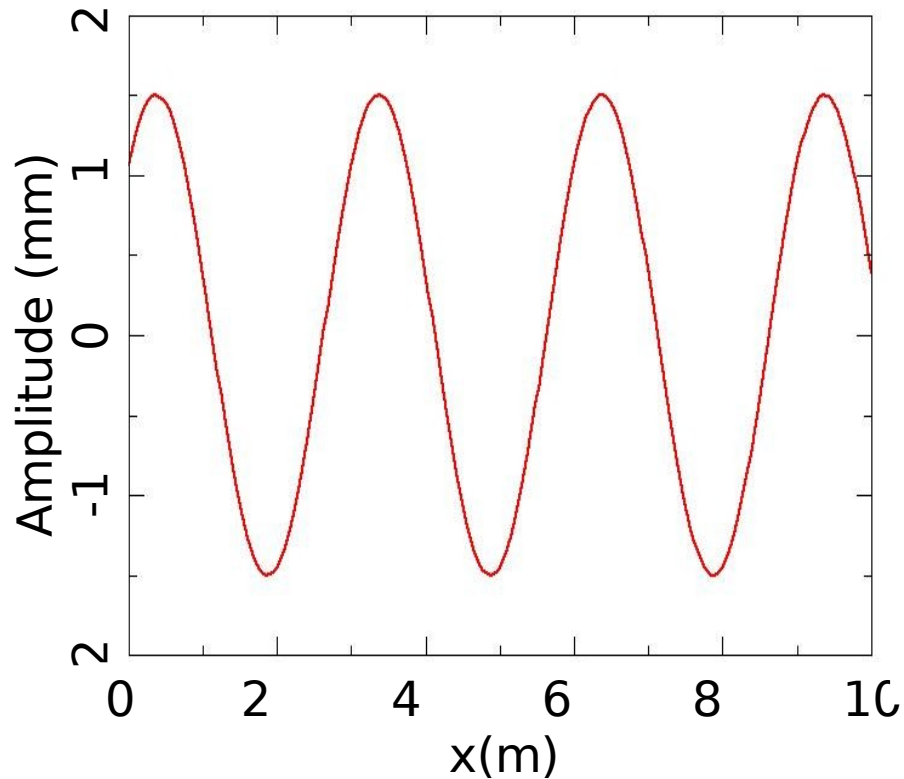
$$\lambda = 0.2 \text{ m} \quad \text{This is a } \cos^2 \text{ -graph!} \quad f = \frac{1}{T} = \frac{1}{2\text{ms}} = 500\text{Hz}$$

Use:

$$v = \lambda f \quad \omega = 2\pi f \quad \Rightarrow \quad 10W = 2\pi^2 \mu \lambda f^3 y_m^2$$

$$\Rightarrow \quad y_m = \sqrt{\frac{10W}{2\pi^2 \mu \lambda f^3}} = 3.2 \text{ mm}$$

HITT 1



The two graphs show a harmonic wave at $t=0$ in space and at $x=0$ in time.

What is the speed of the wave?

Graph 1: Wavelength $\lambda=3\text{m}$

Graph 2: Period $T=1.5\text{s}$ \rightarrow Frequency $f=2/3\text{Hz}$

Velocity: $v=\lambda f=3 \cdot 2/3 \text{ m/s}=2\text{m/s}$

Superposition principle

Formal: Let $y_1(x, t)$ and $y_2(x, t)$ be solutions of the wave equation. Then

$$y'(x, t) = c_1 y_1(x, t) + c_2 y_2(x, t)$$

is also a solution.

Proof:

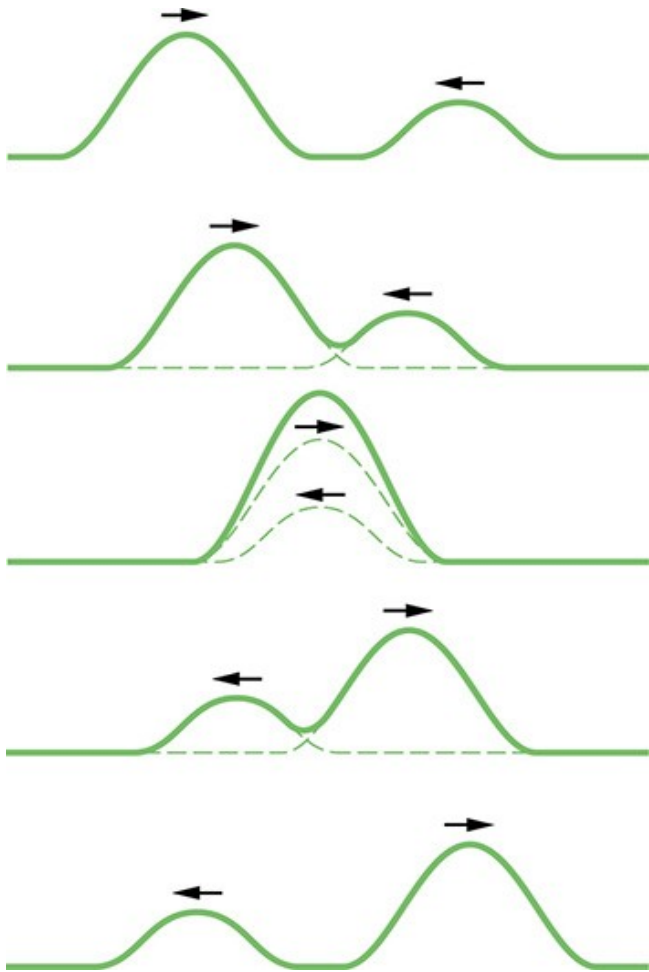
$$\frac{d^2 y'}{dx^2} = c_1 \frac{d^2 y_1}{dx^2} + c_2 \frac{d^2 y_2}{dx^2} = \frac{c_1}{v^2} \frac{d^2 y_1}{dt^2} + \frac{c_2}{v^2} \frac{d^2 y_2}{dt^2} = \frac{1}{v^2} \frac{d^2 (c_1 y_1 + c_2 y_2)}{dt^2} = \frac{1}{v^2} \frac{d^2 y'}{dt^2}$$

What are the consequences?

If two waves travel in the same medium (along the same string), they add algebraically to a resultant net wave.

They travel independent of each other, don't influence each. In other words: They can be treated completely independent from each other (assuming that all approximations we used to derive the wave equation are still valid).

Superposition principle



Example:
Two waves propagating in opposite directions:

$$y_1(x, t) = y_{1m}h(x - vt)$$

$$y_2(x, t) = y_{2m}h(x + vt)$$



$$y'(x, t) = y_{1m}h(x - vt) + y_{2m}h(x + vt)$$

Note: This is not a single wave
because it has two different velocities