

Wave Packets:

Consider 
$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2 \partial^2 \psi}{2m \partial x^2} .$$

This has solutions

$$\begin{aligned} \psi(x, t) &\propto e^{i(kx - \omega t)} && \text{right moving} \\ &\propto e^{i(kx + \omega t)} && \text{left moving.} \end{aligned}$$

Check: 
$$i\hbar \frac{\partial}{\partial t} e^{i(kx - \omega t)} = \hbar \omega e^{i(kx - \omega t)}$$

$$-\frac{\hbar^2 \partial^2}{2m \partial x^2} e^{i(kx - \omega t)} = \frac{\hbar^2 k^2}{2m} e^{i(kx - \omega t)}$$

$$\rightarrow \hbar \omega = \frac{\hbar^2 k^2}{2m} \checkmark$$

$|\psi|^2 \propto |e^{i(kx - \omega t)}|^2 = 1$  has no time dependence.

To get time dependence we must add states of different energies.

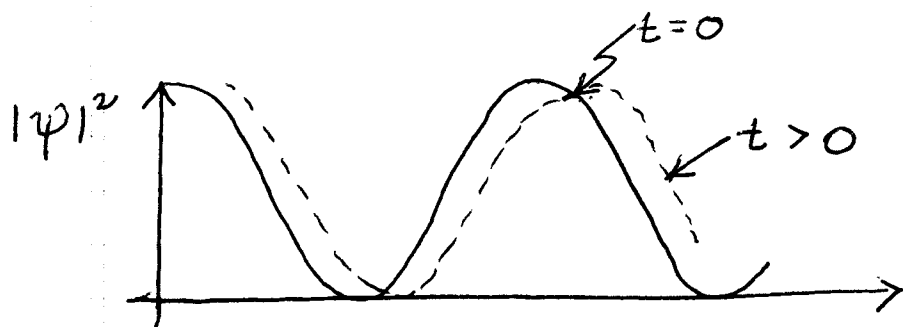
Let  $\psi_1(x, t) = e^{i(k_1 x - \omega_1 t)}$  , where  $\hbar \omega_1 = \frac{\hbar^2 k_1^2}{2m}$

and  $\psi_2(x, t) = e^{i(k_2 x - \omega_2 t)}$  , where  $\hbar \omega_2 = \frac{\hbar^2 k_2^2}{2m}$ .

$$|\psi_1(x, t) + \psi_2(x, t)|^2 = 1 + 1 + e^{i(k_1 - k_2)x} e^{-i(\omega_1 - \omega_2)t} + e^{i(k_2 - k_1)x} e^{-i(\omega_2 - \omega_1)t}$$

$$= 2 + 2 \cos((k_1 - k_2)x - (\omega_1 - \omega_2)t)$$

also a solution



(group) velocity =  $\frac{\omega_1 - \omega_2}{k_1 - k_2}$

This moves to the right but is not normalized.  
Why stop at adding two plane waves:

$$\psi(x, t) = \int \frac{dk}{2\pi} g(k) e^{i(kx - \omega_k t)}, \quad \hbar\omega_k = \frac{\hbar^2 k^2}{2m}$$

This is also a solution. You may wish to write it as:

$$\psi(x, t) = \int \frac{dp}{2\pi\hbar} \psi(p, 0) e^{i(\frac{p}{\hbar}x - \omega_p t)}, \quad \text{where}$$

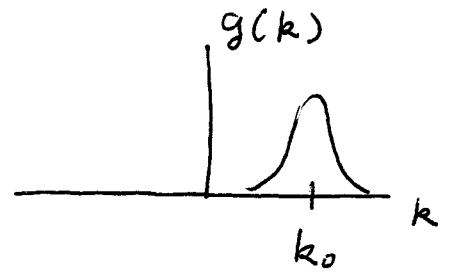
$$\left. \begin{aligned} \psi(x, 0) &= \int \frac{dp}{2\pi\hbar} \psi(p, 0) e^{i\frac{p}{\hbar}x} \\ \psi(p, 0) &= \int dx \psi(x, 0) e^{-i\frac{p}{\hbar}x} \end{aligned} \right\} \text{Fourier relations}$$

Normalization:

$$\begin{aligned} 1 &= \int dx |\psi(x, t)|^2 = \int \frac{dp}{2\pi\hbar} |\psi(p, 0) e^{-i\omega_p t}|^2 \\ &= \int \frac{dp}{2\pi\hbar} |\psi(p, 0)|^2 \end{aligned}$$

An example:

$$g(k) = \frac{\sqrt{2}}{(2\pi\alpha)^{1/4}} e^{-\alpha(k-k_0)^2}$$



$$|\psi(x,t)|^2 = \frac{1}{\sqrt{2\pi}} \frac{\alpha^{1/2}}{\sqrt{\alpha^2 + (\frac{\hbar t}{2m})^2}} \exp\left(-\frac{\alpha(x - \frac{\hbar k_0}{m}t)^2}{2(\alpha^2 + (\frac{\hbar t}{2m})^2)}\right).$$

See last year's lecture 3 for details.

This wave packet moves with velocity  $\hbar k_0/m$ , and the width changes with time.